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The final state swing

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Outline

- ▶ Parton showers, “Pre-confinement” and the size of N_C
- ▶ Colour reconnections
- ▶ The dipole swing in the shower
- ▶ Outlook (pA & AA)

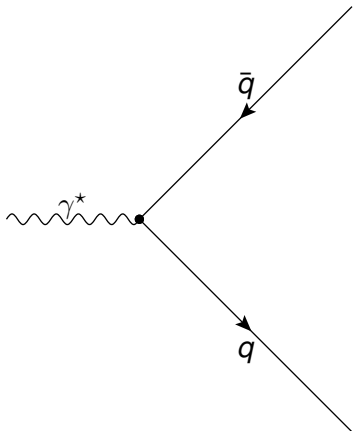


The importance of colour connections

- ▶ All hadrons are colour singlets.
- ▶ Any realistic hadronisation model must ensure this.
- ▶ Exact treatment of colour structures in LHC events is impossible(?)
- ▶ All parton shower approaches use the $N_C \rightarrow \infty$ approximation which gives a unique colour structure.



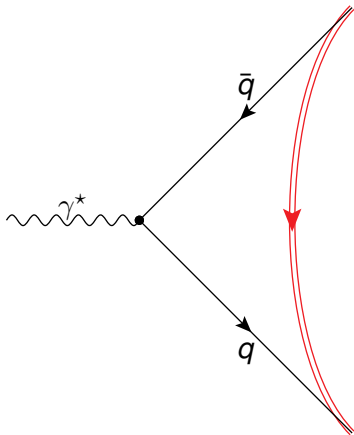
Parton and dipole showers



- ▶ Parton splitting
- ▶ Dipole splitting
- ▶ Pre-confinement: partons close in phase space are likely to be colour-connected.
Nature likes short strings.
- ▶ $N_C \rightarrow \infty$ gives a unique colour flow.
- ▶ But $N_C = 3$.



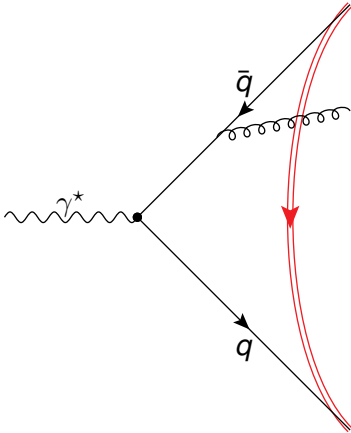
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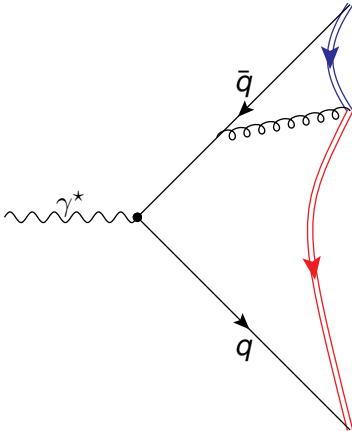
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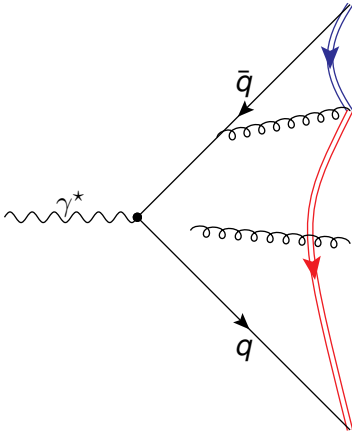
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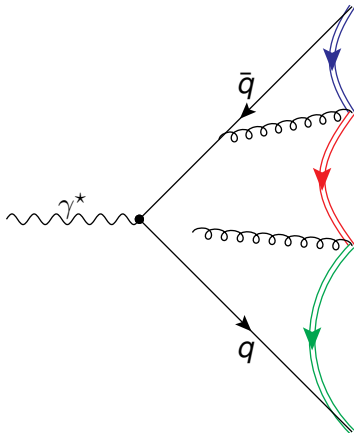
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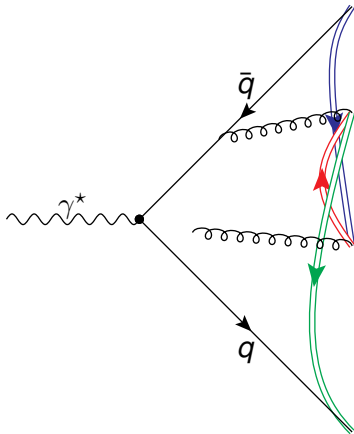
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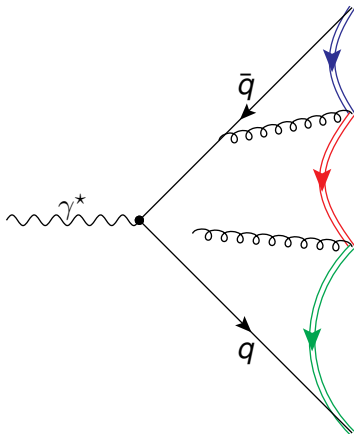
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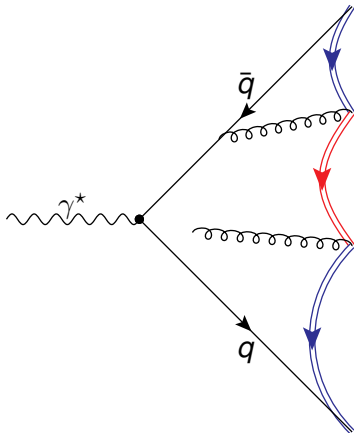
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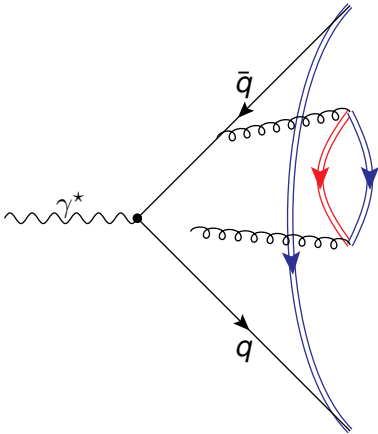
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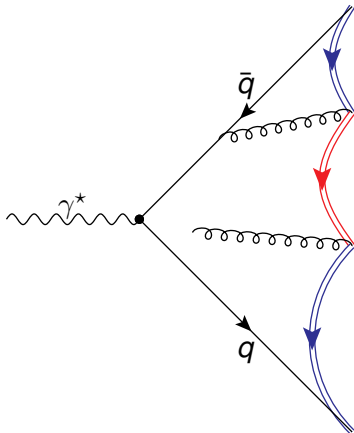
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Colour reconnections

Colour reconnections is a way to include effects of $N_C < \infty$.
The guiding principles are:

- ▶ Probability to reconnect $\sim 1/N_C^2$
- ▶ Nature likes short strings
- ▶ There are no colour-singlet gluons.

[Sjöstrand, Khoze, Gustafson, Zerwas, Lönnblad, Edin, Ingelman, Rathsman, Gieseke, Kirchgaßer, ...]



Short strings?

We typically measure the string lengths in terms of the λ -measure

For a string consisting of n dipoles between a quark and an anti-quark connected with $n - 1$ gluons:

$(q_0 - g_1 - g_2 - \dots - g_{n-1} - \bar{q}_n)$

$$\lambda = \sum_{i=0}^{n-1} \log \left(1 + \frac{m_{i,i+1}^2}{m_0^2} \right)$$



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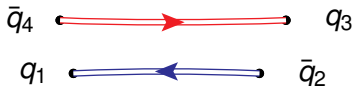
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Simple reconnections

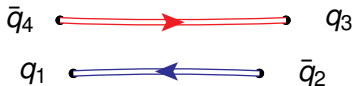


Reconnect?

- ▶ with probability $1/N_C^2$
- ▶ only if $m_{14}m_{23} < m_{12}m_{34}$



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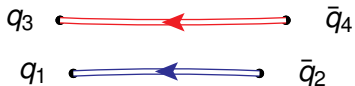


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Junctions

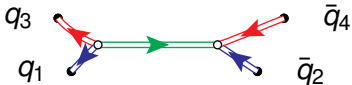


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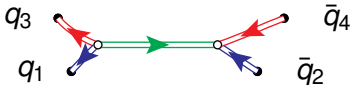


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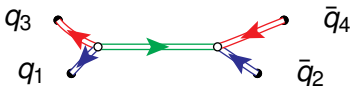


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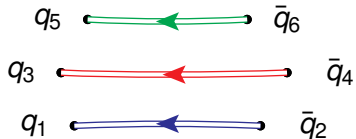


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More Junctions

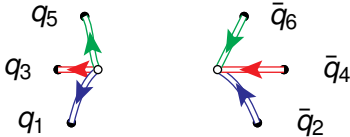


Reconnect?

- ▶ with probability $1/N_C^3$
- ▶ only if λ -measure is reduced
- ▶ (accessible with two subsequent reconnections)



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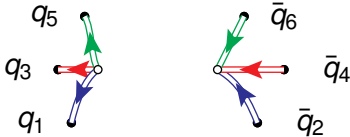


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Perturbative effects

We expect effects of $N_C = 3 < \infty$ also on the perturbative level.

We want a full-colour parton shower, but this probably requires an amplitude-level parton shower scheme, which can become very messy.

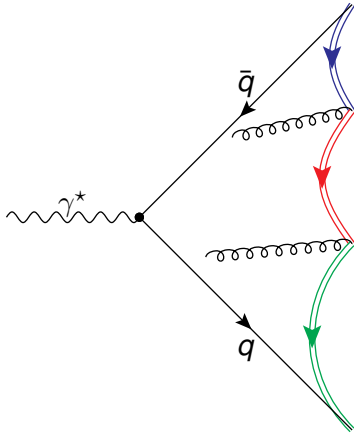
Instead modify what we have: the dipole shower.

Amend it with dipole reconnections between each emission.

Let's put some **swing** into the the dipole shower!



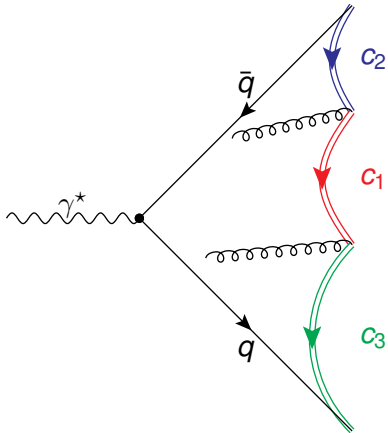
The Dipole Swing



- ▶ Assign a colour index (1-9) to each dipole
- ▶ Dipoles connected with a gluon must have $c_i \neq c_j$
- ▶ New colour index between the emitted gluon and the emitter
- ▶ Only dipoles with the same index may swing
- ▶ Let's Swing



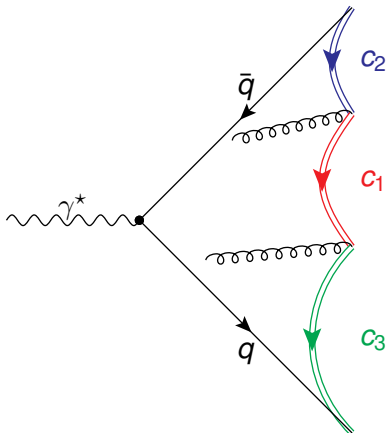
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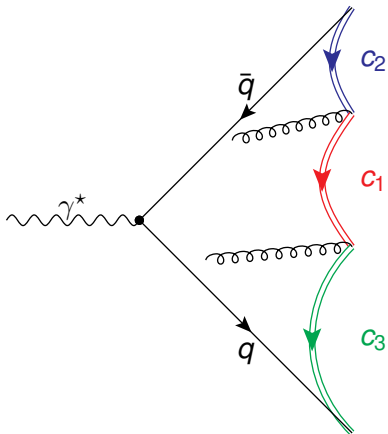
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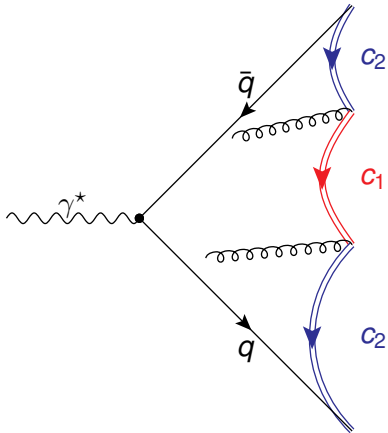
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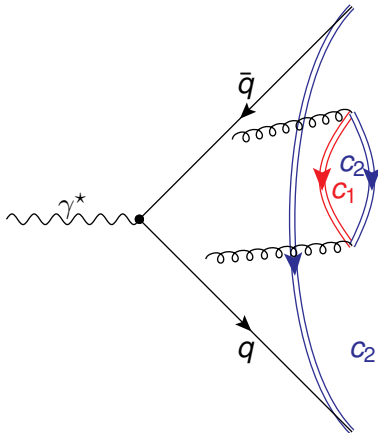
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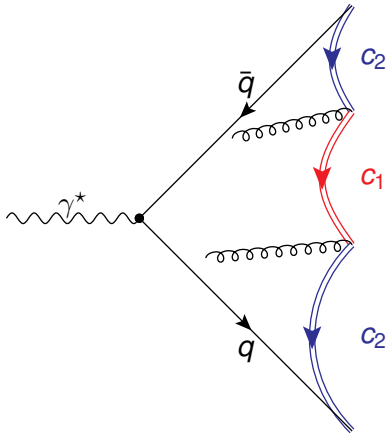
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- ▶ Only dipoles with the same index may swing
- ▶ Let's Swing both ways



The dipole emissions are limited by the dipole mass (cf. angular ordering)

The dipole shower is ordered in transverse momentum, k_{\perp}

The distribution of the *next* emission is given by

$$\frac{d\mathcal{P}}{dk_{\perp}^2} = \frac{\alpha_S}{k_{\perp}^2} \sum_i \int dz P_i(z) \times \Delta(k_{\perp\max}^2, k_{\perp}^2)$$

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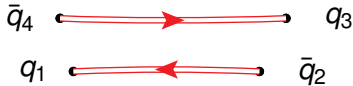
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Add the probability that a dipole may swing

$$\frac{d\mathcal{P}_{\text{swing}}}{dk_{\perp}^2} = \lambda \frac{m_{12}^2 m_{34}^2}{m_{14}^2 m_{32}^2} \times \Delta_{\text{swing}}(k_{\perp\max}^2, k_{\perp}^2)$$

where λ is a strength parameter





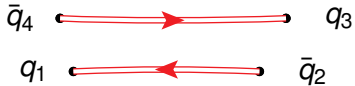
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- ▶ The weighted average of the radiation from the two dipole pair configuration emulates quadrupole radiation.
- ▶ Prefers small mass dipoles giving less radiation





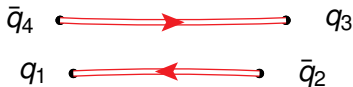
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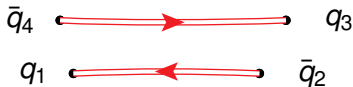
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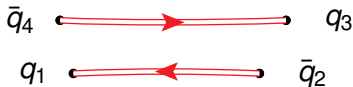
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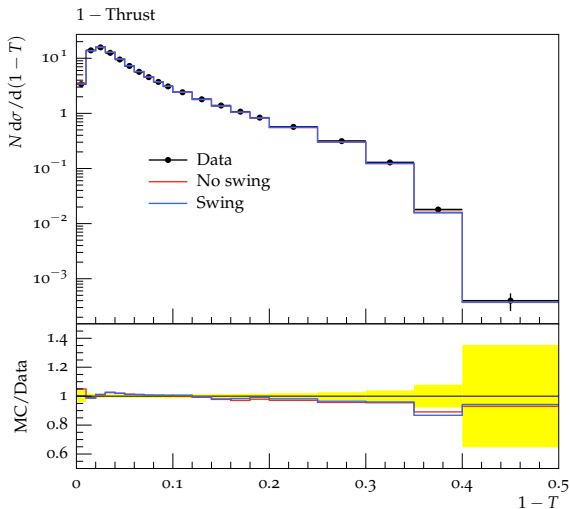




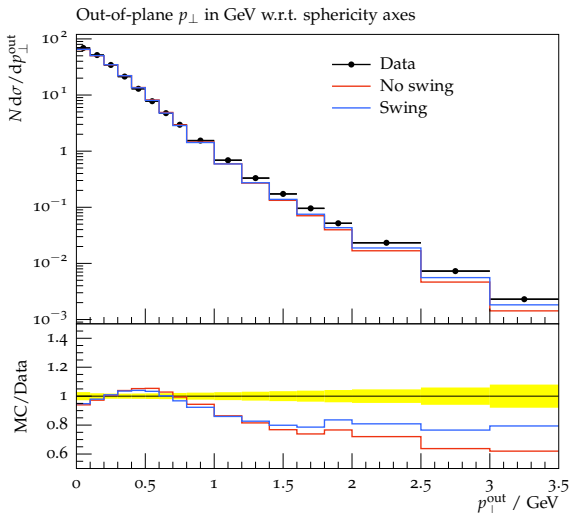
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Small effects in e^+e^- (after retuning)



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Outlook

- ▶ Implemented in Ariadne (DIPSY)
- ▶ Will be implemented in Pythia8 (Angantyr)
- ▶ Need to include a space-time picture in pA & AA
- ▶ Will affect flow and jet shapes



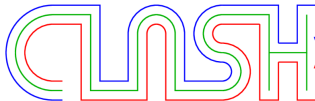
Thanks!



Vetenskapsrådet



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Colour Reconnections

- ▶ Sjöstrand et al., Phys.Rev. D36 (1987) 2019
- ▶ Gustafson et al., Z.Phys. C64 (1994) 659-664
- ▶ Sjöstrand et al., Phys.Rev.Lett. 72 (1994) 28-31
- ▶ Edin et al., Phys.Lett. B366 (1996) 371-378
- ▶ Lönnblad, Z.Phys. C70 (1996) 107-114
- ▶ Gieseke et al., JHEP 1811 (2018) 149

