

# Can we calculate and measure jet quenching in small systems?

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## Qualitative considerations: theory

- 1 Observing jet quenching (JQ) establishes existence of final state interactions (FSI).
- 2 Favored (hydro/ transport) explanations of collectivity ( $v_n$ -flow) are based on FSI.
- 3 Small systems (pp, pA, periph. AA) show soft collectivity but (so far) no unambiguous JQ.

There is logical tension between these statements:  
**Either:** No JQ in small systems. Then interpretation of their soft collectivity needs to involve no-FSI mechanisms.

**Or:** Improved experiments identify JQ in small systems. Then improved theory needs to relate this small JQ to soft collectivity.

**The search for JQ in small systems can inform the quest for a unified dynamical description of hard and soft medium-effects.**

## What theory says? (imho)

Qualitatively, JQ and soft collectivity are related:

- 1 In weakly coupled systems:
  - hydrodynamization via bottom-up thermalization, governed by same  $2 \rightarrow 2$  and  $1 \rightarrow 2$  (LPM) collision kernels as JQ, efficient on short timescales [?].
- 2 In strongly coupled systems:
  - (almost) perfect fluidity and JQ-drag both unavoidable.

Quantitative statements uncertain or missing:

- 1 Many JQ codes not (yet?) explored for pp/pA
  - Does embedding jets in spatio-temporal evolution of pp or pA have higher uncertainties?
  - Do JQ models rely on resumming geometrically enhanced higher-twist effects ( $A^{1/3}/Q^2$ )? If so, are they applicable to small systems?
  - In small systems, triggers on event activity like  $dN^{\text{ch}}/d\eta$  influence selection of hard probes. Many issues: What are suitable JQ observables? How to avoid / utilize trigger biases?  $\Rightarrow$  see experimental part
- 2 BDMPS-type JQ sensitive to  $\hat{q} = \frac{\langle k_{\perp}^2 \rangle_{\text{med}}}{\lambda_{\text{mfp}}}$ 
  - insensitive to whether  $\hat{q}$  built up by many soft or very few harder interactions  $N_{\text{int}}$ .
  - but  $\langle N_{\text{int}} \rangle$  determines whether all “jets” suffer small or whether few jets suffer sizeable medium modifications  $\Rightarrow$  information beyond fluid dynamic averages ( $T^{\mu\nu}$ ) likely relevant for quantitative description of JQ.

## The medium in small systems

- Both almost perfect fluid models and transport with few scatterings [?] reproduce collectivity in small systems [?].
- Close to equilibrium, all QFTs carry hydro- and non-hydro degrees of freedom.
- Transport and Israel-Stewart theory can carry same hydro but different non-hydro excitations [?].

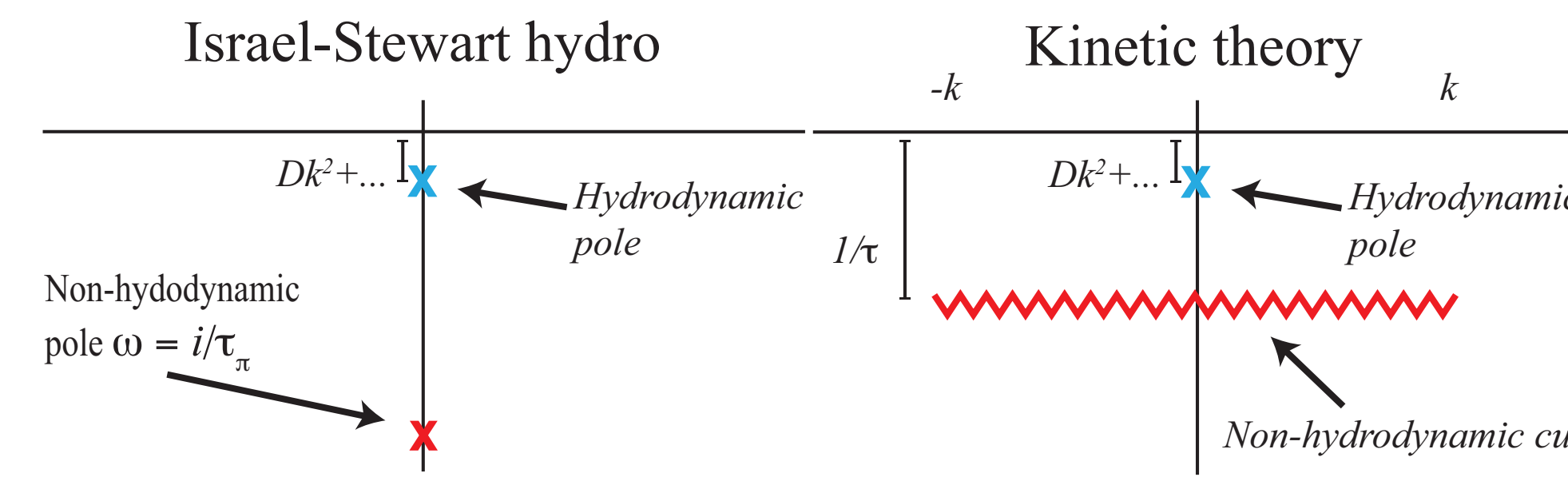


Figure 1: Pole structure of retarded correlator  $G_R(\omega, k)$  (shear mode). **Non-hydro modes become more important in smaller systems [?].**

- In conformal systems without other intrinsic scale, opacity  $\hat{\gamma} = \gamma R^{3/4} (\varepsilon_0 \tau_0)^{1/4}$ , ( $\gamma = \frac{\text{const.}}{\eta/s}$ ), is only scaling variable [?],

$$\hat{\gamma} = \frac{0.11}{\eta/s} \left( \frac{1}{\pi f_{\text{work}}(\hat{\gamma})} \right)^{1/4} \left( R \frac{dE_{\perp}}{d\eta_s} \right)^{1/4}$$

- $\hat{\gamma} \gg 1 \Rightarrow \hat{\gamma} = c_{\text{hyd}} \left( \frac{dN}{d\eta} \right)^{1/3}$  multiplicity scaling.
- $\hat{\gamma} \lesssim 1 \Rightarrow \hat{\gamma} = c_{\text{hyd}} \left( R \langle p_{\perp} \rangle \frac{dN}{d\eta} \right)^{1/4}$  no multiplicity scaling.

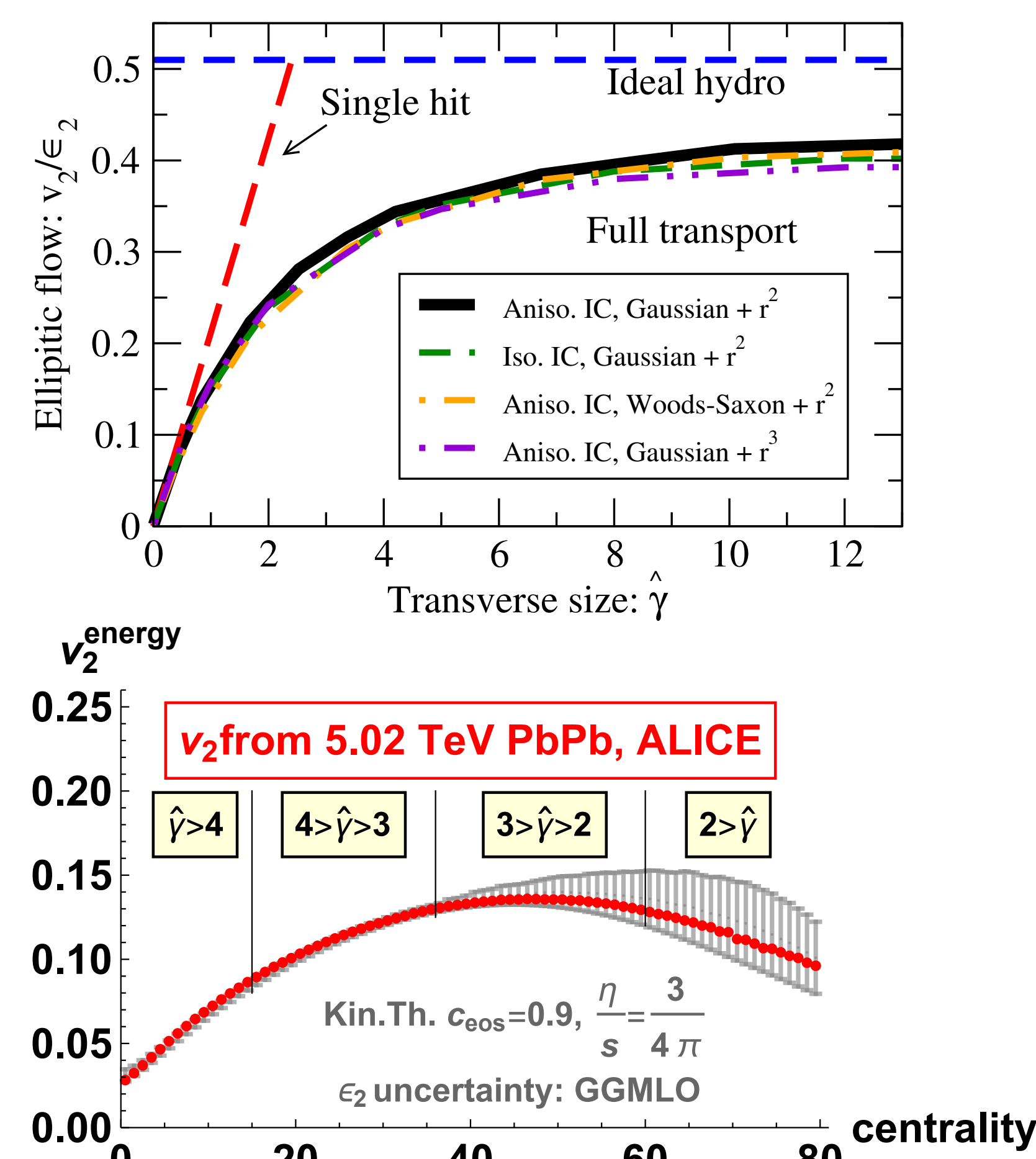


Figure 2: Small system soft collectivity may originate not at fluid dynamic limit. **Kinetic transport: enhanced escape probability of particles/jets without scattering [?].** Figs. from [?].

## Qualitative considerations: experiment

- 1 Jet quenching effects in small systems expected to be small; little guidance from theory
- 2 characterization of collision geometry in small systems is challenging: large relative fluctuations of Event Activity  $\rightarrow$  large model dependence of  $\langle T_{\text{pA}} \rangle$

Choice of observables: jets, hadrons

- inclusive suppression
- triggered coincidence: recoil suppression and deflection
- jet substructure

$\rightarrow$  require precise control over backgrounds, including MPIs

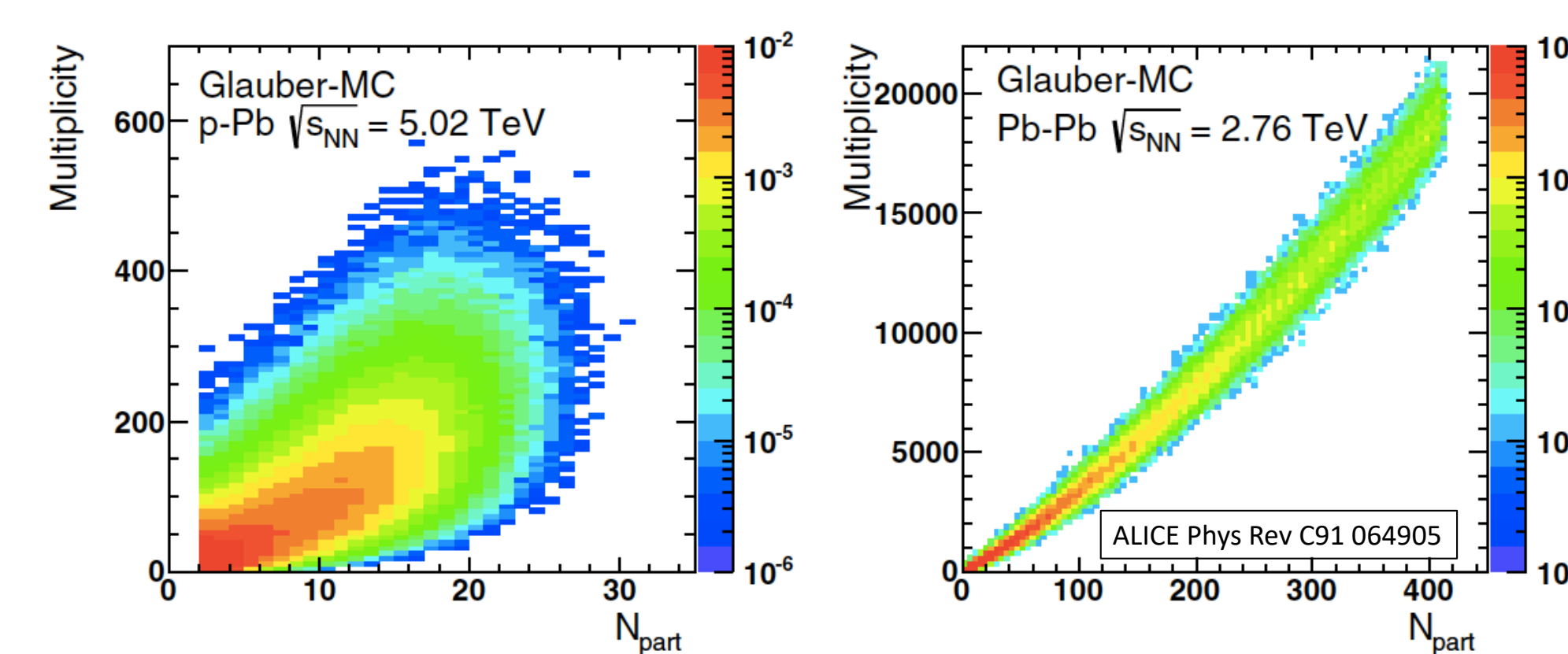
$\rightarrow$  independence of  $\langle T_{\text{pA}} \rangle$ : coincidence, substructure observables

$\rightarrow$  kinematic range: low  $p_{\text{T}}^{\text{jet}}$  may have largest relative modification/largest sensitivity

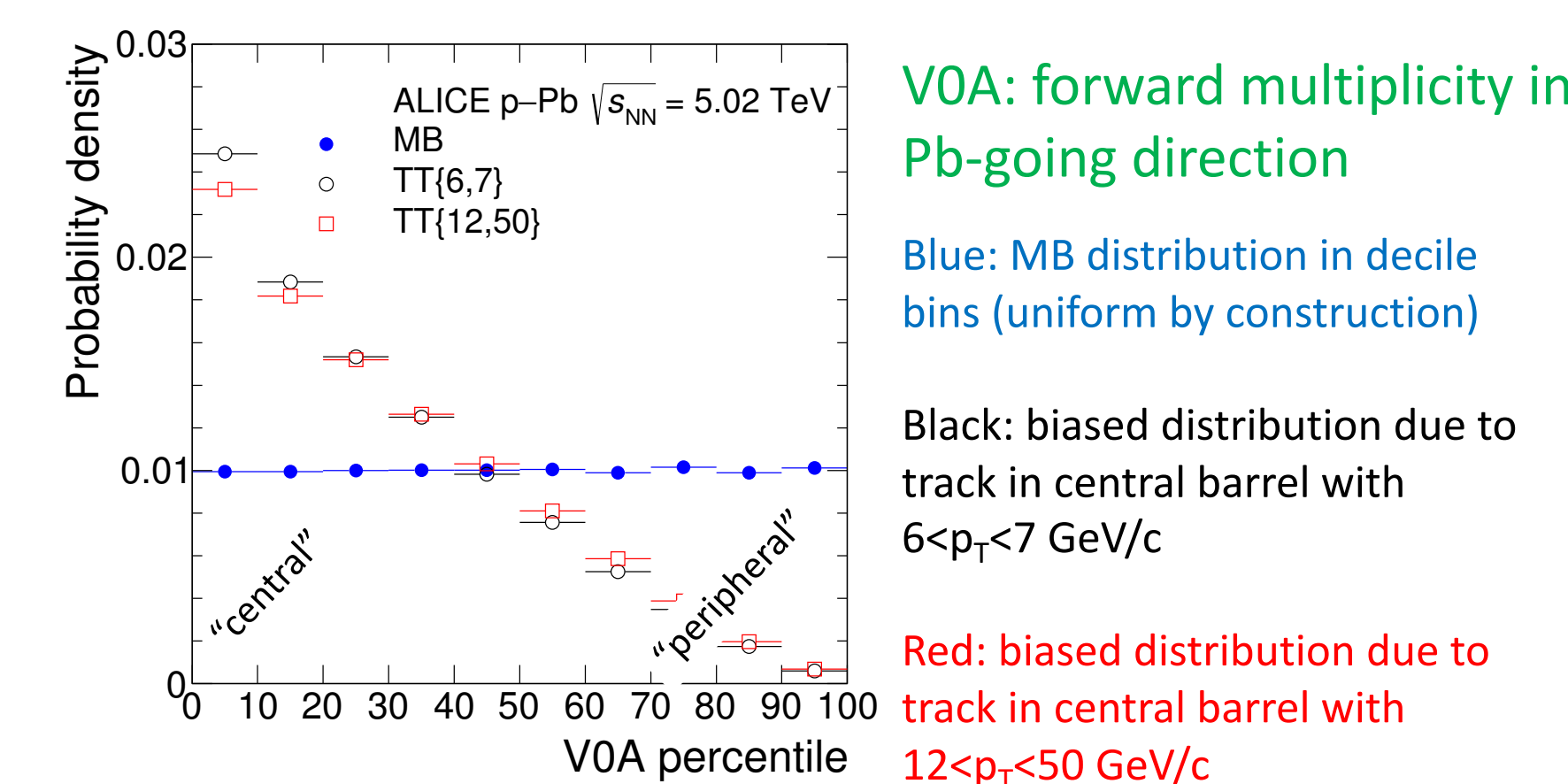
**Optimum strategy: utilize tension between multiple observables with different parametric dependencies on e.g.  $\hat{q}$  and  $L$**

## EA and “centrality”

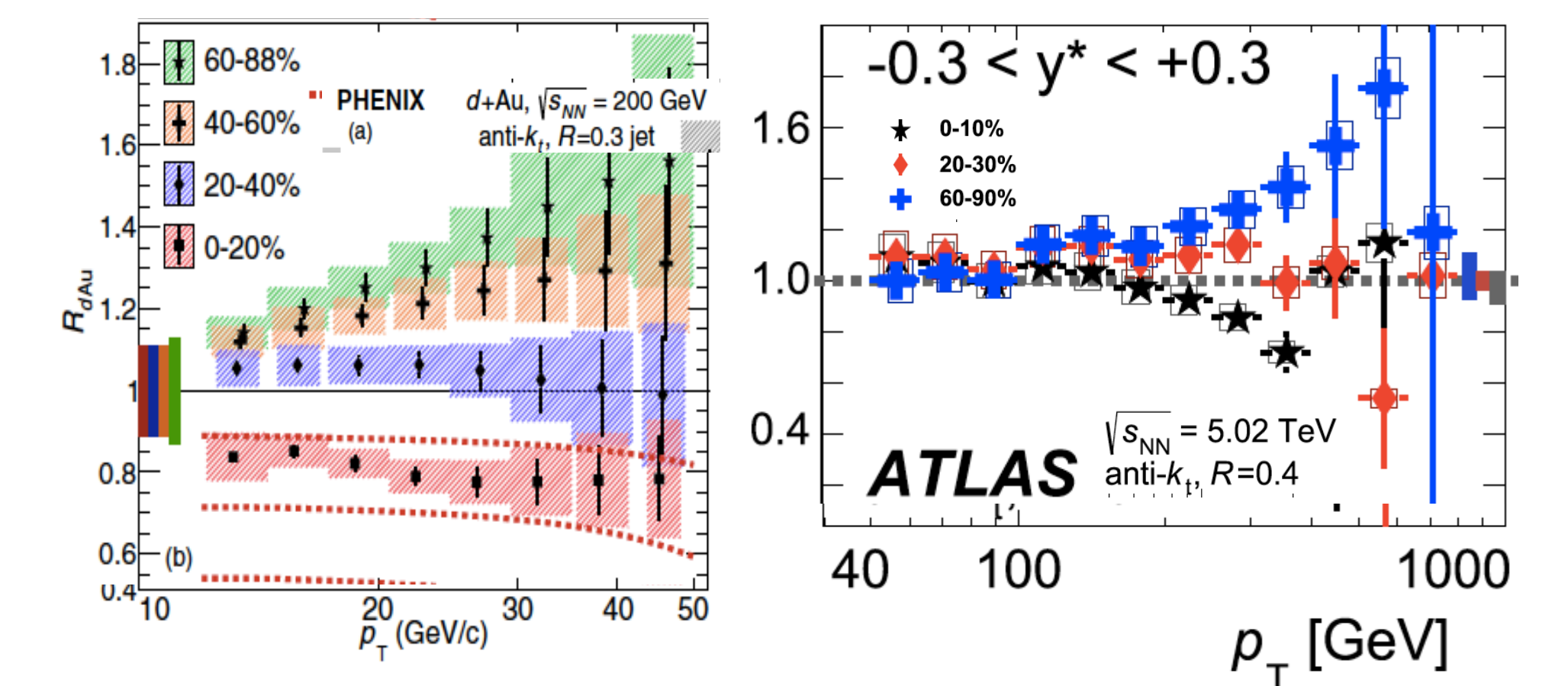
Glauber calculations showing correlation of forward multiplicity with calculated  $N_{\text{part}}$  for p+Pb and Pb+Pb; p+Pb has much larger relative fluctuations.



EA bias due to the presence of a hard process:



## Inclusive $R_{\text{pA}}$ and semi-inclusive jet measurements



### Semi-inclusive jet quenching observables

- Trigger-normalized yield suppression  $\text{BDMPS} \sim \langle \hat{q} L^2 \rangle$

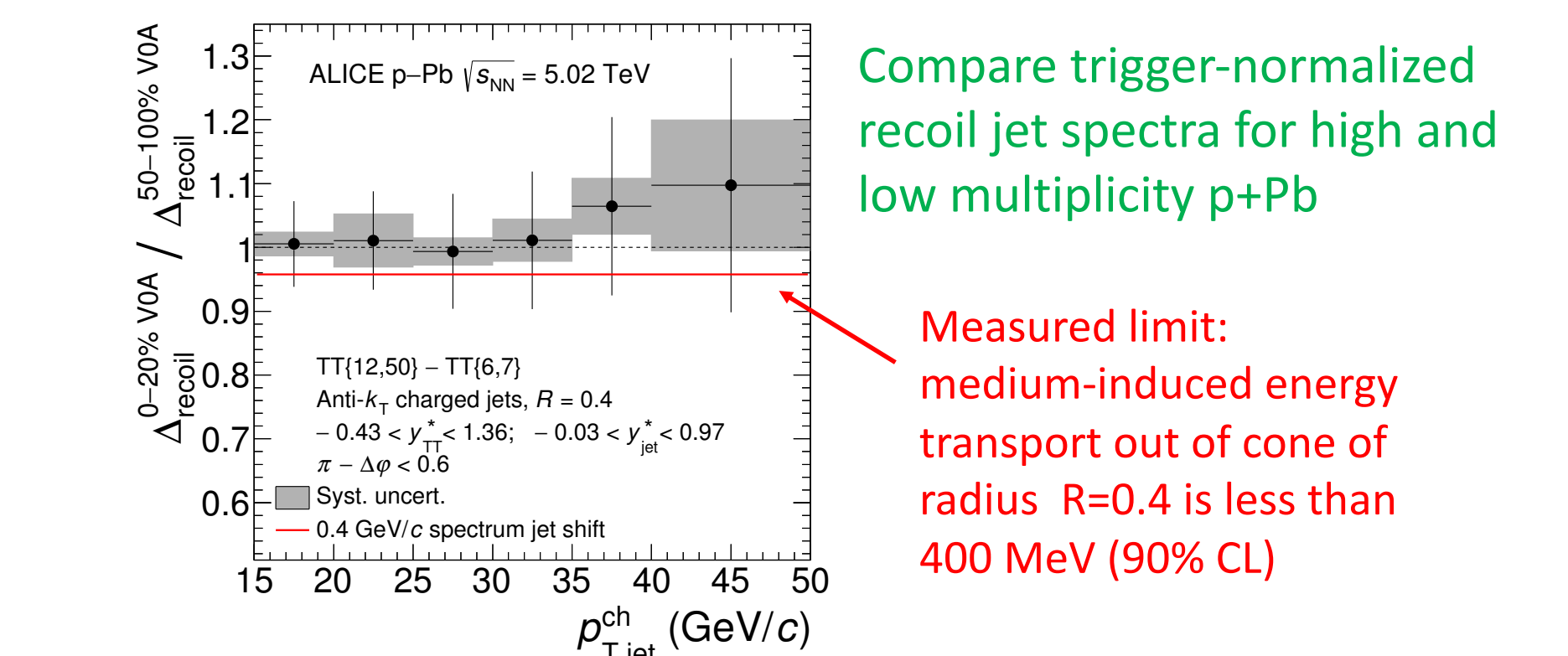
- Azimuthal deflection  $\sim \langle \hat{q} L \rangle$

No requirement to associate Event Activity and geometry:

$$\frac{1}{N_{\text{trig}}^{\text{AA}}} \frac{d^2 N_{\text{jet}}^{\text{AA}}}{d p_{\text{T,jet}}^{\text{ch}} d \eta_{\text{jet}}} \Big|_{p_{\text{T,trig}} \in \text{TT}} = \left( \frac{1}{\sigma^{\text{AA} \rightarrow \text{h} + \text{X}}} \cdot \frac{d^2 \sigma^{\text{AA} \rightarrow \text{h} + \text{jet} + \text{X}}}{d p_{\text{T,jet}}^{\text{ch}} d \eta_{\text{jet}}} \right) \Big|_{p_{\text{T,h}} \in \text{TT}}$$

In case of no nuclear effects

$$\frac{1}{N_{\text{trig}}^{\text{AA}}} \frac{d^2 N_{\text{jet}}^{\text{AA}}}{d p_{\text{T,jet}}^{\text{ch}} d \eta_{\text{jet}}} \Big|_{p_{\text{T,trig}} \in \text{TT}} = \left( \frac{1}{\sigma^{\text{pp} \rightarrow \text{h} + \text{X}}} \cdot \frac{d^2 \sigma^{\text{pp} \rightarrow \text{h} + \text{jet} + \text{X}}}{d p_{\text{T,jet}}^{\text{ch}} d \eta_{\text{jet}}} \right) \Big|_{p_{\text{T,h}} \in \text{TT}} \times \frac{T_{\text{AA}}}{T_{\text{AA}}}$$



Inclusive  $R_{\text{pA}}$  and semi-inclusive limit are not compatible [?]. My guess: uncorrected biases in  $\langle T_{\text{pA}} \rangle$  due to QCD correlations, beyond the increase in pp UE[?].

## References

- [1] A. Kurkela, E. Lu, Phys. Rev. Lett. **113** (2014)
- [2] L. He, T. Edmonds, Z. W. Lin, F. Liu, D. Molnar and F. Wang, Phys. Lett. B **753** (2016) 506
- [3] A. Kurkela, U. A. Wiedemann and B. Wu, arXiv:1905.05139 [hep-ph].
- [4] S. Acharya *et al.* [ALICE Collaboration], Phys. Lett. B **783**, 95 (2018)
- [5] A. Adare *et al.* [PHENIX Collaboration], Phys. Rev. C **90**, no. 3, 034902 (2014)