

Novelties of Lund strings in Heavy ion Collisions

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With Christian Bierlich [†], Gösta Gustafson, Leif Lönnblad

PhD talks

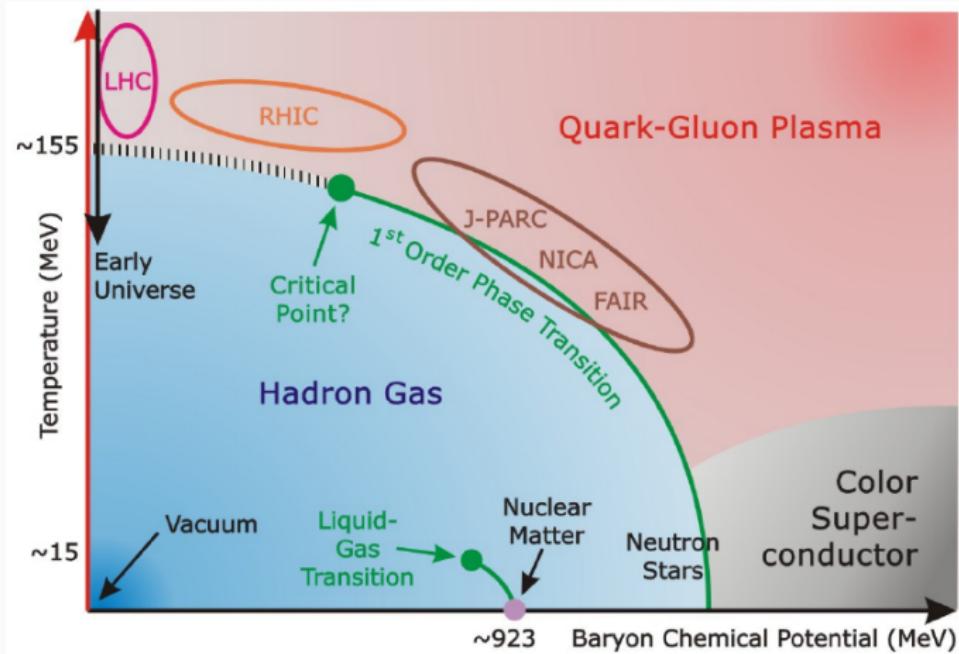
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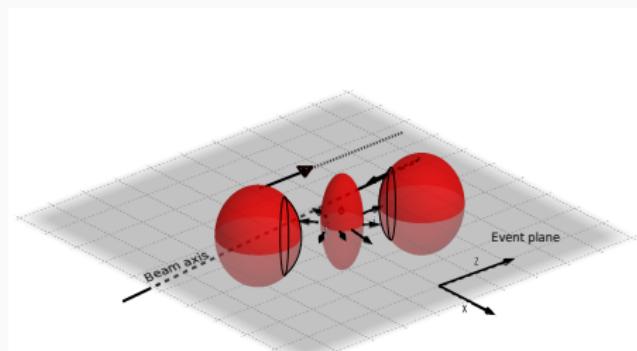


Why heavy ions and what are we looking for?



Conjectured QCD phase diagram

Anisotropy in collision geometry



Heavy ion collision in the lab frame

1. Initial eccentricities



2. Anisotropic flow
coefficients



3. Transport properties of the
produced medium

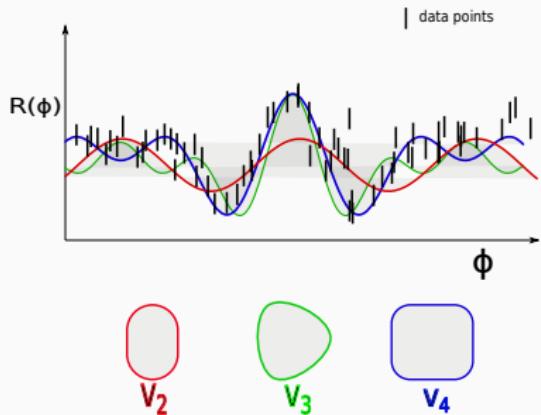
Fourier coefficients in particle distributions

$$\frac{dN}{d\phi} \sim R(\phi) \propto 1 + 2\sum_{n=1}^{\infty} v_n \cos(n\phi)$$

$$R(\phi) \propto 1 + 2v_2 \cos(2\phi)$$

$$R(\phi) \propto 1 + 2v_2 \cos(2\phi) + v_3 \cos(3\phi)$$

$$R(\phi) \propto 1 + 2v_2 \cos(2\phi) + v_3 \cos(3\phi) + v_4 \cos(4\phi)$$



Angular particle distribution and
geometrical interpretation of Fourier
coefficients

Fourier coefficients in particle distributions

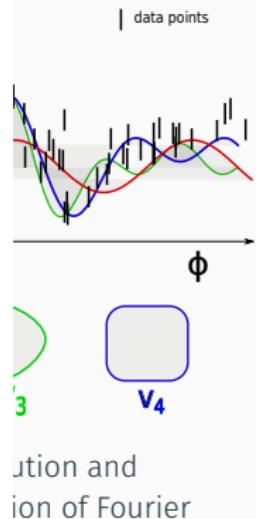
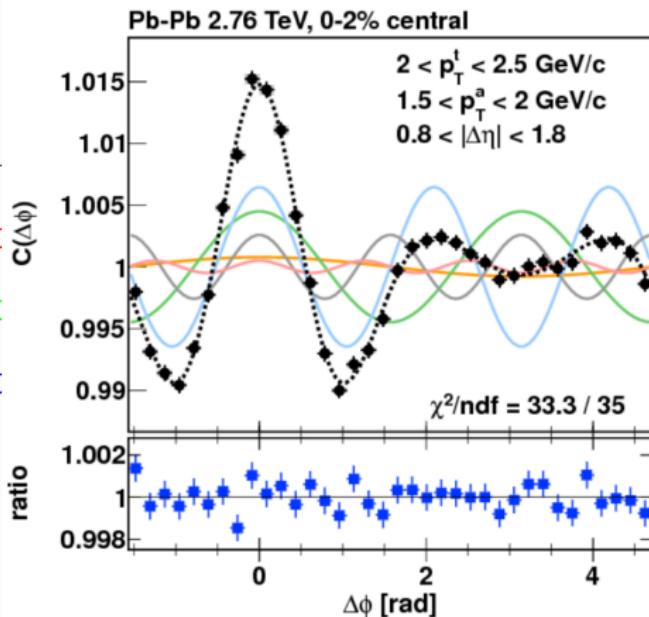
$$\frac{dN}{d\phi} \sim R(\phi) \propto 1 +$$

$$R(\phi) \propto 1 + 2v_2cc$$

$$R(\phi) \propto 1 + 2v_2cc$$

$$R(\phi) \propto 1 + 2v_2cc$$

$$+ v_4 \cos(4\phi)$$



The quest for the perfect fluid: Quark Gluon Plasma

To search for explanation of possible collective effects (fluid-like behaviour) in high multiplicity p-p collisions with string model

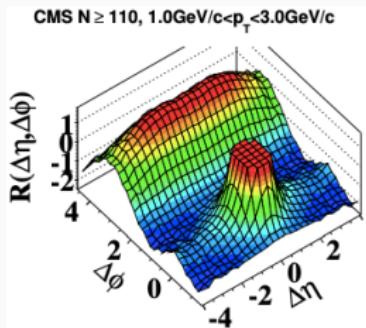


Figure 3: p-p

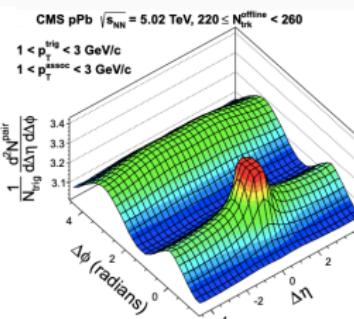


Figure 4: p-Pb

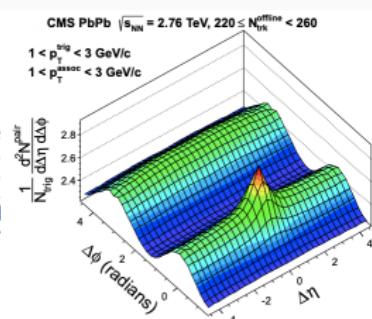
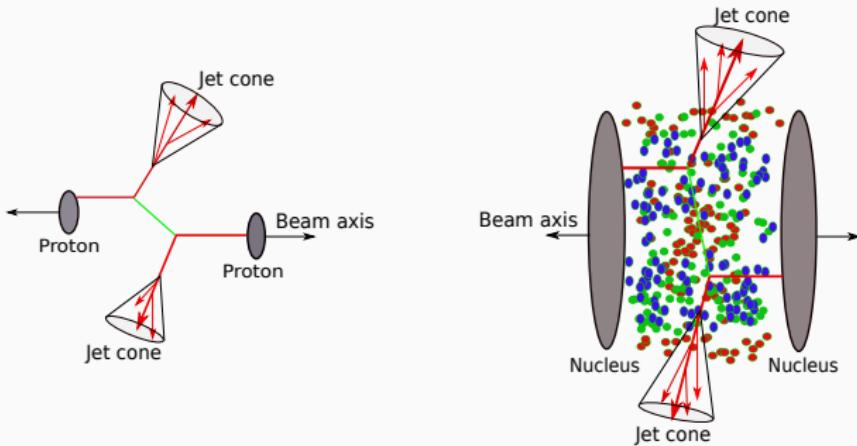


Figure 5: Pb-Pb

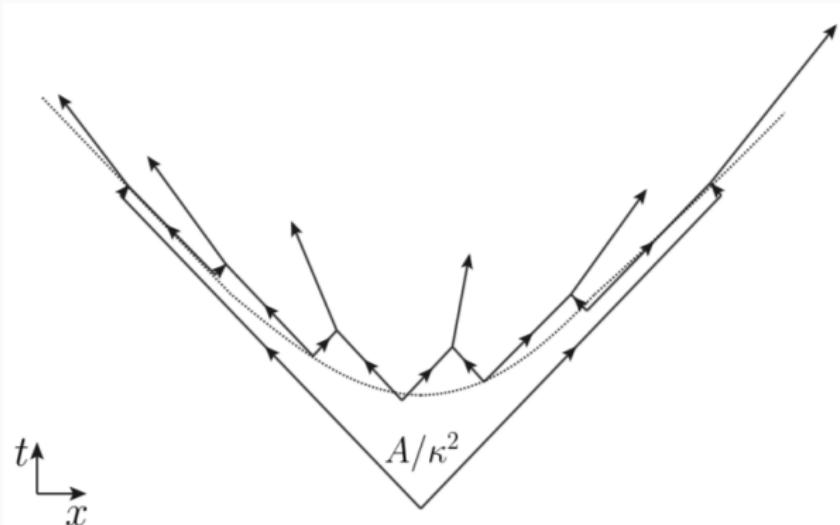
Messengers from the Deep: Jets



Jets in p-p and in A-A collisions

- Change in jet energies (quenching) observed in p-A and A-A collisions
- Is there any jet quenching in high multiplicity p-p events?
Explanation in string model?

Lund string model



x - t diagram of a Lund string, with fragmentation hyperbola. The area enclosed by the quark lines is the coherence area A in units of the string tension κ

Shoving in Angantyr and advancements

1. Aspects of Angantyr:

- A-A is treated as a collection of overlaid p-p collisions
- Modifications needed when one nucleon in one nucleus collides with several nucleons in the other
- No collective effects

Shoving in Angantyr and advancements

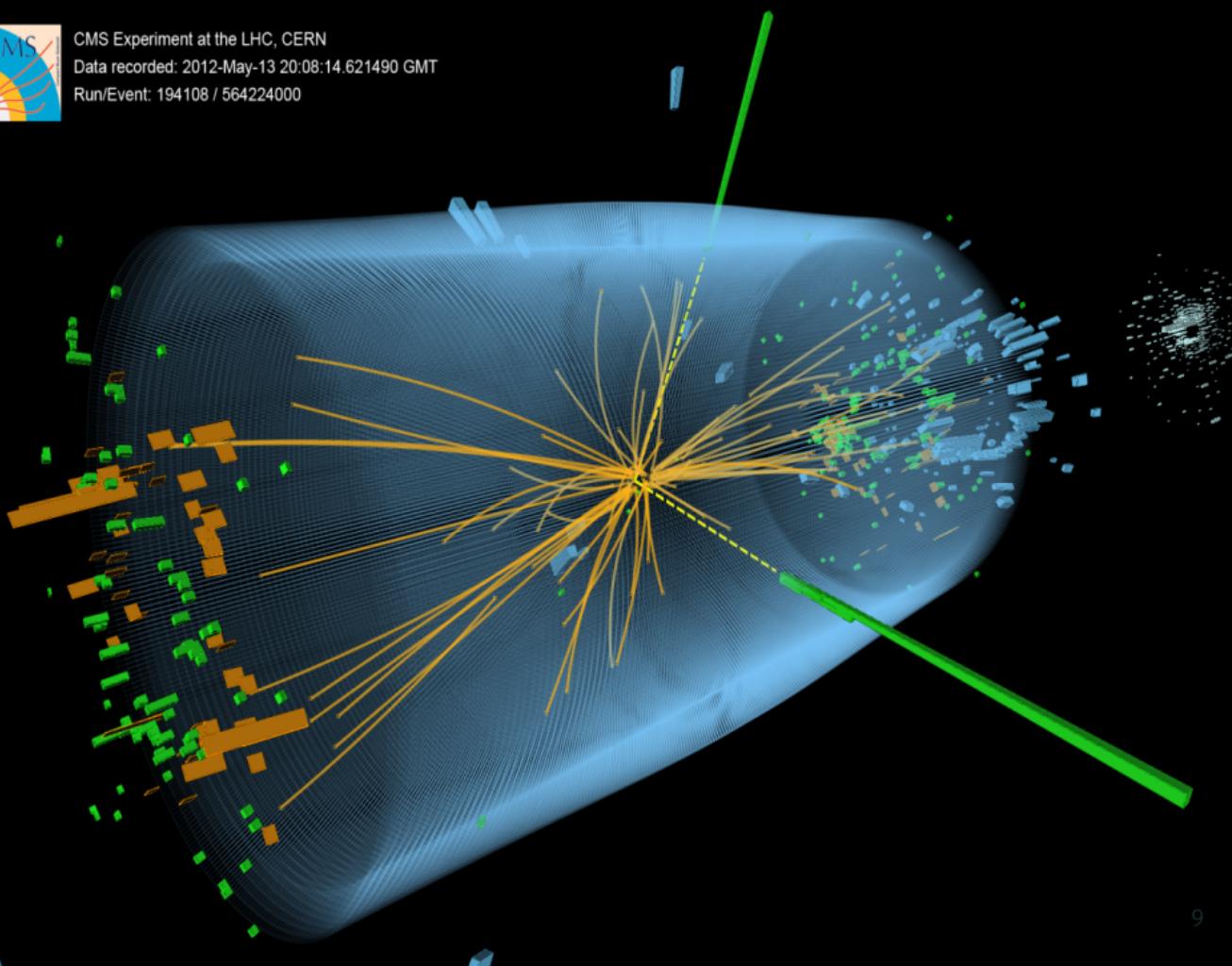
1. Aspects of Angantyr:
 - A-A is treated as a collection of overlaid pp collisions
 - Modifications needed when one nucleon in one nucleus collides with several nucleons in the other
 - No collective effects
2. Mechanisms to study high-multiplicity p-p and A-A behaviours :
 - **String shoving** → Final-state collective effects?

Methodology

1. Parallel frame
2. Interaction force between two strings
3. Push distribution among hadrons

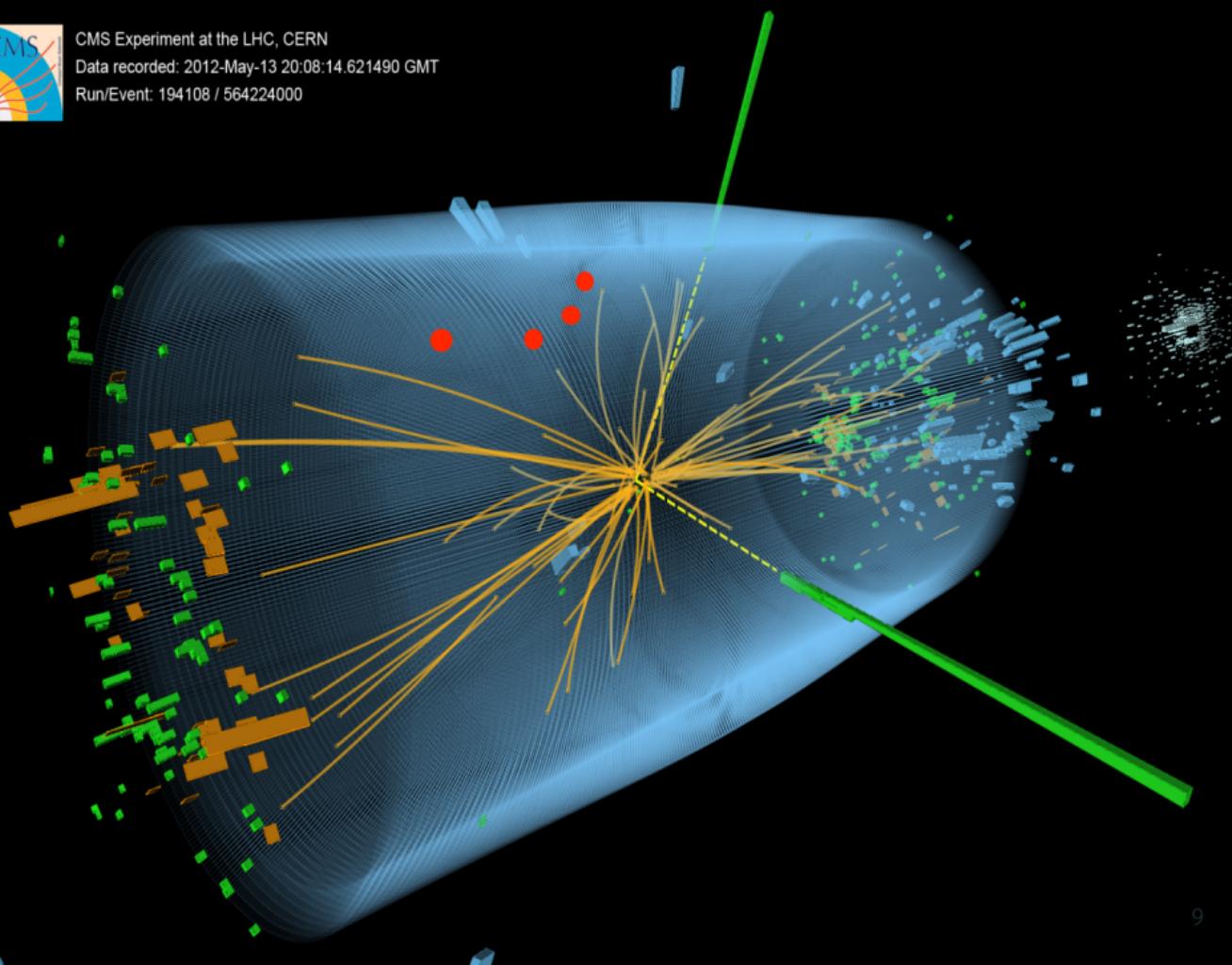


CMS Experiment at the LHC, CERN
Data recorded: 2012-May-13 20:08:14.621490 GMT
Run/Event: 194108 / 564224000

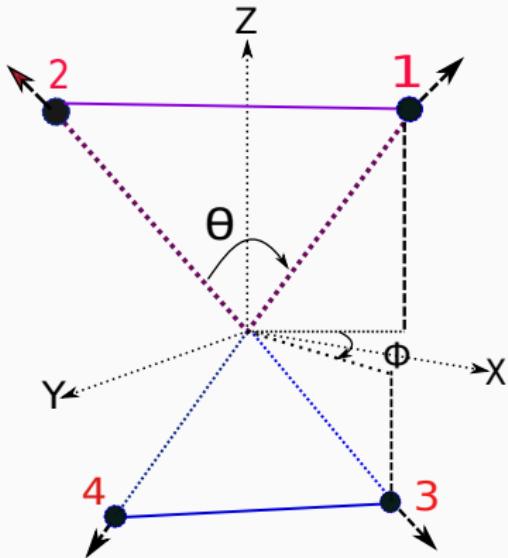




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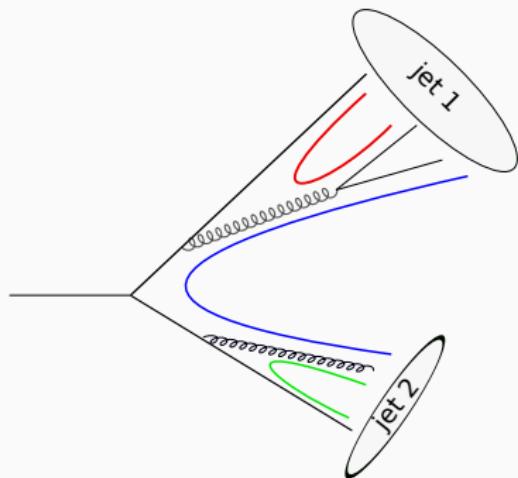


1. Lorentz invariant frame - the parallel frame



Schematic view of two strings in the parallel frame,
1,2,3,4 are partons(string-ends),
 θ = opening angle, ϕ = skew angle.
All partons are considered as massless!!

Jets in parallel frames

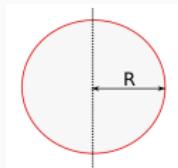


- Jets → quarks and gluons
- Interaction with partons following rule of least string length → modifies initial energy of jets

2. Interaction energy between strings

1. A string of radius R :

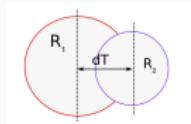
$$\text{Field } E(r_{\perp}) = C \exp\left(-\frac{r_{\perp}^2}{2R^2}\right) \quad (1)$$



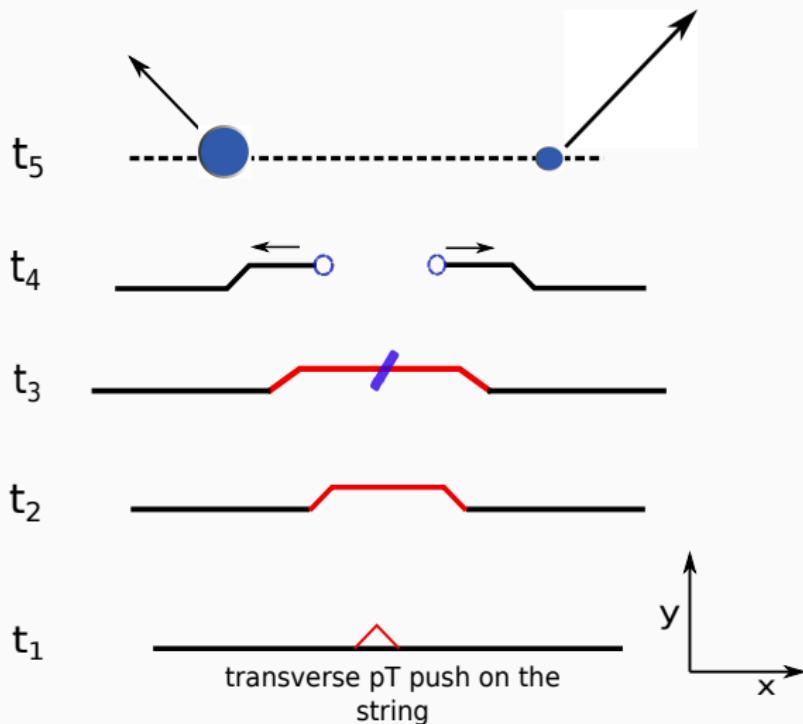
2. Force $f(d_{\perp})$ per unit length:

$$f(d_{\perp}) = \frac{dE_{int}}{dd_{\perp}} = \frac{g\kappa d_{\perp}}{R^2} \exp\left(-\frac{d_{\perp}^2(t)}{4R^2}\right) \quad (2)$$

where g is a tunable parameter.

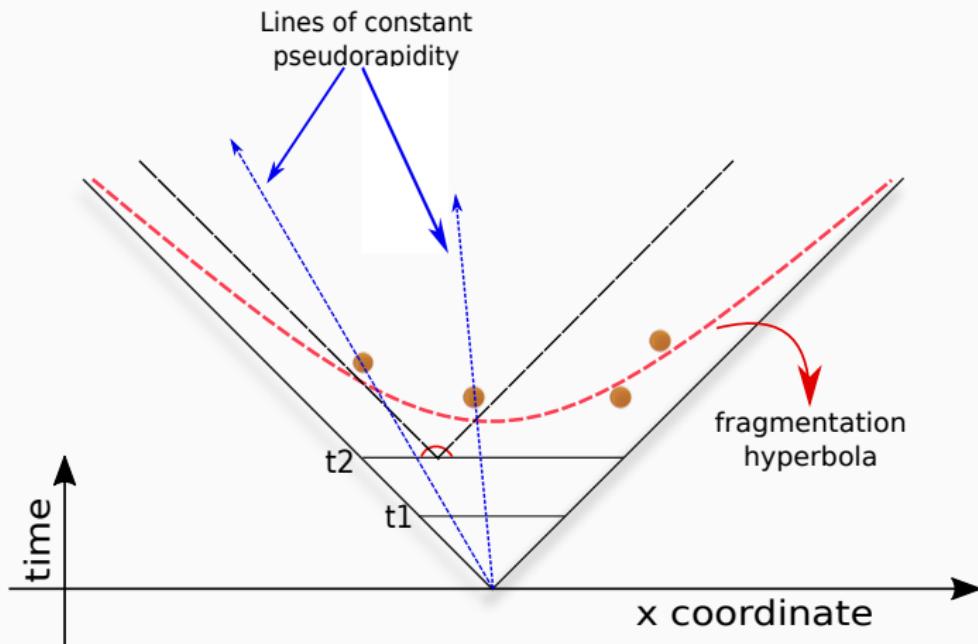


3. 'Push' distribution among hadrons



Parton vertices and hadronization

$$\text{pseudorapidity } \eta = -\ln(\tan \frac{\theta}{2})$$



Preliminary results

What are we looking at?

$$1. S_N = \frac{1}{N(N-1)} \frac{d^2 N^{signal}}{d\Delta\phi d\Delta\eta}$$

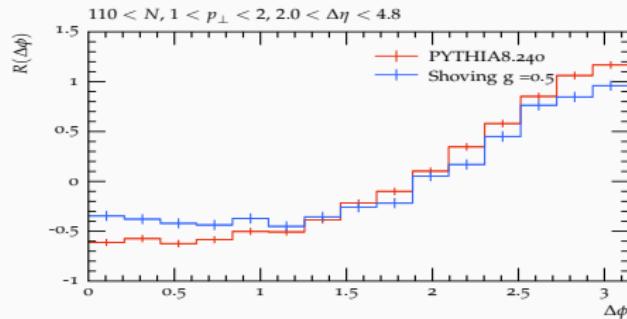
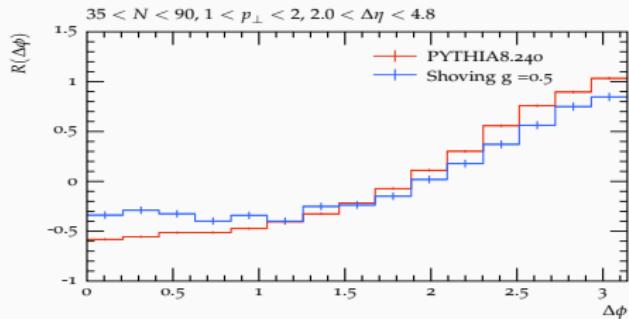
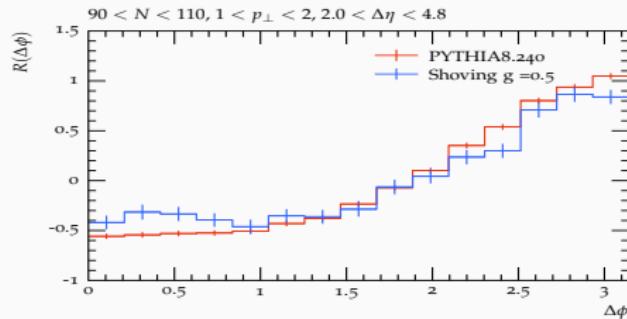
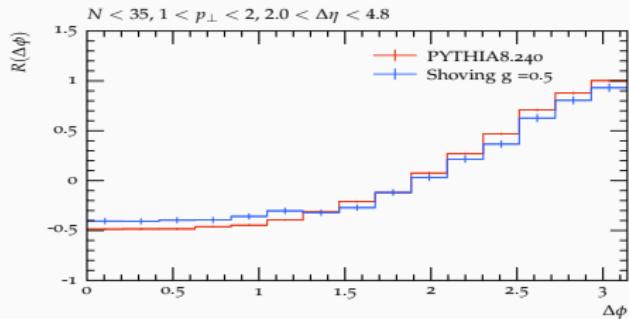
$$2. B_N = \frac{1}{N^2} \frac{d^2 N^{mixed}}{d\Delta\phi d\Delta\eta}$$

$$3. R(\phi) = \left\langle (\langle N \rangle - 1) \left(\frac{S_N}{B_N} - 1 \right) \right\rangle$$

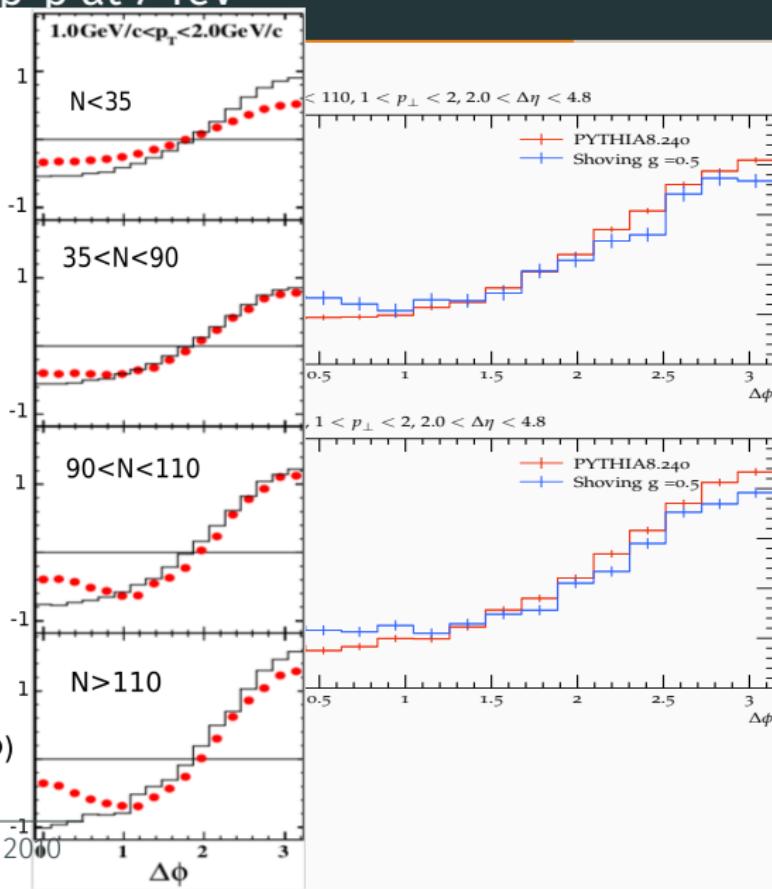
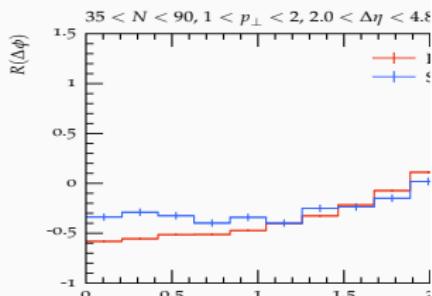
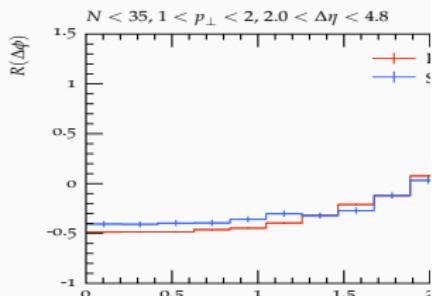
where $\langle N \rangle$ is the number of tracks per event averaged over the multiplicity bin, and the final $R(\Delta\eta, \Delta\phi)$ is found by averaging over multiplicity bins

- Analysis follows from: **Observation of Long-Range, Near-Side Angular Correlations in Proton-Proton Collisions at the LHC**, CMS Collaboration, arXiv:1009.4122v1 [hep-ex] 21 Sep 2010.

Di-hadron correlations in p-p at 7 TeV

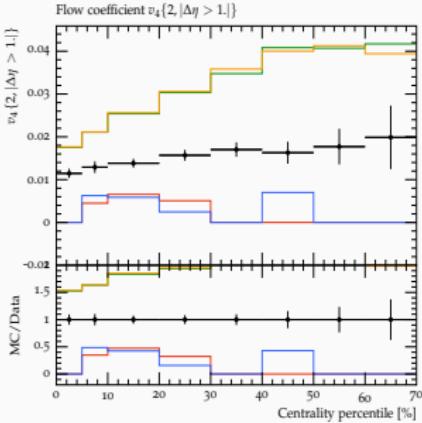
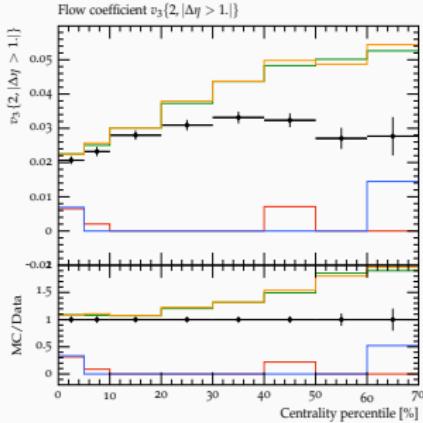
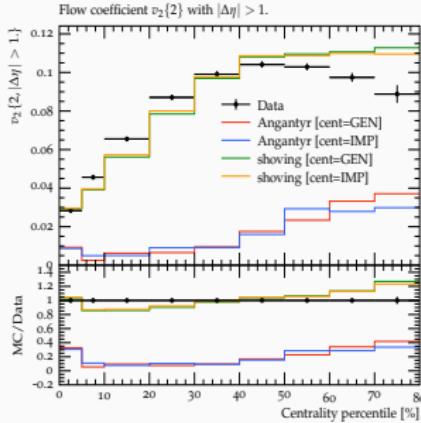


Di-hadron correlations in p-p at 7 TeV



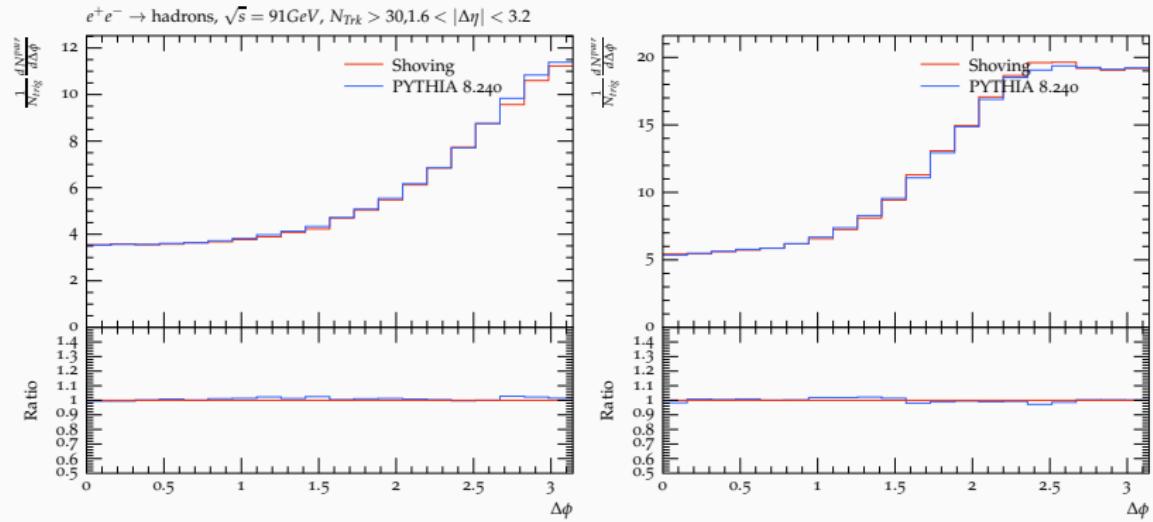
¹arXiv:1009.4122v1 [hep-ex] 21 Sep 2010

Flow coefficients in Pb-Pb at 5.02 TeV

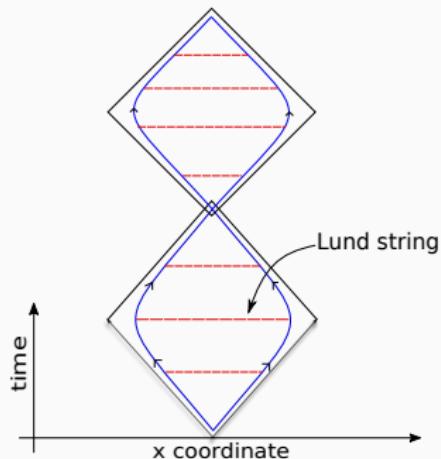


v_2 , v_3 and v_4 coefficients calculated using Angantyr and shoving
 $g=0.25$, for generated centrality bins and impact parameter bins.

Two particle correlations in $e^+ - e^-$ at 91 GeV



Heavy quark scenario



1. **Motivation:** Strangeness and charmonium production in string picture
2. **Aim:** To study hadron production ratios
3. **Changes:** Mass term in kinematics, $v < c$

Conclusions

1. Summary

- Shoving gives an observable collective effect in high multiplicity p-p, p-A and A-A events
- No corresponding effect observed for $e^+ - e^-$ collisions

2. Next steps

- Study of initial state anisotropy and its effects
- Study with jets: search for jet quenching
- Heavy quarks and hadron production ratios

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Extras

Fourier series and coefficients

1. Azimuthal angle ϕ is in the range 0 to 2π
2. $R(\phi) = \frac{dN}{d\phi} \rightarrow$ periodic function representing azimuthal distribution
3. Fourier series representation:
$$r(\phi) = \frac{a_0}{2\pi} + \frac{1}{\pi} \sum_{n=1}^{\infty} [a_n \cos(n\phi) + b_n \sin(n\phi)],$$
where $a_n = \int_{-\pi}^{\pi} r(\phi) \cos(n\phi) d\phi, n \geq 0$ and
$$b_n = \int_{-\pi}^{\pi} r(\phi) \sin(n\phi) d\phi, n \geq 1$$
4. When there is a finite number of particles, these becomes
$$a_n = \sum r(\phi_i) \cos(n\phi_i), n \geq 0$$
 and $b_n = \sum r(\phi_i) \sin(n\phi_i), n \geq 1$ where ϕ_i is the azimuthal angle of the i-th particle.

Fourier series and coefficients(contd.)

1. Introduce w_n and Ψ_n : $w_n = \sqrt{a_n^2 + b_n^2}$, $\frac{-\pi}{n} \leq \Psi_n < \frac{\pi}{n}$,
2. $a_n = w_n \cos(n\Psi_n)$ $b_n = w_n \sin(n\Psi_n)$
3. $a_n \cos(n\phi) + b_n \sin(n\phi) = w_n (\cos(n\Psi_n) \cos(n\phi) + \sin(n\Psi_n) \sin(n\phi))$
= $w_n \cos(n(\phi - \Psi_n))$
4. Therefore,

$$a_0 = w_0 \cos(0\Psi_n) = w_0 = \int_{-\pi}^{\pi} \frac{dN}{d\phi} d\phi = N \quad (3)$$

5. The Fourier series representation can be written as:

$$r(\phi) = \frac{w_0}{2\pi} (1 + \sum_{n=1}^{\infty} 2 \frac{w_n}{w_0} \cos(n(\phi - \Psi_n))) \quad (4)$$

Fourier series and coefficients(contd.)

1. The average of $\cos(n\phi) = \cos(n(\phi - \Psi_n))$ gives the coefficients of the series:

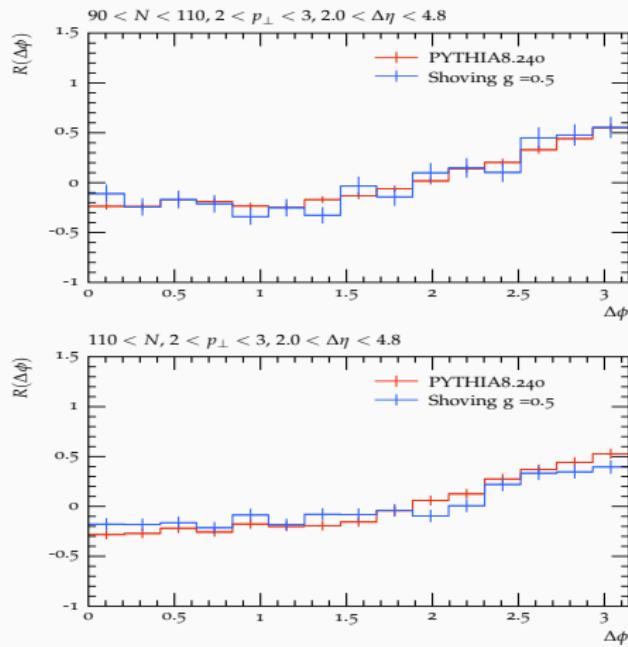
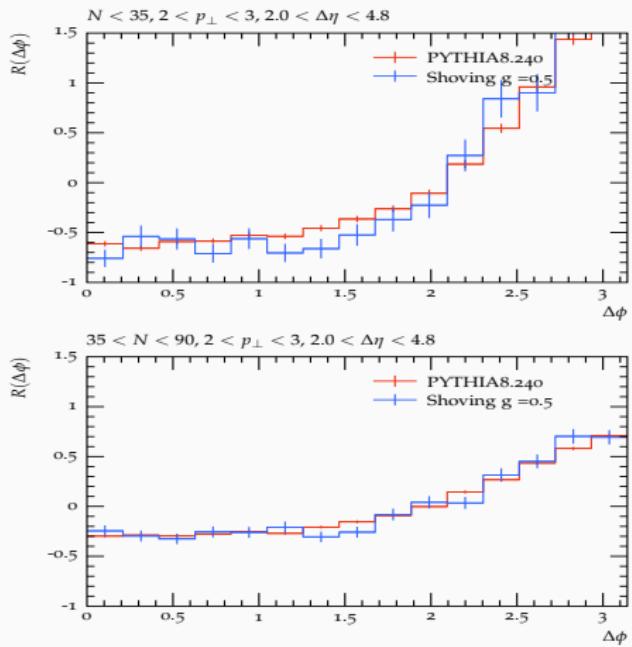
$$\langle \cos(n\Delta\phi) \rangle = \frac{\int_{-\pi}^{\pi} \cos(n\Delta\phi)(1 + \sum_m 2v_m \cos(m\Delta\phi)) d\phi}{\int_{-\pi}^{\pi} (1 + \sum_m 2v_m \cos(m\Delta\phi)) d\phi} \quad (5)$$

2. We get $\int_{-\pi}^{\pi} \cos(n\Delta\phi)(1 + \sum_m 2v_m \cos(m\Delta\phi)) d\phi = 2v_n\pi$ and $\int_{-\pi}^{\pi} (1 + \sum_m 2v_m \cos(m\Delta\phi)) d\phi = 2\pi$

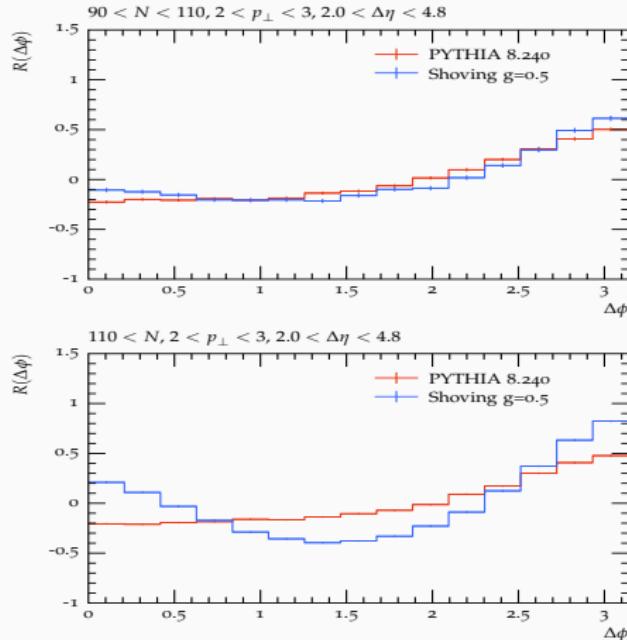
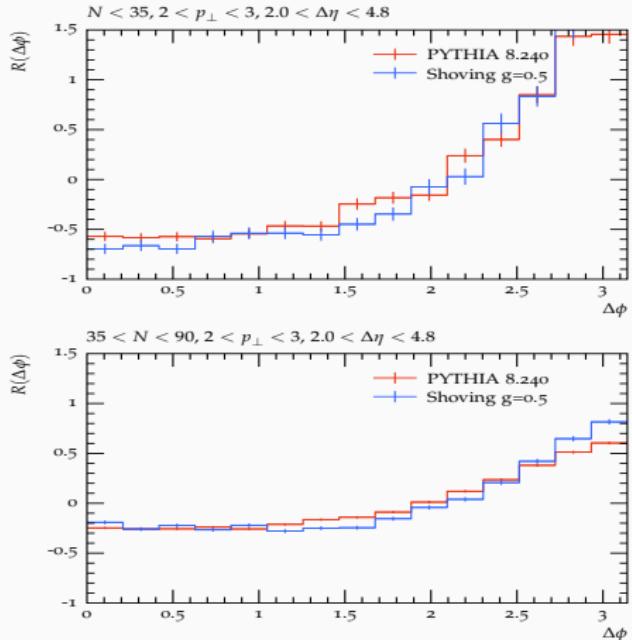
3.

$$\langle \cos(n\Delta\phi) \rangle = \frac{2v_n\pi}{2\pi} = v_n \quad (6)$$

Di-hadron correlations in p-p at 7 TeV

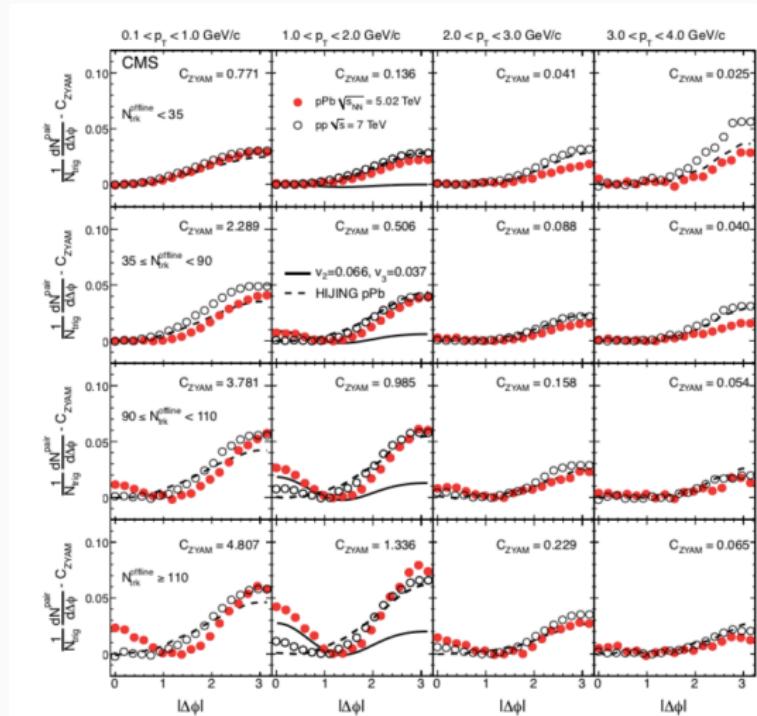


Di-hadron correlations in p-Pb

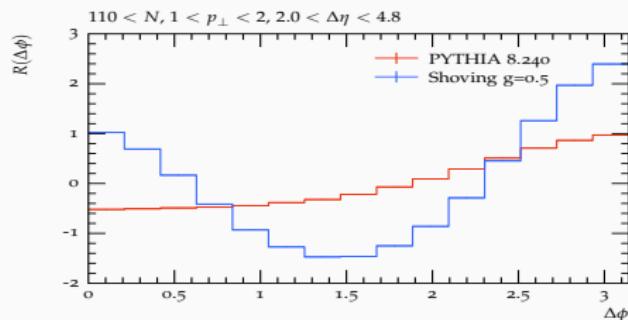
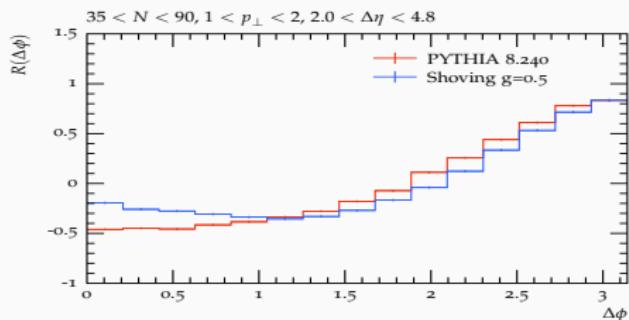
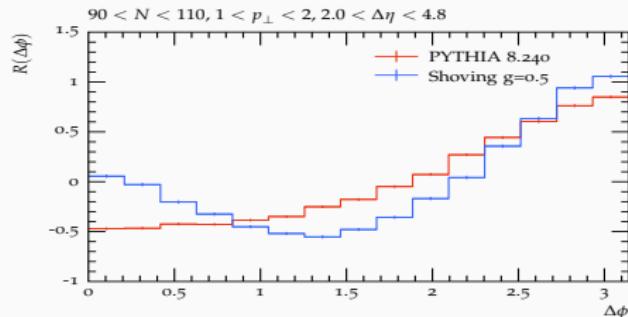
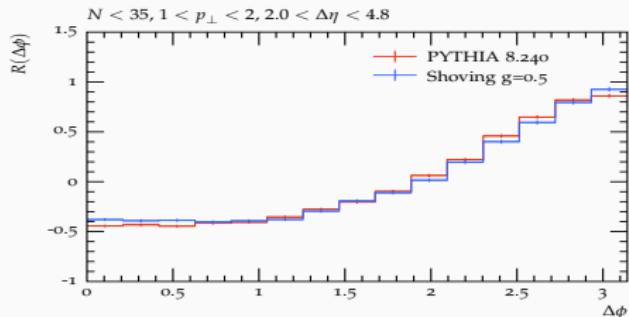


Note: Rivet analysis used is for p-p!

Di-hadron correlations in p-Pb



Di-hadron correlations in p-Pb



Note: Rivet analysis used is for p-p!