Probability to emit a gluon:

$$dP(z,k_t) = \frac{\alpha_s}{\pi} \frac{2C_R}{z} \frac{1}{k_t} dz dk_t$$

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The probability of radiation  $P = \int dP \gg 1$ ! (from large logarithms of  $\frac{p_{T,hard}}{\Lambda_{QCD}}$ ) To resum these emissions: parton showers = jets.



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How will jets change in heavy-ion collisions?



Adam Takacs



Adam Takacs

## Quenching effects in the cumulative jet spectrum

### <u>Adam Takacs</u><sup>\*</sup> University of Bergen (Norway) Konrad Tywoniuk University of Bergen (Norway)



\*adam.takacs@uib.no

• Definition:

$$R_{AA}(p_T) = \frac{\frac{d\sigma^{med}}{dp_T}(p_T)}{\left| \frac{d\sigma^{vac}}{dp_T}(p_T) \right|}$$

- $R_{AA}$ : Compares jets in vacuum to jets in medium at the **same**  $p_T$ .
- Jet with  $p_T$  in medium loose energy and ends up with  $p_T - \varepsilon$ .
- Complication 1:  $R_{AA}$  doesn't compare the "same" jets!
- The spectrum is steeply falling  $n \gg 1$ .

$$\frac{d\sigma}{dp_T} \sim p_T^{-n}$$

• Complication 2:

 $R_{AA}$  is sensitive to n (bias on energy loss)!





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### Other possibilities

- To decrease the bias: Use the cumulative distribution instead.
- What is the "original"  $p_{\rm T}$ ? Instead of the spectrum, use the probability (cumulative).
- Quantile procedure

[J.Brewer et. al. Phys.Rev.Lett. 122, 222301 (2019)]



$$Q_{med}(p_T^{q,v}) = \frac{p_T^{q,m}}{p_T^{q,v}}$$



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### Building up the quenched jet spectrum

### 1. The quenching weight



### The quenched spectrum: the quenching weight

[Baier, Dokshitzer, Mueller, Schiff (1998), Salgado, Wiedemann (2001)]

The quenched spectrum (probability  $\mathcal P$  of loosing  $\boldsymbol \varepsilon$  energy)

$$\frac{d\sigma^{med}}{dp_T}(p_T) \equiv \int_0^\infty d\varepsilon \,\mathcal{P}(\varepsilon) \frac{d\sigma^{vac}}{dp_T}(p_T + \varepsilon) \approx \frac{d\sigma^{vac}}{dp_T}(p_T) \int_0^\infty d\varepsilon \,\mathcal{P}(\varepsilon) e^{-\frac{n\varepsilon}{p_T}} \frac{d\sigma^{vac}}{dp_T}(p_T) \sim p_T^{-n}$$

The  $R_{AA}$  is the quenching weight

$$R_{\rm med}(p_T) \equiv \frac{d\sigma^{\rm med}}{dp_T}(p_T) / \frac{d\sigma^{\rm vac}}{dp_T}(p_T) \approx \int_0^\infty d\varepsilon \,\mathcal{P}(\varepsilon) e^{-\frac{n\varepsilon}{p_T}} \equiv \mathcal{Q}_{med}(p_T)$$



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What is  $\mathcal{P}(\varepsilon)$ ?



### Building up the quenched jet spectrum

### 1. Medium Induced Emissions



Single parton, single medium induced emission

$$\mathcal{P}_{>}^{(0)}(\varepsilon) \approx \frac{dI_{>}}{d\varepsilon}$$

Single parton, multiple induced emission [JHEP09 (2001) 033]

$$\mathcal{P}_{>}^{(0)}(\varepsilon) = \sum_{n=0}^{\infty} \frac{1}{n!} \left[ \prod_{j=1}^{n} \int d\omega_{j} \frac{dI_{>}}{d\omega_{j}} \right] \delta\left(\varepsilon - \sum_{j=1}^{n} \omega_{j}\right) e^{-\int d\omega_{j} \frac{dI_{>}}{d\omega_{j}}}$$
$$\mathcal{Q}_{>}^{(0)}(p_{T}) = \exp\left[ -\int_{0}^{\infty} d\omega \left(1 - e^{-\frac{n\omega}{p_{T}}}\right) \frac{dI_{>}}{d\omega} \right]$$

Multi parton (jet), multiple induced emission [Phys.Rev.D98 (2018) 051501]

$$Q_{>}^{jet}(p_T) \approx Q_{>}^{(0)}(p_T)\mathcal{C}(p_T,R)$$



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Adam Takacs

#### Lund ATP Coffee seminar 2021

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(Pool



Single parton, single medium induced emission





### Comparison to data

To include more effects:

- Assuming factorization: PDF and nPDF effects, quark/gluon ratio
- Medium resolution and color coherence effects
- Broadening of the induced gluons in/out of the cone
- Thermalizing soft gluons
- Energy loss from elastic scattering

Missing:

- Geometry and time dependence
- Fluctuations

## 0-10% PbPb @ 5.02TeV, $|\eta| < 2.8$

see RAA also in arXiv:2101.01742]





### 0-10% PbPb @ 5.02TeV, $|\eta| < 2.8$





### Quenching weight





## Different observable possibilities

	Name	def.	n  sensit.	statistics	cartoon
•	Nuclear modification:	$R_{AA} = \frac{\sigma_{AA}}{\sigma_{pp}}\Big _{p_T}$	ok	ok	
•	Quantile ratio:	$Q_{AA} = \frac{p_T^{AA}}{p_T^{pp}} \bigg _{\Sigma^{eff}}$	smaller	better	
•	Cumulative- $R_{AA}$ : Behaves as $\sim Q_{AA}^{n-1}$	$\tilde{R}_{AA} = \frac{\Sigma_{AA}}{\Sigma_{pp}} \bigg _{p_T}$	ok	better	
•	Pseudo-Quantile: Behaves as $\sim R_{AA}^{-1/n}$	$\tilde{Q}_{AA} = \frac{p_T^{AA}}{p_T^{pp}} \bigg _{\sigma^{eff}}$	smaller	ok	
•	Spectrum shift: Equivalent to $ ilde{Q}_{AA}$	$\sigma^{AA}(p_T) = \sigma^{pp}(p_T + S)$	smaller	ok	

### Thank you for your attention!



## Modern Jet Quenching - Medium Induced Emission

[Zakharov, BDMPS, GLV (1996-2000) - Blaizot, Iancu, Salgado (2012-)]

Scattering on the medium (of an energetic parton):

- Colored background  $\mathcal{A}_0(t, x)$
- Energy is conserved  $(p^+)$ , transverse kick (p)
- Multiple scatterings

Keeping space-time: partial Fourier space  $(p^+, p, p^-) \rightarrow (p^+, x, t)$ 

2-d Schrodinger equation with imaginary potential

$$\begin{bmatrix} i\partial_t + \frac{\partial_i^2}{2p^+} + g\mathcal{A}_0(t, \mathbf{x}) \end{bmatrix} G(t, \mathbf{x}, t_0, \mathbf{x}_0 | p^+) = i\delta(\mathbf{x} - \mathbf{x}_0)\delta(t - t_0)$$
$$G(t, \mathbf{x}, t_0, \mathbf{x}_0 | p^+) = \int_{\mathbf{x}_i}^{\mathbf{x}_f} \mathcal{D}\mathbf{r} \ e^{i\frac{p^+}{2}\int_{t_i}^{t_f} ds \ \dot{r}^2(s)} \mathcal{T}e^{-ig\int_{t_i}^{t_f} ds \ \mathcal{A}_0(r(s))}$$





### Modern Jet Quenching - Medium Induced Emissions

Medium induced emission:

$$\omega \frac{dI}{d\omega d^2 \mathbf{k}} = \frac{2\alpha_s C_R}{(2\pi)^2 \omega^2} \operatorname{Re} \int_0^L dt \int_t^L d\bar{t} \int d^2 \mathbf{x} \, e^{-\int ds \, v(\mathbf{x})} \partial_{\mathbf{x}} \partial_{\mathbf{y}} \, G(\bar{t}, \mathbf{x}, t, \mathbf{y} \to 0 \mid \omega)$$
$$G(t, \mathbf{x}, t_0, \mathbf{x}_0 \mid p^+) = \int_{\mathbf{x}_i}^{\mathbf{x}_f} \mathcal{D} \mathbf{r} \, e^{\int_{t_i}^{t_f} ds \left[ i \frac{p^+}{2} \dot{r}^2 - v(r, s) \right]}$$

Elastic scattering potential:

$$v(\mathbf{x}) = N_c \int \frac{d^2 \mathbf{q}}{(2\pi)^2} \frac{d^2 \sigma_{el}}{d^2 \mathbf{q}} (1 - e^{i\mathbf{x}\mathbf{q}})$$
$$\hat{\sigma}_c(t) = \frac{d^2 \sigma_{el}}{d^2 \mathbf{q}} (1 - e^{i\mathbf{x}\mathbf{q}})$$

$$v(\mathbf{x}) = \frac{\hat{q}_0(t)}{4} \mathbf{x}^2 \qquad \qquad \frac{d^2 \sigma_{el}^{G-W}}{d^2 \mathbf{q}} = \frac{4\pi}{N_c} \frac{\hat{q}_0(t)}{(\mathbf{q}^2 + \mu^2)^2}$$

Medium resolution:



### Modern Jet Quenching - Medium Induced Emissions

[Blaizot, Iancu, Mehtar-Tani (2014)]

Medium induced parton shower

$$\partial_t D(x,t) = \frac{1}{t_*} \int dz \,\mathcal{K}(z) \left[ \sqrt{\frac{z}{x}} D\left(\frac{x}{z},t\right) - \frac{z}{\sqrt{x}} D(x,t) \right]$$
$$t_*^{-1} = \bar{\alpha} \sqrt{\frac{\hat{q}}{E}} \qquad \mathcal{K}(z) = \frac{[1-z+z^2]^{5/2}}{[z(1-z)]^{3/2}}$$

With transverse broadening  $D(x, \mathbf{k}, t) = (2\pi)^2 x \frac{dN}{dxd^2 \mathbf{k}}$ 

$$\partial_t D(\mathbf{x}, \mathbf{k}, \mathbf{t}) = \frac{1}{\mathbf{t}_*} \int dz \,\mathcal{K}(z) \left[ \frac{1}{z^2} \sqrt{\frac{z}{x}} D\left(\frac{x}{z}, \frac{\mathbf{k}}{z}, t\right) - \frac{z}{\sqrt{x}} D(x, \mathbf{k}, t) \right] + \int \frac{d^2 \mathbf{q}}{(2\pi)^2} \mathcal{C}(\mathbf{q}) D(x, \mathbf{k} - \mathbf{q}, t)$$



## Comparison to data

### To include more effects:



- Geometry and time dependence
- Fluctuations

## Quenching weight



