Hadron yields and fluctuations at CERN SPS: system size dependence from p+p to Pb+Pb

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In collaboration with: V. Begun, V. Vovchenko, M.I. Gorenstein, H. Stoecker Based on 1711.07789 + 1811.10645



# **Outline**

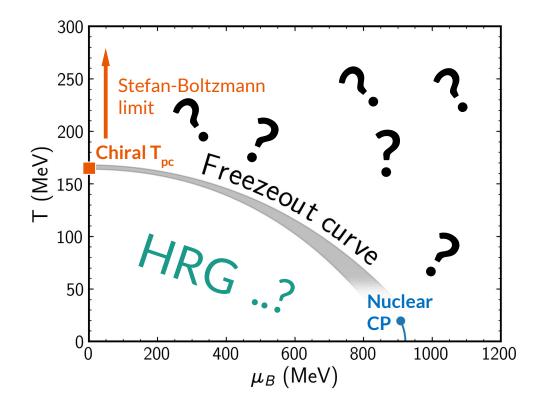
## What is done:

Systems: p+p, Be+Be, Ar+Sc, Pb+Pb Energies:  $E_{lab}$ =30(31) — 150(158) GeV Tools: HRG — Thermal-FIST, transport codes — UrQMD Observables: K<sup>+</sup>/ $\pi$ <sup>+</sup> ratio,  $\omega$ [N\_],  $\Omega$ [N\_,  $E_{n}$ ]

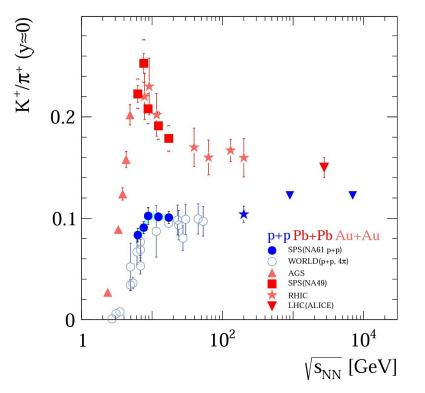
## **Contents:**

- 1. Why it's done;
- 2. CE HRG results for p+p, Be+Be, Ar+Sc, Pb+Pb;
- 3. Transport model for p+p, Be+Be, Ar+Sc, Pb+Pb:
  - a. How b=0 collisions are related to data;
  - b. Centrality effects on fluctuations;
  - c. Comparison with the NA61/NA49 data;
  - d. Centrality for p+p (!);
  - e.  $K^+/\pi^+$  ratio.

One of the main goals of Heavy lons physics is to study **phase structure** of QCD matter in bulk. However, yet there's still not so much known:



# **Observables motivation**



40 ω 0.5 35 - 10 30 2 (MeV) 20 25 ⊢ <sub>15</sub>′ 10 liquid gas 0.1 5 (b) 0.01 0 880 890 900 910 920 930 μ (MeV)

The horn structure in  $K^+/\pi^+$  serves as experimental evidence for the onset of deconfinement and quark-gluon plasma formation [1,2] Particle **fluctuations** reflect **critical behavior** of the system and so are useful measures to detect transitions between phases.

#### Non-monotonous fluctuations ⇒ criticality

[1] N. Abgrall et al., [NA61/SHINE Collaboration]: SPSC-SR-145. CERN-SPSC-2014-031

[2] M. Gazdzicki, M.I. Gorenstein, hep-ph/9803462.

[3] V. Vovchenko, D.V. Anchishkin, M.I. Gorenstein, and R.V. Poberezhnyuk, 1506.05763

# **Recent NA61/SHINE data**

**Onset of Fireball** — the rapid change of hadron production properties that start when moving from Be+Be to Ar+Sc collisions discovered in recent preliminary (checks are still in ()=^() 0.25 150A - 158A GeV/c process) NA61/SHINE 30A GeV/c ×() 0.25 (¢ µ/+µ results [1,2]. **NA49** Pb+Pb 0.2 Scaled variance *w* for negatively Pb+Pb charged particles N\_: **NA49** 0.15 0.15 Be+Be  $\omega[N_{-}] = \frac{\langle N_{-}^2 \rangle - \langle N_{-} \rangle^2}{\langle N_{-} \rangle}$ Be+Be WNM SHINE preliminary  $10^{2}$  $10^{3}$  $10^{3}$ 10<sup>2</sup> 10 10 <W> <W> Z 31.2 N 31.2 h at 30A GeV/c h at 150A-158A GeV/c What does theory have to do Be+Be WNM with this data? 1.1 1.1 Are there any effects not Poisson Poisson

Are there any effects not related to QCD physics but to **imperfectiontions of the experiment**? 

 Between the second s

 $\begin{array}{c} \text{S1.2} \\ 1.1 \\ 1$ 

[1] A. Aduszkiewicz et al. [NA61/SHINE Collaboration], CERN-SPSC-2017-038[2] M. Gazdzicki et al. [NA61/SHINE Collaboration], PoS CPOD2017 (2018) 012

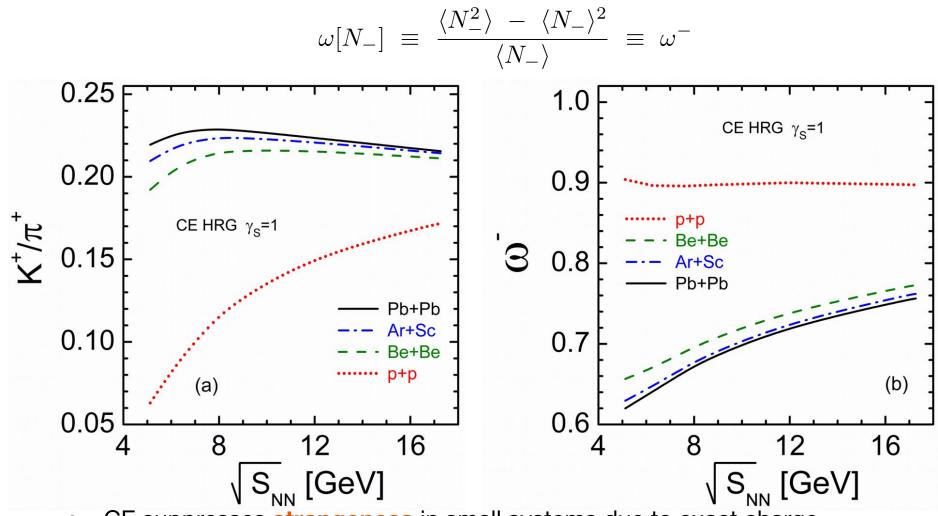
# Hadron Resonance Gas: chemical freeze-out curve

We use HRG within <u>Thermal-FIST</u> [1] package imposing **Canonical Ensemble** (baryon, electric, and strange charges conserved in every microstate), so: temperature T, volume V, and strangeness suppression  $\gamma_{s}$ are free parameters that are fixed from the data We obtain a continuous  $f(A+A, \sqrt{s_{NN}})$  freeze-out curve in CE: 0.16 12 (b) 0.15 Pb+Pb 0.15 9 **S** 0.14 **9** 0.13 Xe+La 0.10 R [fm] [fm<sup>-'</sup>] 6 Ar+Sc 0.12 0.05 3 Be+Be 0.11 (a) 0 0 12 16 12 12 8 8 16 8 16 4 4  $\sqrt{s_{NN}}$  [GeV]  $\sqrt{S_{NN}}$  [GeV]  $\sqrt{S_{NN}}$  [GeV]  $: \frac{T(\mu_B)}{\mu_B(\sqrt{s_{NN}})} = \frac{a - b \,\mu_B^2 - c \,\mu_B^4}{1 + e \,\sqrt{s_{NN}}}$  $R \equiv [3V/(4\pi)]^{1/3}$  and **T** is from [2,3]: Where canonical volume V can be found through baryonic density  $\rho_{R}$  (assumption)  $R = [3B/(4\pi\rho_B)]^{1/3}$ 

[1] V. Vovchenko and H. Stoecker, 1901.05249, see also V.Vovchenko talks on Tuesday + Wednesday
[2] J. Cleymans, H. Oeschler, K. Redlich, and S. Wheaton, hep-ph/0511094
[3] V. Vovchenko, V. V. Begun, and M. I. Gorenstein ,1512.08025

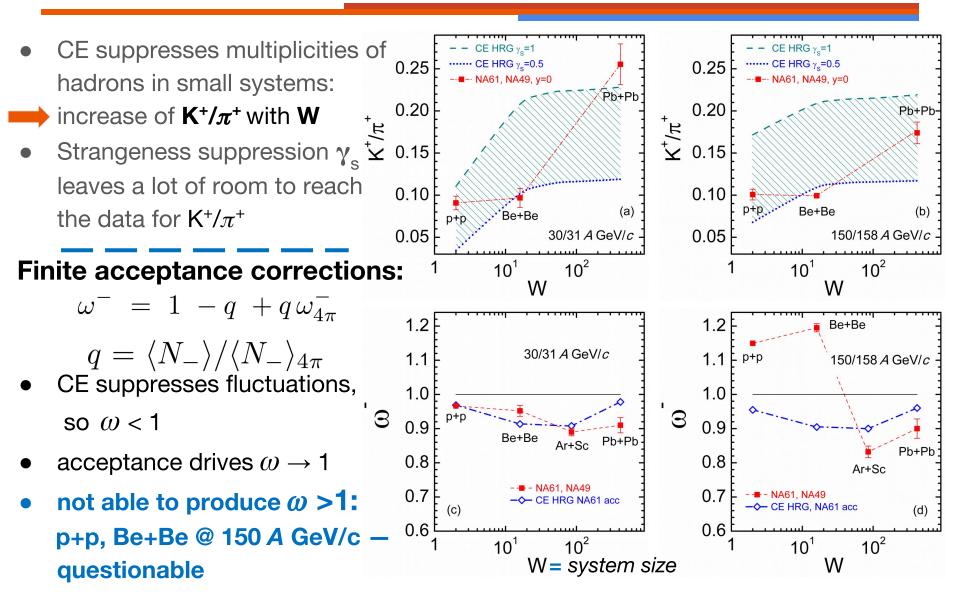
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# **HRG: results**



- CE suppresses strangeness in small systems due to exact charge conservation;
- CE suppresses fluctuations in large systems;

# **HRG: comparison with the data**



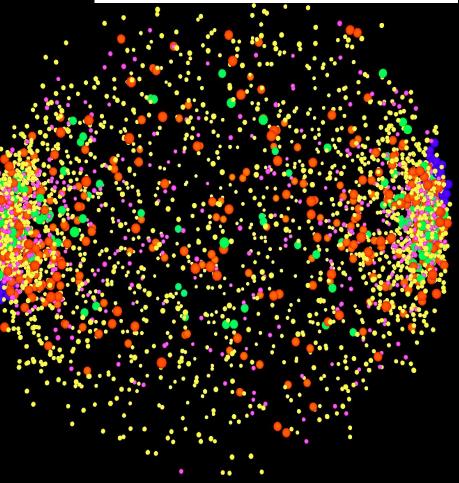
# **Microscopical simulations**

The Ultra relativistic Quantum Molecular Dynamics (UrQMD) model [1,2] is a microscopic transport theory based on the covariant propagation of all hadrons on classical trajectories in combination with stochastic binary scatterings, color string formation (by Pythia) and resonance decays.

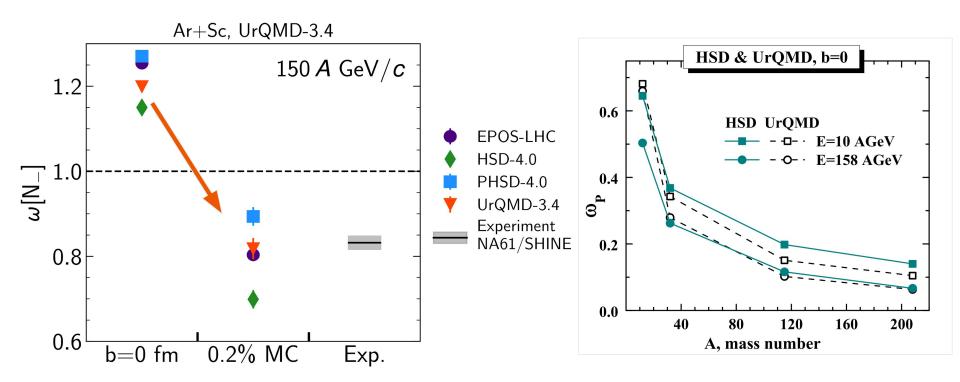
In this work: no hydro, no EoS effects, no mean fields, only cascade + strings.

[1] S. A. Bass et al., nuclth/9803035[2] M. Bleicher et al., hep-ph/9909407





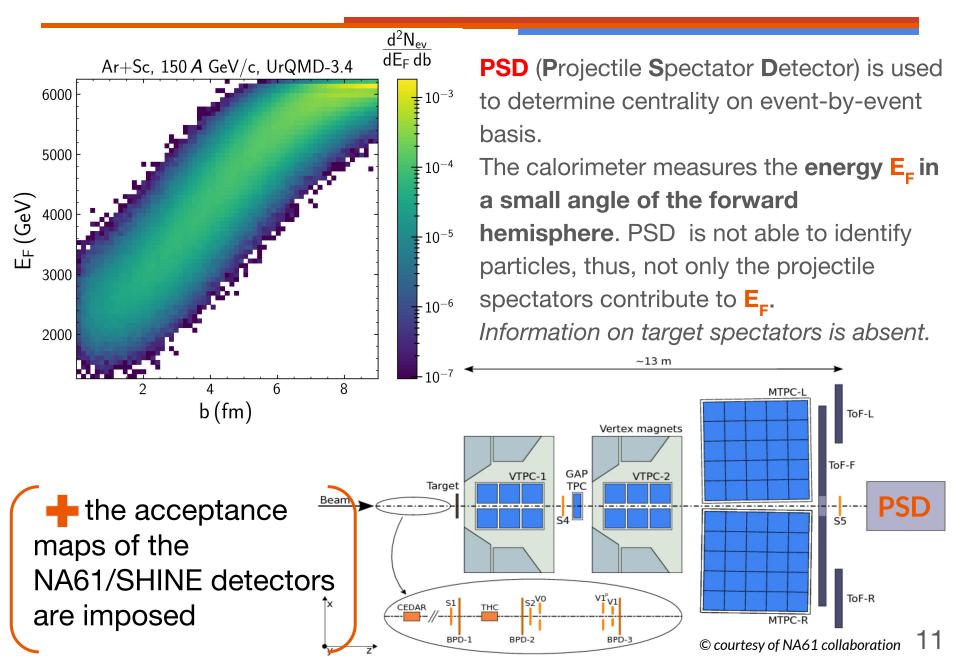
# Importance of centrality in intermediate systems



Most central events for intermediate systems cannot be approximated by b=0 fm collisions. In b=0 fluctuations of participant number are still significant.

Centrality selection by multiplicity will bias the event sample: study of multiplicity fluctuations in a sample of events that are preselected by multiplicity. [1] A. Motornenko, K. Grebieszkow, E. Bratkovskaya, M. I. Gorenstein, M. Bleicher, and K. Werner, 1711.07789 [2] V.P. Konchakovski, M.I. Gorenstein, E.L. Bratkovskaya and W. Greiner, 1001.3085 10

# **Centrality selection: as it's done in the experiment**



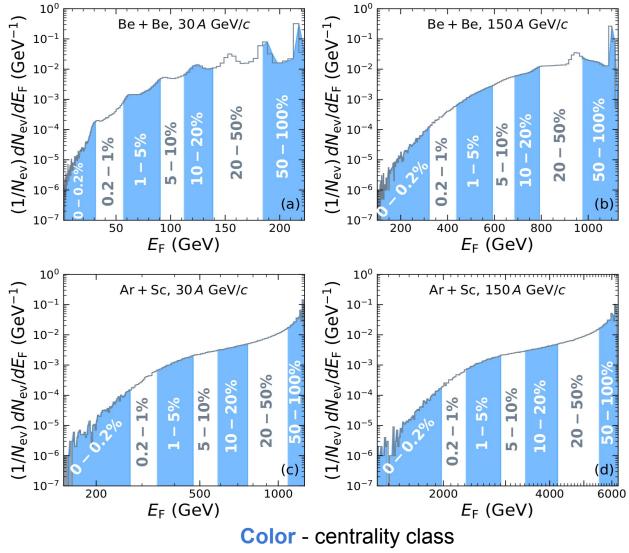
# **Energy in PSD: unbiased centrality for fluctuations**

Single peaks in spectra (for a single spectator) are smeared:

• energy of a spectator is **smeared by Fermi-motion** in nucleus (which is then boosted to the collision frame);

• PSD is a spectrometer: not able to distinguish spectator from forward particle – all particles in forward direction are counted

E unbiased centrality is fluctuations for measure analysis. Centrality by multiplicity bias puts on studied event sample artificial reduction of fluctuations.

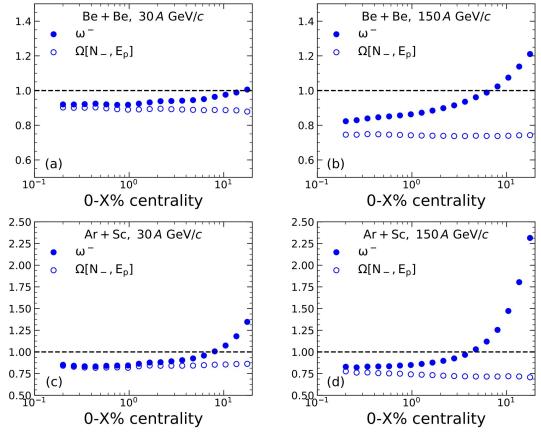


# **Omega in Be+Be and Ar+Sc**

To remove effects of volume fluctuations one can construct strongly intensive scaled variance  $\Omega[1,2]$ :

 $\Omega[N, E_P] = \omega[N] - (\langle N \cdot E_P \rangle - \langle N \rangle \cdot \langle E_P \rangle) / \langle E_P \rangle , E_P = E_{BEAM} - E_F$   $\Omega$  approaches scaled variance in limit of most central collision. if  $\Omega$  approaches  $\omega$  — collisions are violent enough.

At 1% centrality class all systems converge to a reasonable agreement between  $\Omega$  and  $\omega$ .



[1] M. Gazdzicki and S. Mrowczynski, Z. Phys. C 54, 127 (1992)
[2] M. I. Gorenstein and M. Gazdzicki, Phys. Rev. C 84, 014904 (2011)

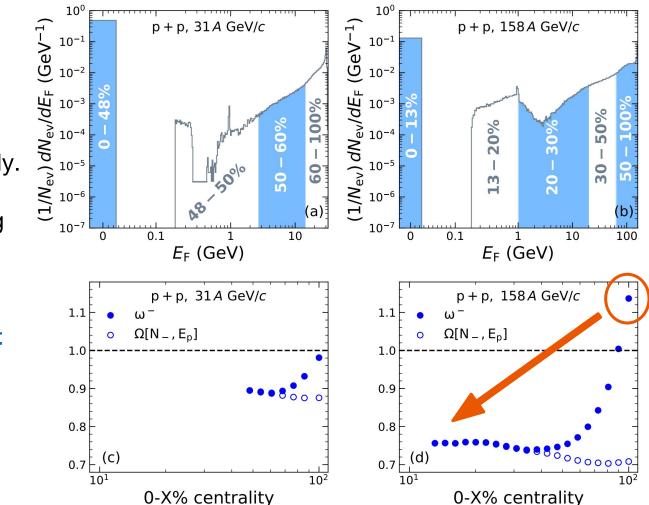
# **Centrality for p+p**

Impact parameter b is an ill defined quantity for p+p, but one may consider most "violent" p+p collisions defined by **small E**<sub>fwd</sub>:

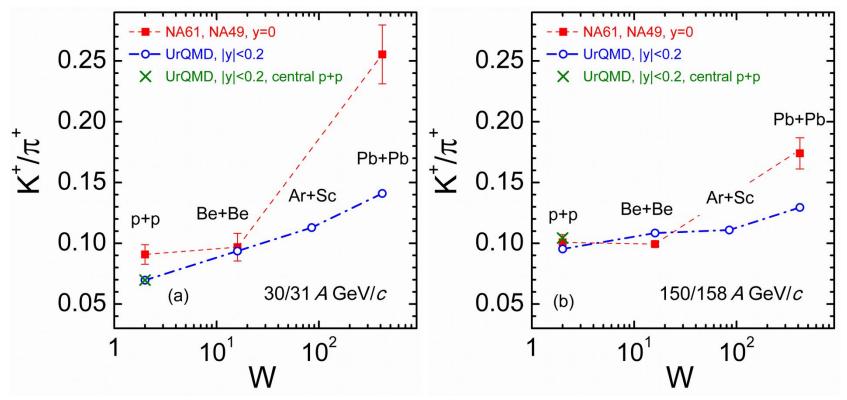
Peaks correspond to emission of single pion and nucleon respectively.

There's a non-vanishing probability for **E<sub>F</sub>=0** (no particles are emitted in the PSD). **We define those events as "most central"** 

For "most central" p+p one obtains  $\omega_{<1}$ .

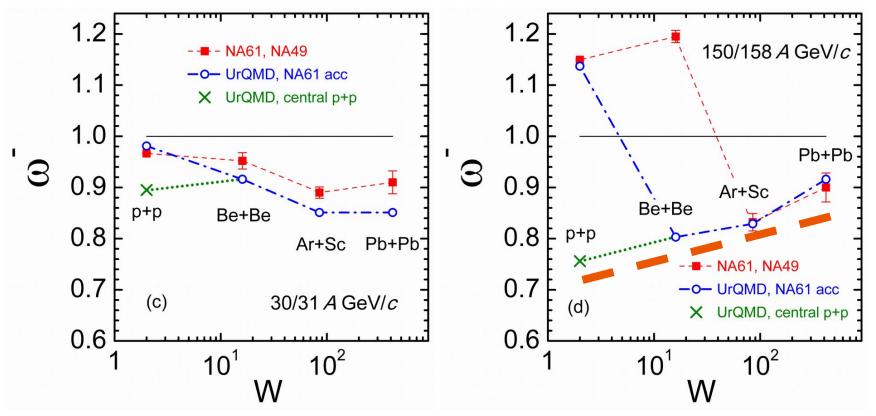


# K<sup>+</sup>/ $\pi$ <sup>+</sup> from UrQMD



- UrQMD underestimates strange production for large systems known issue;
- Small systems are Ok;
- Centrality in p+p doesn't affect ratio.

# **Fluctuations in UrQMD**



- Centrality in p+p reduces fluctuations, so  $\omega < 1$ ;
- UrQMD describes the data except Be+Be @ 150 GeV there are still exp. checks of this datapoint;
- Comparison of inelastic p+p with most central A+A apples with oranges;
- Centrality in p+p provides a monotonous increase of  $\omega$  @ 150 GeV/c.

# **Summary**

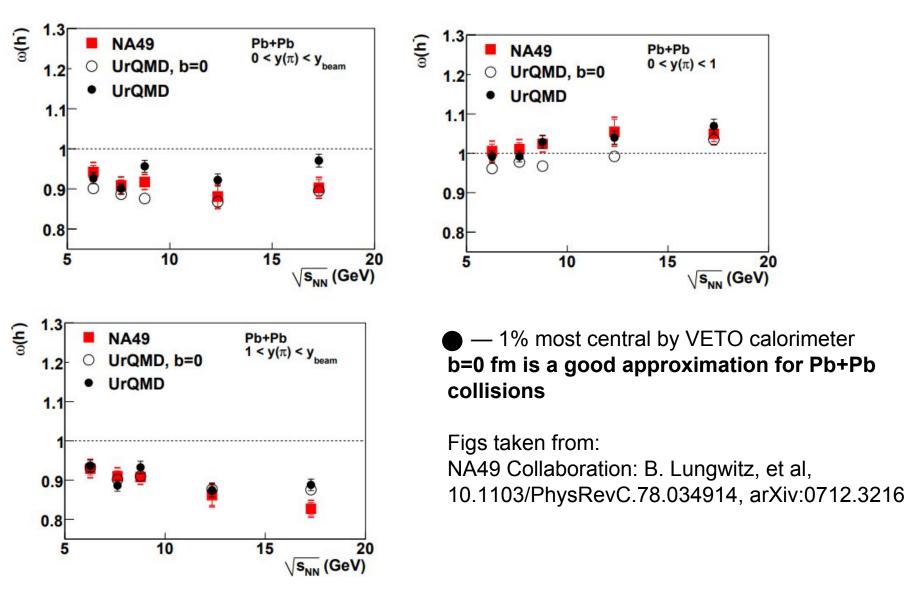
- Canonical effects are important in small systems;
- The CE HRG gives qualitative **agreement for K**<sup>+</sup>/ $\pi$ <sup>+</sup>, but cannot provide fluctuations  $\omega > 1$  for small systems;
- UrQMD describes scaled variance  $\omega$ , but:
  - Intermediate size systems require centrality to be done exactly as in the experiment;
- Centrality may be introduced for p+p collisions this may be a reason of discrepancy between the data and CE HRG;

■ centrality → measure of how violent collision is

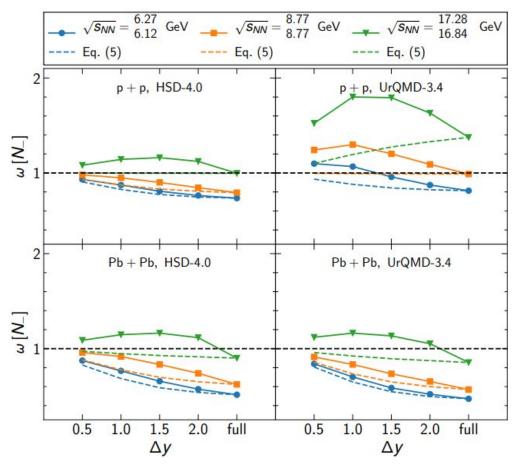
- Experimental procedures (centrality, acceptance) hide underlying physics — so comparison of the data and theory is not transparent;
- Description of  $K^+/\pi^+$  in large systems is still a problem in UrQMD

# Thanks for your attention!

# **Backup: Centrality selection for Pb+Pb**



# **Backup: Acceptance effects**



 $\omega$  [N<sub>ch</sub>] in p+p (top) and Pb+Pb (bottom) collisions calculated in different rapidity regions within HSD (left) and UrQMD (right) models.

A. Motornenko, K. Grebieszkow, E. Bratkovskaya, M. I. Gorenstein, M. Bleicher, and K. Werner, 1711.07789

Scaled variance  $\omega$  if acceptance (rapidity cut  $\Delta y$  ) is applied:

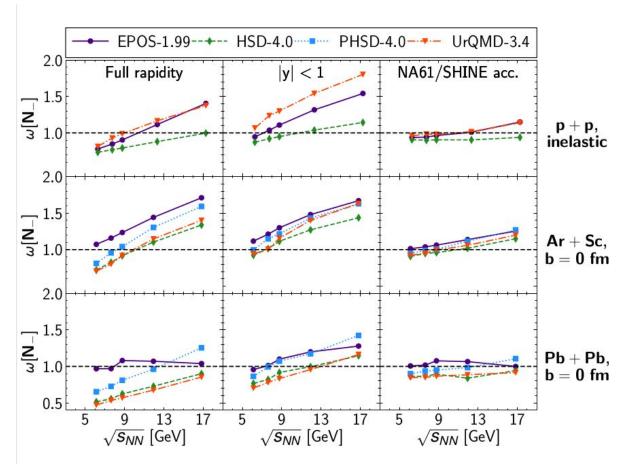
$$_{\mathsf{acc}}[X] = 1 - q + q \omega[X] \, ,$$
 $0 < q = rac{\langle X_{\mathsf{acc}} 
angle}{\langle X 
angle} < 1 \, .$ 

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assumes binomial statistics for particle to be accepted. Too naive.

 $\omega$  grows as function of  $\Delta y$  because conservation laws become suppressed, but then  $\omega \rightarrow 1$  as  $\Delta y \rightarrow 0$  to fulfill Poisson statistics.

# **Backup: b=0 collisions — to examine models**



$$\omega[N_{-}] = \frac{\langle N_{-}^2 \rangle - \langle N_{-} \rangle^2}{\langle N_{-} \rangle}$$

- No non-monotonous behavior;
- Violation of KNO scaling in EPOS  $Pb+Pb (\omega \not\sim \langle N \rangle)$
- significant suppression of ω by NA61/SHINE acceptance.

# b=0 events doesn't show any non-monotonous behavior so no critical phenomena in the models

A. Motornenko, K. Grebieszkow, E. Bratkovskaya, M. I. Gorenstein, M Bleicher, and K. Werner, 1711.07789