

Relativistic Corrections to Electronic Scattering Lengths. Results for Strontium.

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Definition

Scattering length is a parameter used to describe low-energy electron-atom collisions. It is defined as the radius of a rigid sphere in the zero-energy total cross-section. The sign of a scattering length represents the type of interaction: positive for repulsion and negative for attraction.

$$E = \frac{1}{2} \alpha_d \vec{\mathbf{E}}^2$$

$$V(r) = -\frac{1}{2} \frac{\alpha_d}{r^4} = -\frac{1}{2} \frac{\beta^2}{r^4}$$

$$\vec{\mathbf{E}} = \frac{e\vec{\mathbf{r}}}{\sqrt{(r^3 + r_0^3)^2}},$$

$$V_{pol}(r) = -\frac{1}{2} \frac{\alpha_d r^2}{(r^3 + r_0^3)^2}$$

$$V_{pol}(r) = -\frac{1}{2} \frac{\alpha_d r^2}{(r^3 + r_0^3)^2} - \frac{1}{2} \frac{\alpha'_q r^4}{(r^5 + r_0^5)^2}$$

Parameter β is often called characteristic quantum length.

Dirac's equation - zero energy with polarisation

$$\left(\frac{d}{dr} + \frac{\kappa}{r}\right)P_{\kappa}(r) = \left(\frac{2}{\alpha} + \frac{\alpha\beta^2}{2r^4}\right)Q_{\kappa}(r)$$

$$\left(\frac{d}{dr} - \frac{\kappa}{r}\right)Q_{\kappa}(r) = -\frac{\alpha\beta^2}{2r^4}P_{\kappa}(r)$$

$$\frac{d^2P(r)}{dr^2} + \frac{1}{r} \frac{4\alpha^2\beta^2}{2r^4 + \alpha^2\beta^2} \frac{dP(r)}{dr} - \frac{1}{r^2} \frac{4\alpha^2\beta^2}{2r^4 + \alpha^2\beta^2} P(r) = \left(-\frac{\beta^2}{r^4} - \frac{\alpha^2\beta^4}{4r^8}\right)P(r)$$

$$\frac{d^2Q(r)}{dr^2} + \frac{4}{r} \frac{dQ(r)}{dr} + \frac{2}{r^2} Q(r) = \left(-\frac{\beta^2}{r^4} - \frac{\alpha^2\beta^4}{4r^8}\right)Q(r)$$

Solution for small component

$$\begin{aligned} Q(r) &= \frac{\alpha^2}{4} \left(\frac{a}{r} + \frac{\beta^2}{r^2} - \frac{\beta^2}{2!} \frac{a}{r^3} - \frac{\beta^4}{3!} \frac{1}{r^4} + \frac{\beta^4}{4!} \frac{a}{r^5} \right. \\ &\quad \left. + \frac{\beta^6}{6!} \frac{1}{r^6} - \frac{\beta^6}{6!} \frac{a}{r^7} - \frac{4! \cdot \frac{\alpha^2 \beta^2}{4}}{6!} \frac{a}{r^7} + \dots \right) \\ &= \frac{\alpha^2}{4} \frac{a}{r} \left(1 - \frac{\beta^2}{2!} \frac{1}{r^2} + \frac{\beta^4}{4!} \frac{1}{r^4} - \frac{\beta^6}{6!} \frac{1}{r^6} + \dots \right) \\ &\quad + \frac{\alpha^2}{4} \frac{\beta}{r} \left(\frac{\beta}{r} - \frac{\beta^3}{3!} \frac{1}{r^3} + \frac{\beta^5}{5!} \frac{1}{r^5} - \dots \right) - \frac{4! \cdot \frac{\alpha^2 \beta^2}{4}}{6!} \frac{a}{r^7} \dots \\ Q(r) &= \frac{\alpha^2}{4} \frac{a}{r} \cos\left(\frac{\beta}{r}\right) + \frac{\alpha^2}{4} \frac{\beta}{r} \sin\left(\frac{\beta}{r}\right) + \alpha^4 O\left(\frac{1}{r^7}\right) \end{aligned}$$

Solution for large component

$$\begin{aligned}P(r) &= r - a - \frac{\beta^2}{2} \frac{1}{r} + a \frac{\beta^2}{3!} \frac{1}{r^2} + \frac{\beta^4}{4!} \frac{1}{r^3} - a \left(\frac{\beta^4}{5!} + \frac{4! \cdot \frac{\alpha^2 \beta^2}{4}}{5!} \right) \frac{1}{r^4} + \dots \\&= r \left(1 - \frac{\beta^2}{2} \frac{1}{r^2} + \frac{\beta^4}{4!} \frac{1}{r^4} - \dots \right) - \frac{ar}{\beta} \left(\frac{\beta}{r} - \frac{\beta^3}{3!} \frac{1}{r^3} + \frac{\beta^5}{5!} \frac{1}{r^5} - \dots \right) \\&\quad - a \frac{4! \cdot \frac{\alpha^2 \beta^2}{4}}{5!} A \frac{1}{r^4} + \dots \\&= r \cos\left(\frac{\beta}{r}\right) - \frac{ar}{\beta} \sin\left(\frac{\beta}{r}\right) + O\left(\frac{1}{r^4}\right) \\P(r) &= \frac{r}{a} \cos\left(\frac{\beta}{r}\right) - \frac{r}{\beta} \sin\left(\frac{\beta}{r}\right) + \alpha^2 O\left(\frac{1}{r^4}\right)\end{aligned}$$

- ① Check whether the solution tends to straight line
- ② Check whether adding additional coefficients makes better results
- ③ Check whether we can get back parameter β
- ④ Check the correctness of a formula with sine and cosine
- ⑤ Check whether we can spot the difference between non-relativistic and relativistic equations

Some information about strontium:

- Has closed-subshell structure
- $\langle r \rangle = 4.55369$ a.u.
- $\alpha_d = 197.2$ (Schwerdtfeger, P. and Nagle, J. K. 2018)
- $\beta = 14.04279$ ($\beta^2 = \alpha_d$)
- Known scattering lengths: experimental $a = -12.5$ a.u., (Giannakeas, P. and Eiles, Matthew T. and Robicheaux, F. and Rost, Jan M. 2020), $a = -13.2$ a.u (De Salvo 2015) theoretical $a = -18$ a.u (Bartschatt, Sadeghpour 2003)

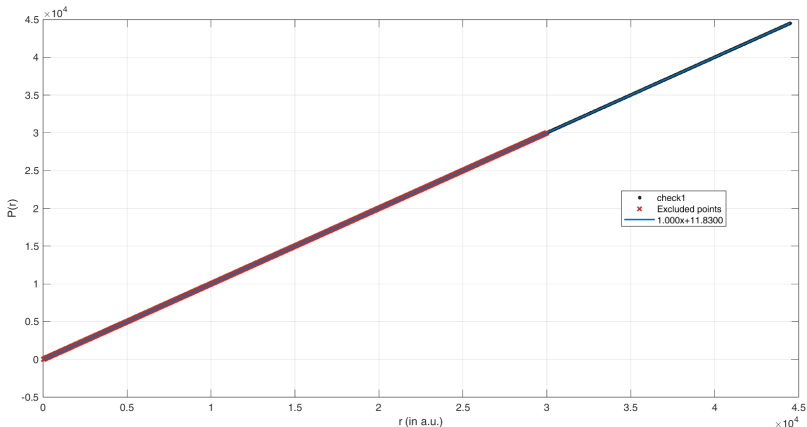


Figure: Claim: Is the straight line approximation correct? Approximation:
 $P(r) = 1.000r + 11.8300 \rightarrow a = -11.8300$

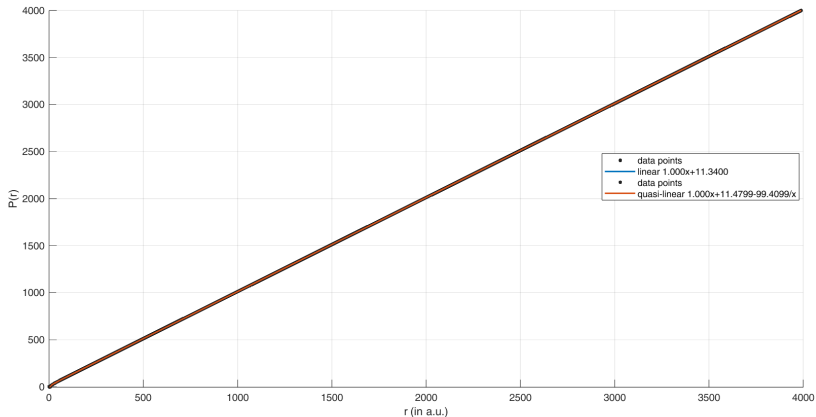


Figure: Claim: Is the longer approximation better? Approximations:
 $P(r) = 1.000r + 11.3400$ and $P(r) = 1.000r + 11.4799 - 99.4099/r$

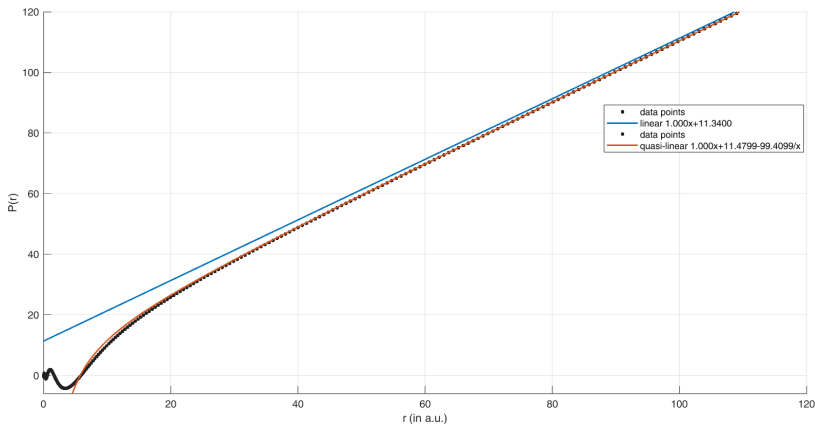


Figure: Claim: Is the longer approximation better? Approximations:
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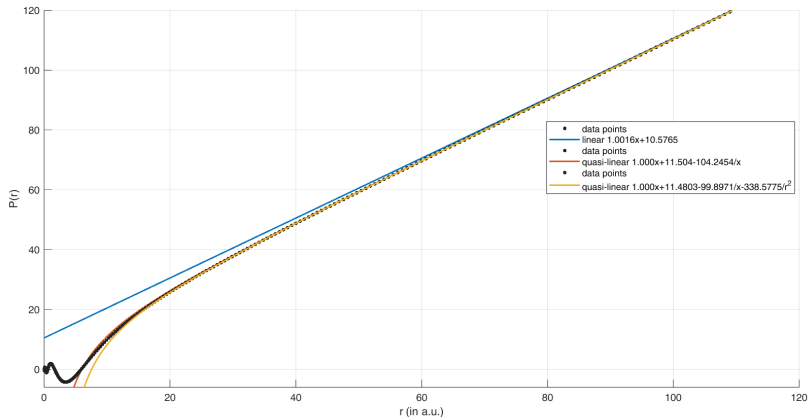


Figure: Claim: Is the longer approximation better? Approximations:
 $P(r) = 1.0016r + 10.5765$, $P(r) = 1.000r + 11.504 - 104.2454/r$ and
 $P(r) = 1.000r + 11.4803 - 99.8971/r - 338.5775/r^2$

We know that the coefficient next to r^{-1} is equal to $\frac{-\beta^2}{2}$. From the first approximation we get $\frac{-\beta^2}{2} = -99.4099 \rightarrow \alpha_d = \beta^2 = 198.8198$
From the second approximation
 $-\frac{\beta^2}{2} = -99.8971 \rightarrow \alpha_d = \beta^2 = 199.7942$. To recall, in the code we have put $\alpha_d = \beta^2 = 197.2$

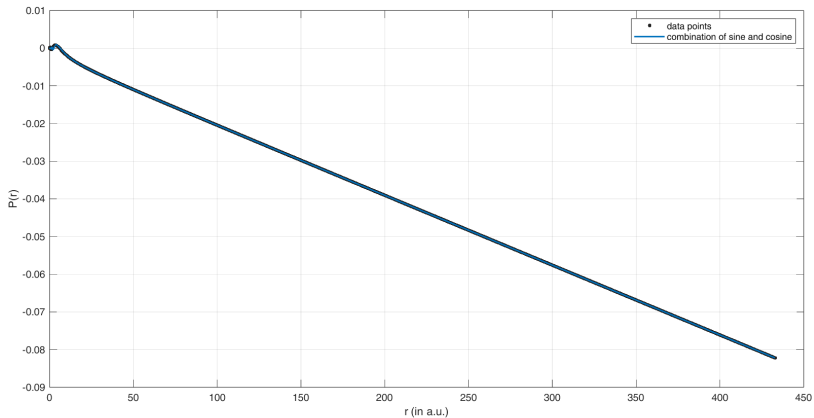


Figure: Claim: Is the sine-cosine approximation better?

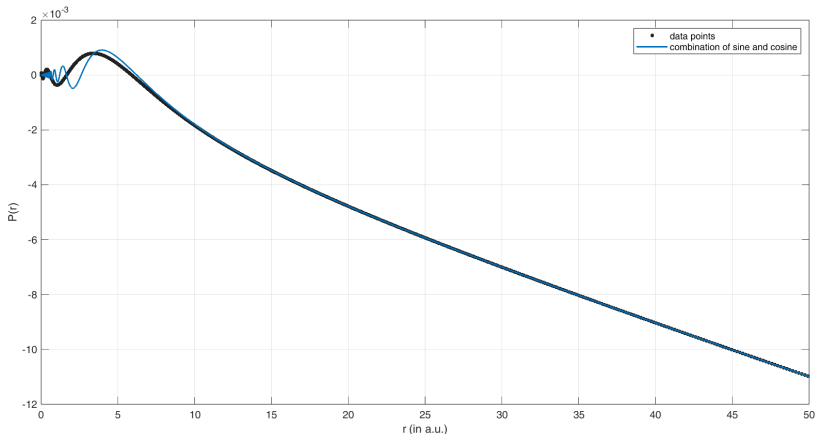


Figure: Claim: Is the sine-cosine approximation better? For this graph we have $P(r) = N\left(\frac{r}{a} \cos \frac{\beta}{r} - \frac{r}{\beta} \sin \frac{\beta}{r}\right)$ with $a = -11.4751$ a.u and $\beta = 14.0358$.

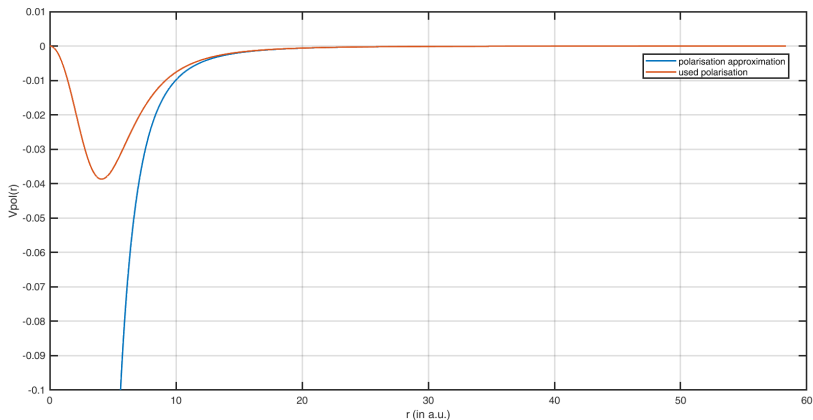


Figure: Comparison between model polarisation potential and potential used in calculations.

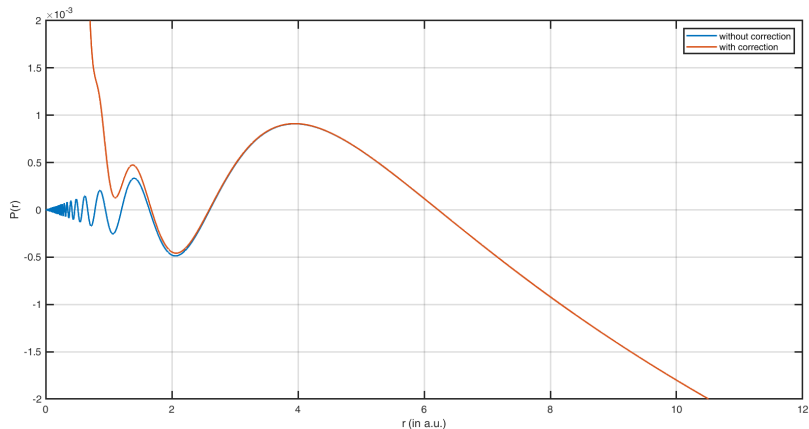


Figure: Comparison between the sine-cosine result with and without relativistic correction.

Thank you for your attention. We will work more on the continuum solutions to Dirac equations.