
Radiative data calculations for electric dipole transitions in the Yb I isoelectronic sequence (Ta IV - Pt IX) of interest to nuclear fusion diagnostics

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1. Introduction

- Ionic impurities in the fusion plasma
- Plasma diagnostics

2. HFR+CPOL preliminary results

- Pseudo-relativistic Hartree-Fock method with the core polarization corrections (HFR+CPOL)
- Model + least squares adjustment
- Radiative decay rates
- The configuration interactions and the core-valence correlations

3. MCDHF preliminary results

- Multiconfigurational Dirac-Hartree-Fock (MCDHF) method
- Model and convergence
- Radiative decay rates (comparisons)

Introduction

Ionic impurities in the fusion plasma

- Divertor will be made of tungsten: high fluxes of heat and particles (neutrons)
- Transmutation products after few years irradiation time: rhenium, osmium, tantalum, hafnium, iridium and platinum → impurities
- Brittleness of pure tungsten → alloying elements (tantalum, rhenium, titanium or vanadium)

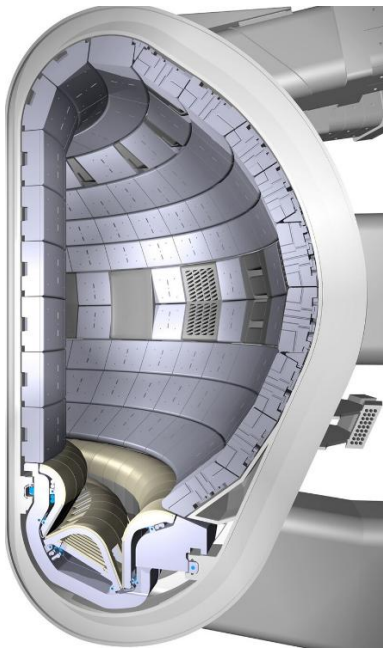
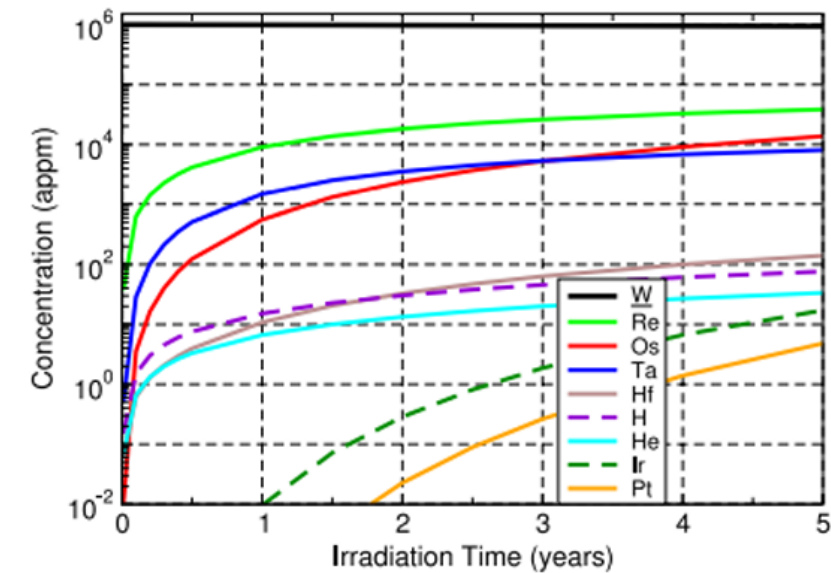


Illustration of the inside of the Tokamak.
(<https://encyclopedia.pub/entry/37629>)



Evolution of the concentration (Atomic Part Per Millions) of all the elements produced by the transmutation of W for a five-years irradiation [Gilbert, M. R and Sublet, J.Ch , 2011]

Plasma diagnostics

- Ionic impurities will contribute to the radiation losses/allow to diagnose the fusion plasma
- From observed intensity ratio and radiative parameters [HJ. Kunze, 2009] :

$$\blacktriangleright \frac{\varepsilon_Z(p \rightarrow q)}{\varepsilon_Z(p' \rightarrow q')} = \frac{\lambda_{p'q'}}{\lambda_{pq}} \frac{A_Z(p \rightarrow q)}{A_Z(p' \rightarrow q')} \frac{A_Z(p' \rightarrow)}{A_Z(p \rightarrow)} \frac{X_Z(g \rightarrow p)(T_e)}{X_Z(g \rightarrow p')(T_e)} \Rightarrow T_e$$

$$\blacktriangleright \frac{\varepsilon_Z(p \rightarrow q)}{\varepsilon_Z(p' \rightarrow q')} = \frac{\lambda_{p'q'}}{\lambda_{pq}} \frac{A_Z(p \rightarrow q)}{A_Z(p' \rightarrow q')} \frac{A_Z(p' \rightarrow)}{A_Z(p \rightarrow)} \frac{X_Z(g \rightarrow p)(T_e)}{X_Z(g \rightarrow p')(T_e)} \left[1 + n_e \frac{X_Z(p' \rightarrow)}{A_Z(p' \rightarrow)} \right] \Rightarrow n_e$$

HFR+CPOL preliminary results

HFR method

General procedure [Cowan, R. D., 1981]

- Solve $H\Psi = E\Psi$ where $H = \sum_{i=1}^N \left(-\frac{1}{2} \Delta_i + V(r_i) \right)$ (central field approximation)
- $H_i \varphi_i = E_i \varphi_i \rightarrow \varphi_i(r_i, \theta_i, \phi_i, s_i) = \frac{1}{r_i} P_{n_i l_i}(r_i) Y_{l_i}^{m_i}(\theta_i, \phi_i) \sigma_{m_{s_i}}(s_i)$
- $P_{n_i l_i}(r_i) ? \rightarrow$ solve Hartree-Fock equations (Self-Consistent Field method)
- HF equations obtained by variationnal principle on the average energy of each electronic configuration
- Atomic State Functions (ASFs) : $\Psi(\alpha, P, J, M_J) = \sum_{r=1}^{n_c} c_r \Phi(\alpha_r, P, L_r, S_r, J, M_J)$

Configuration State Functions (CSFs) are built thanks to Slater determinants

HFR+CPOL method

Core polarisation correction

- Valence electron correlations represented by configuration interactions (CI) and other correlations by core-polarisation potential
- Quinet, P. et al [1999, 2002]: pseudo potentiel have one-body and two-body part :

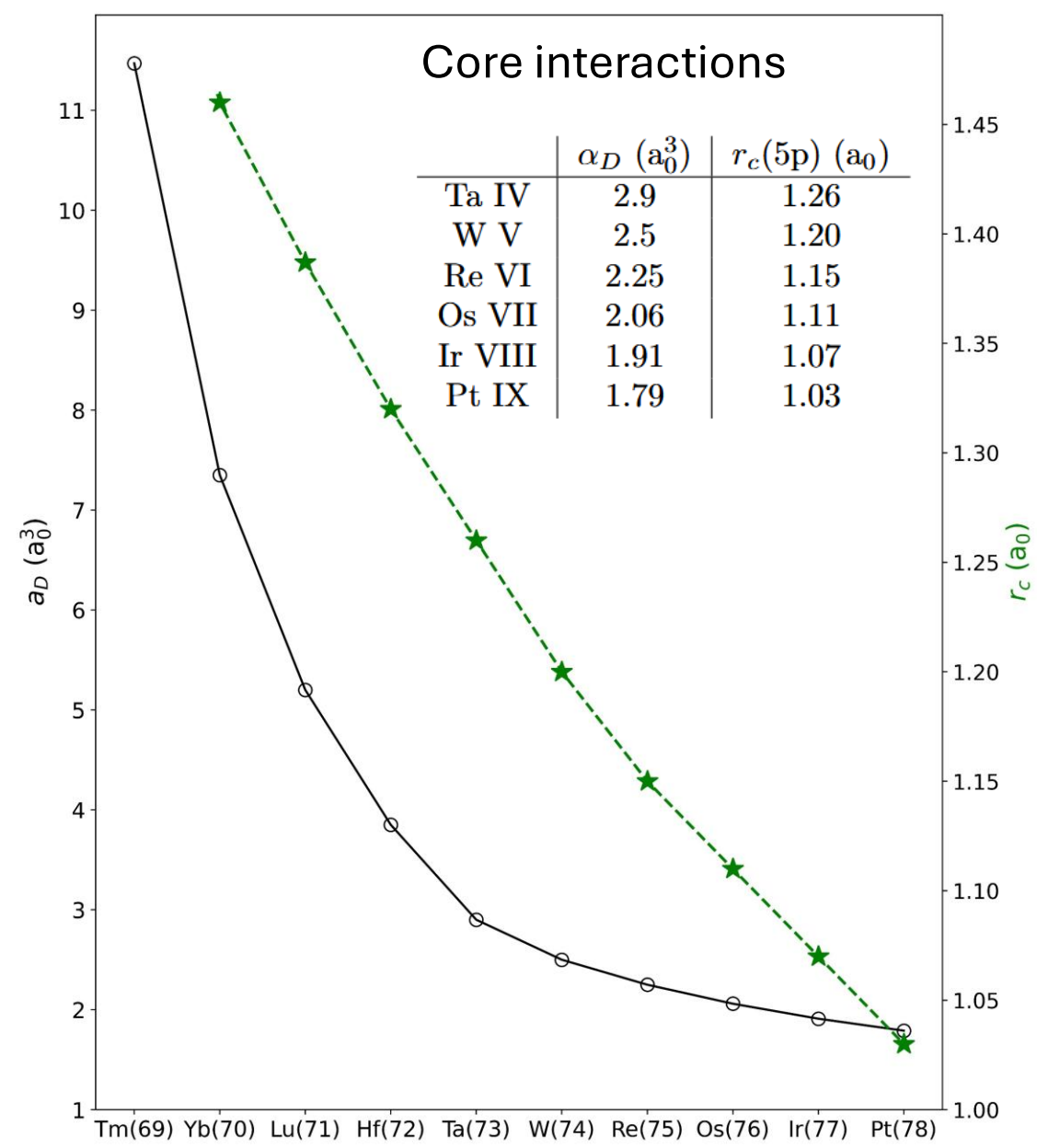
$$\text{➤ } V_{P1} = -\frac{1}{2}\alpha_D \sum_{i=1}^N \frac{r_i^2}{(r_i^2 + r_c^2)^3} \text{ and } V_{P2} = -\alpha_D \sum_{i>j} \frac{\vec{r}_i \cdot \vec{r}_j}{[(r_i^2 + r_c^2)(r_j^2 + r_c^2)]^{3/2}}$$

➤ α_D : dipole polarisability; r_c : ionic core radius

HFR+CPOL Model

Valence-Valence interactions

Even parity		Odd parity	
5d ²	6p6f	5d6p	6p6d
5d6s	6p7f	6s6p	6p7d
5d7s	6d ²	5d7p	6d7p
6s ²	6d7s	5d5f	6d5f
5d6d	6d7d	5d6f	6d6f
5d7d	7s ²	5d7f	6d7f
6s7s	7p ²	6s7p	7s5f
6s6d	7s7d	6s5f	7s6f
6s7d	7p5f	6s6f	7s7f
6p7p	7p6f	6s7f	7s7p
6p ²	7p7f	6p7s	7s7d
6p5f			



Least squares method

- Minimise difference between computed energy levels and observed ones with spin-orbit and Slater parameters
- Accuracy of the fit: $\sigma = \left[\frac{\sum_k (E^k - T^k)^2}{N_k - N_p} \right]^{1/2}$, where E^k : energy eigenvalues; T^k : observed energies; N_k : number of fitted levels; N_p : number of fitted parameters

	Even parity		Odd parity	
	nbr of fitted levels	σ (cm ⁻¹)	nbr of fitted levels	σ (cm ⁻¹)
Ta IV	14	23	44	145
W V	13	20	44	132
Re VI	13	30	29	347

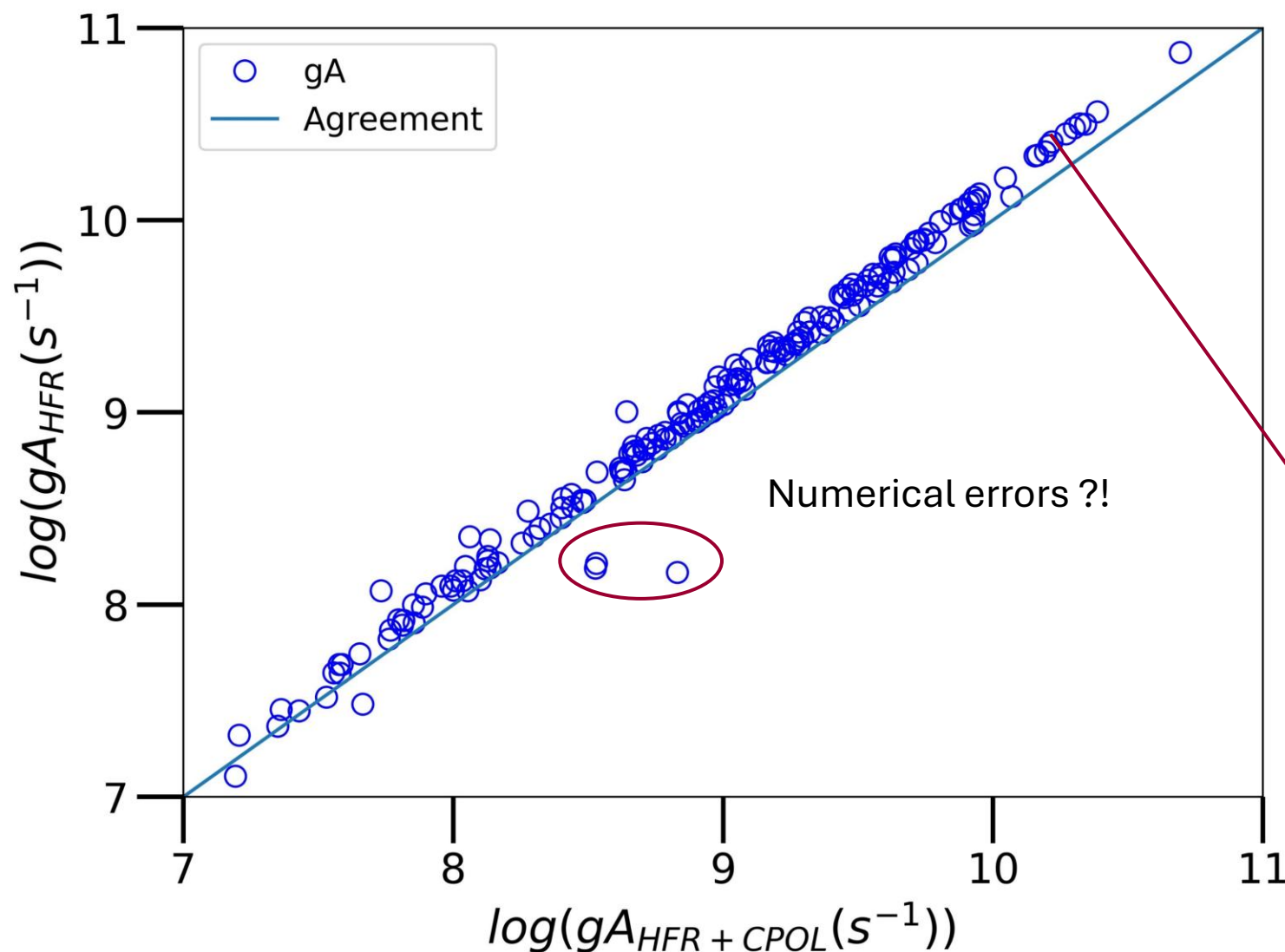
[Churilov, S.S. et al, 1996 ; Meijer, F.G and Metsch, 1978 ; Kildiyarova, R.R et al, 1996 ; Kramida , A. et al, 2024 ; Sugar, J. et al, 1994 ; Yoca, S.E et al, 2012]

Radiative decay rates

Ta IV

[Churilov, S.S. et al, 1996 ; Meijer, F.G and Metsch, 1978 ; Kildiyarova, R.R et al, 1996]

213 observed electric dipole lines
between 551.422 – 3076.060 Å



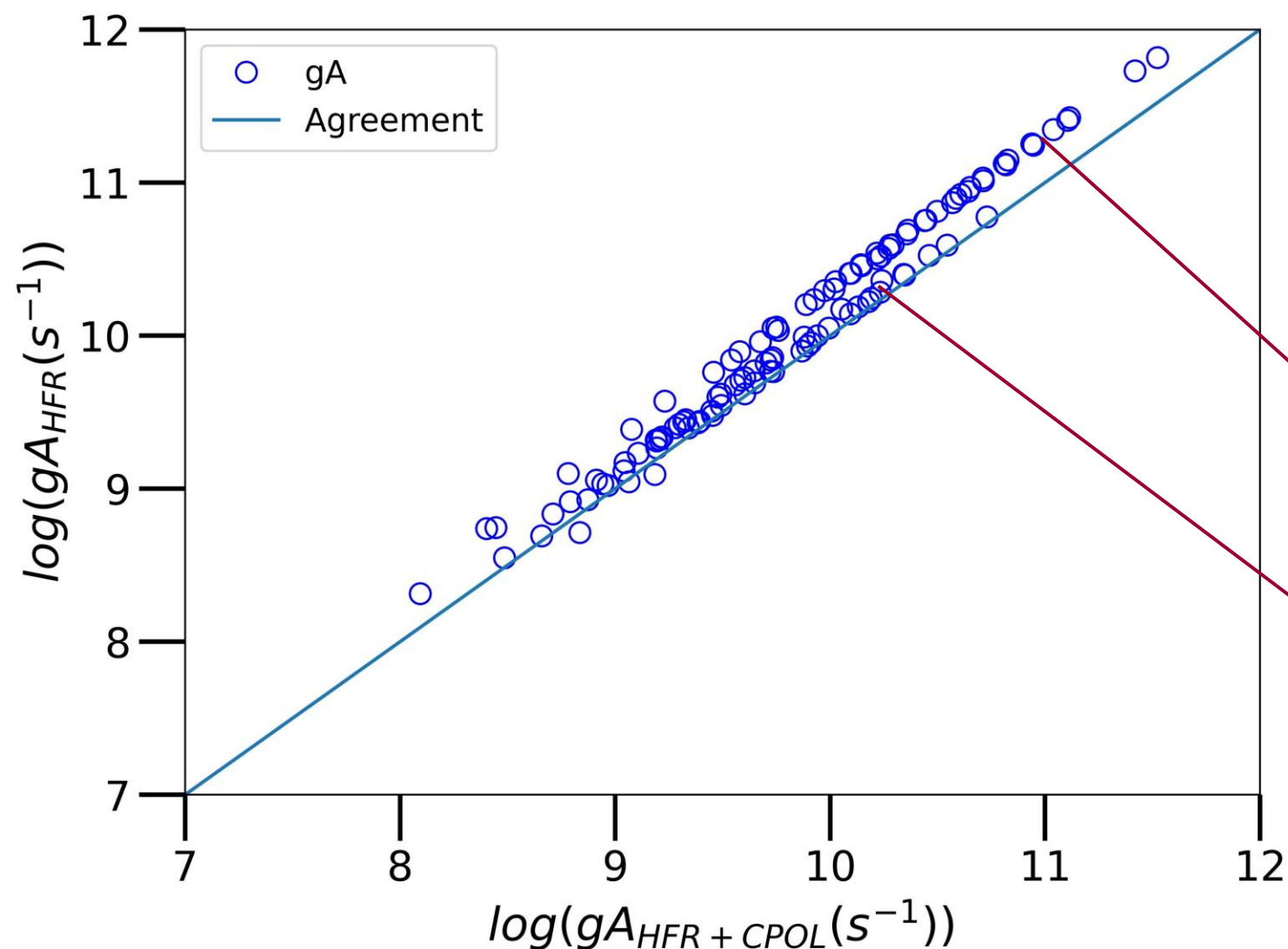
Mean ratio $gA_{HFR+CPOL}/gA_{HFR}=0.83$

CPOL corrections reduce gA-values of
about 20%, as expected.

For 5d → 5f, 7p (5d² → 5d5f, 5d7p) :
reduction is on average 30%

Re VI

112 observed electric dipole lines between 332.044 – 1513.957 Å [Sugar, J. et al, 1994]



Mean ratio $gA_{HFR+CPOL}/gA_{HFR}=0.69$

CPOL corrections reduce more gA-values of about 30%

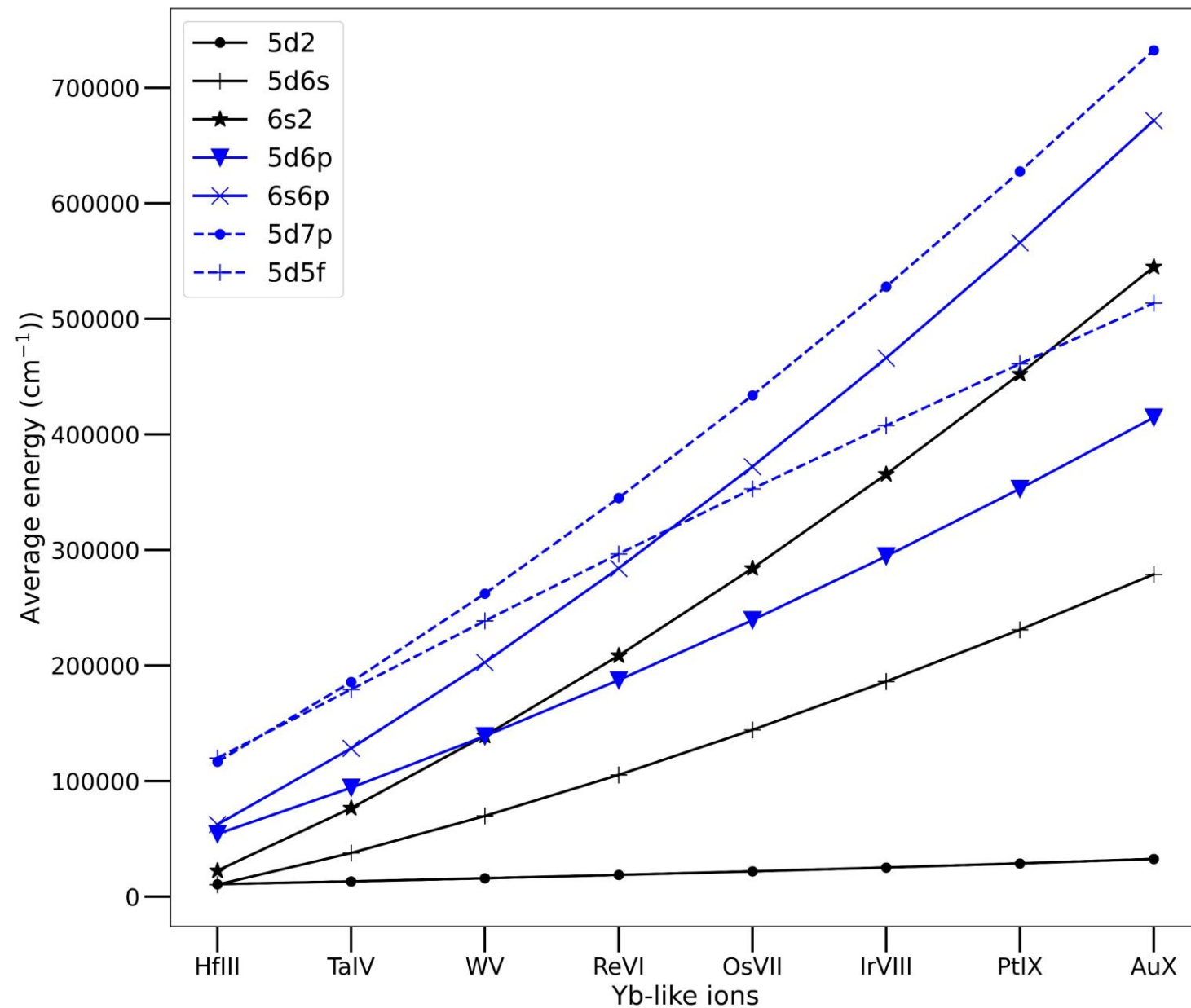
For $5d \rightarrow 5f$ ($5d^2 \rightarrow 5d5f$):
 mean ratio $gA_{HFR+CPOL}/gA_{HFR} = 0.50$

For $5d,6s \rightarrow 6p$ ($5d^2,5d6s \rightarrow 5d6p$):
 mean ratio $gA_{HFR+CPOL}/gA_{HFR} = 0.85$

$\Rightarrow E_{av}$ of configurations with open 4f,5p core orbital closer to valence configurations

Configuration Interactions

The atomic level structure more spread out with increasing ionic charge state → on average less mixing in the level composition

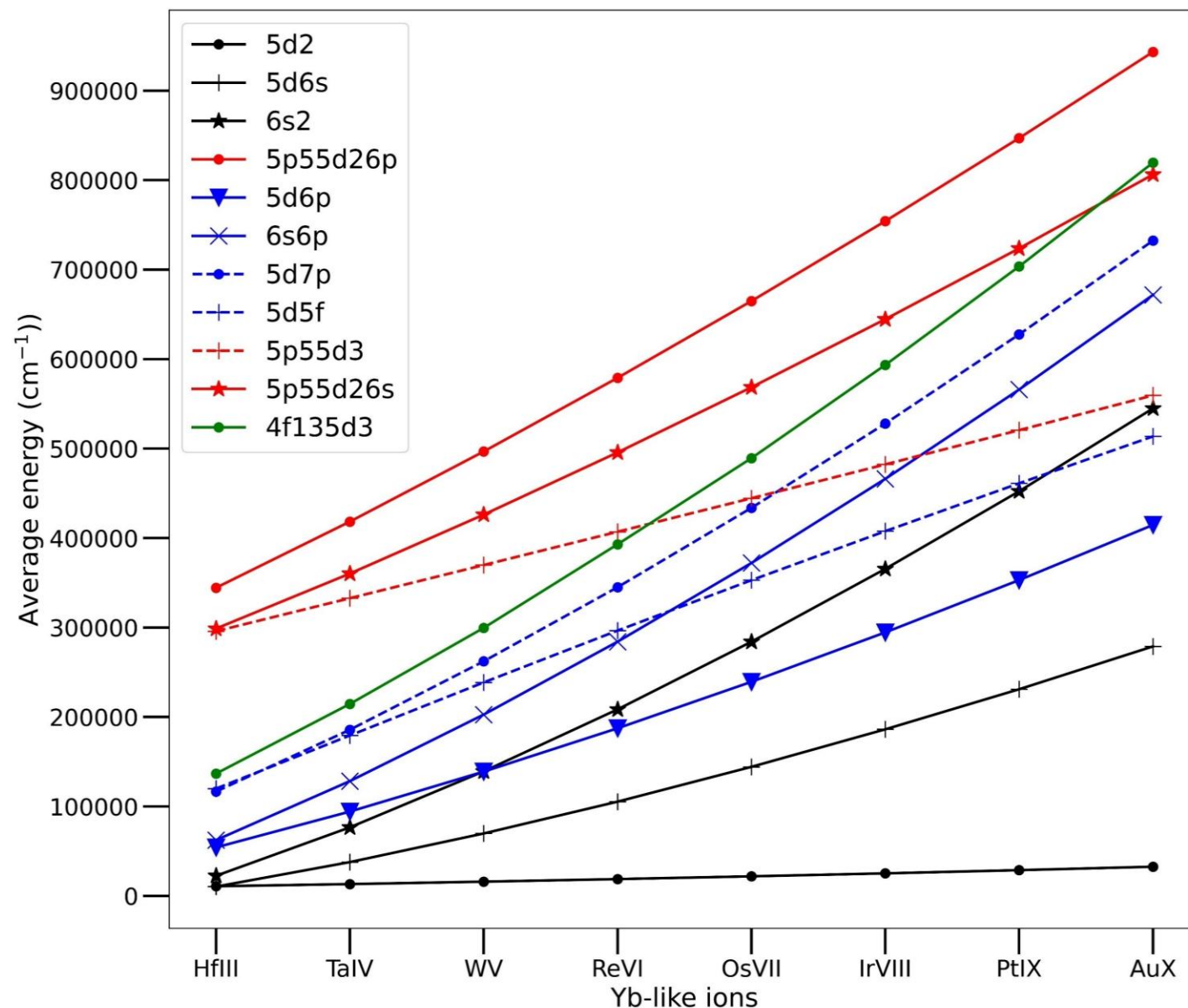


Core-valence correlations

$5p^5 5d^3$ closer to (and overlap) $5d7p$,
 $6s6p$ and $5d5f$

→ core-valence correlations with a
pseudo potential no longer valid

→ need to introduce explicitly the
configurations with an open $5p, 4f$ core
orbitals in our model (through CI)



MCDHF preliminary results

MCDHF method

General procedure [Grant, I. P., 2007]

- $H_{DC} = \sum_{i=1}^N h_{D_i}$ with $h_{D_i} = c\vec{\alpha} \cdot \vec{p}_i + (\beta - 1)c^2 + V(r_i)$ ($\alpha^j = \gamma^0 \gamma^j$ and $\beta = \gamma^0$)
- Each electron: $h_D \varphi = E \varphi \rightarrow \varphi(r, \theta, \phi) = \frac{1}{r} \begin{pmatrix} P_{n,\kappa}(r) \chi_{\kappa,m}(\theta, \phi) \\ i Q_{n,\kappa}(r) \chi_{\kappa,m}(\theta, \phi) \end{pmatrix}$ where $P_{n,\kappa}(r)$ and $Q_{n,\kappa}(r)$ are **large** and **small radial part**, respectively.
- $P_{n,\kappa}(r), Q_{n,\kappa}(r) ? \rightarrow$ solve MCDHF equations (Self-Consistent Field method)
- CI: (ASF) $\Psi(P, J, M) = \sum_{r=1}^{n_c} c_r \Phi(\gamma_r, P, J, M)$

MCDHF Models

- The Atomic State Functions (ASFs) are built with the Active Set approach ($n_{max}l, n'_{max}l', \dots$):
- Optimisation of all orbitals (5s,5p,5d,4f) on the $5d^2\ ^3F_2$ ground state
- MR : (re)optimize valence orbitals on all levels
- VV1,2,3 : ONLY « new » correlation orbitals are optimized on all levels of the MR
- CV,CC : Relativistic Configuration Interactions (RCI) calculations

Even parity		Number of CSFs
MR (5d ² ,5d6s,6s ²)	(6s,5p,5d,4f)	14
VV1	SD(MR) → (7s,6p,6d,5f)	81
VV2	SD(MR) → (8s,7p,7d,6f)	225
VV3	SD(MR) → (9s,8p,8d,7f)	446
CV1 (RCI)	SrD (MR{4f}) → (9s,8p,8d,7f)	29 249
CV2 (RCI)	SrD (MR{4f,5p}) → (9s,8p,8d,7f)	47 273
CV3 (RCI)	SrD (MR{4f,5p,5s}) → (9s,8p,8d,7f)	53 778
CC1 (RCI)	SD (MR{4f(2)}) → (9s,8p,8d,7f)	274 515
CC2 (RCI)	SD (MR{4f,5p}) → (9s,8p,8d,7f)	367 215
CC3 (RCI)	SD (MR{4f,5p}) + SrD (MR{5s}) → (9s,8p,8d,7f)	373 720
Odd parity		Number of CSFs
MR (5d6p,6s6p,5d5f,5d7p)	(6s,7p,5d,5f)	47
VV1	SD(SR) → (7s,8p,6d,6f)	188
VV2	SD(SR) → (8s,9p,7d,7f)	399
VV3	SD(SR) → (9s,10p,8d,8f)	688
CV1 (RCI)	SrD (MR{4f}) → (9s,10p,8d,8f)	115 980
CV2 (RCI)	SrD (MR{4f,5p}) → (9s,10p,8d,8f)	184 664
CV3 (RCI)	SrD (MR{4f,5p,5s}) → (9s,10p,8d,8f)	208 892
CC1 (RCI)	SD (MR{4f(2)}) → (9s,10p,8d,8f)	1 980 405
CC2 (RCI)	SD (MR{4f,5p}) → (9s,10p,8d,8f)	2 574 277
CC3 (RCI)	SD (MR{4f,5p}) + SrD (MR{5s}) → (9s,10p,8d,8f)	2 602 097

Convergence of the models

Ta IV

	Even parity		Odd parity	
	$\Delta E/E_{obs}$	$\Delta E/E_{prev}$	$\Delta E/E_{obs}$	$\Delta E/E_{prev}$
MR	9.99%		1.53%	
VV1	4.12%	5.91%	1.52%	0.35%
VV2	3.75%	0.39%	1.56%	0.05%
VV3	3.69%	0.06%	1.57%	0.01%
CV1	5.25%	1.86%	0.71%	1.44%
CV2	7.71%	3.35%	2.11%	2.08%
CV3	6.40%	1.31%	2.33%	0.21%
CC1	6.59%	1.60%	10.60%	9.48%
CC2	6.90%	0.86%	5.10%	2.83%
CC3	5.65%	0.88%	5.23%	2.74%

303 transitions over 317 have a ratio $gA_{HFR+CPOL}/gA_{MCDHF} < 10$ with a mean ratio 1.39

299 transitions over 317 have a ratio $gA_{HFR+CPOL}/gA_{MCDHF} < 10$ with a mean ratio 1.26

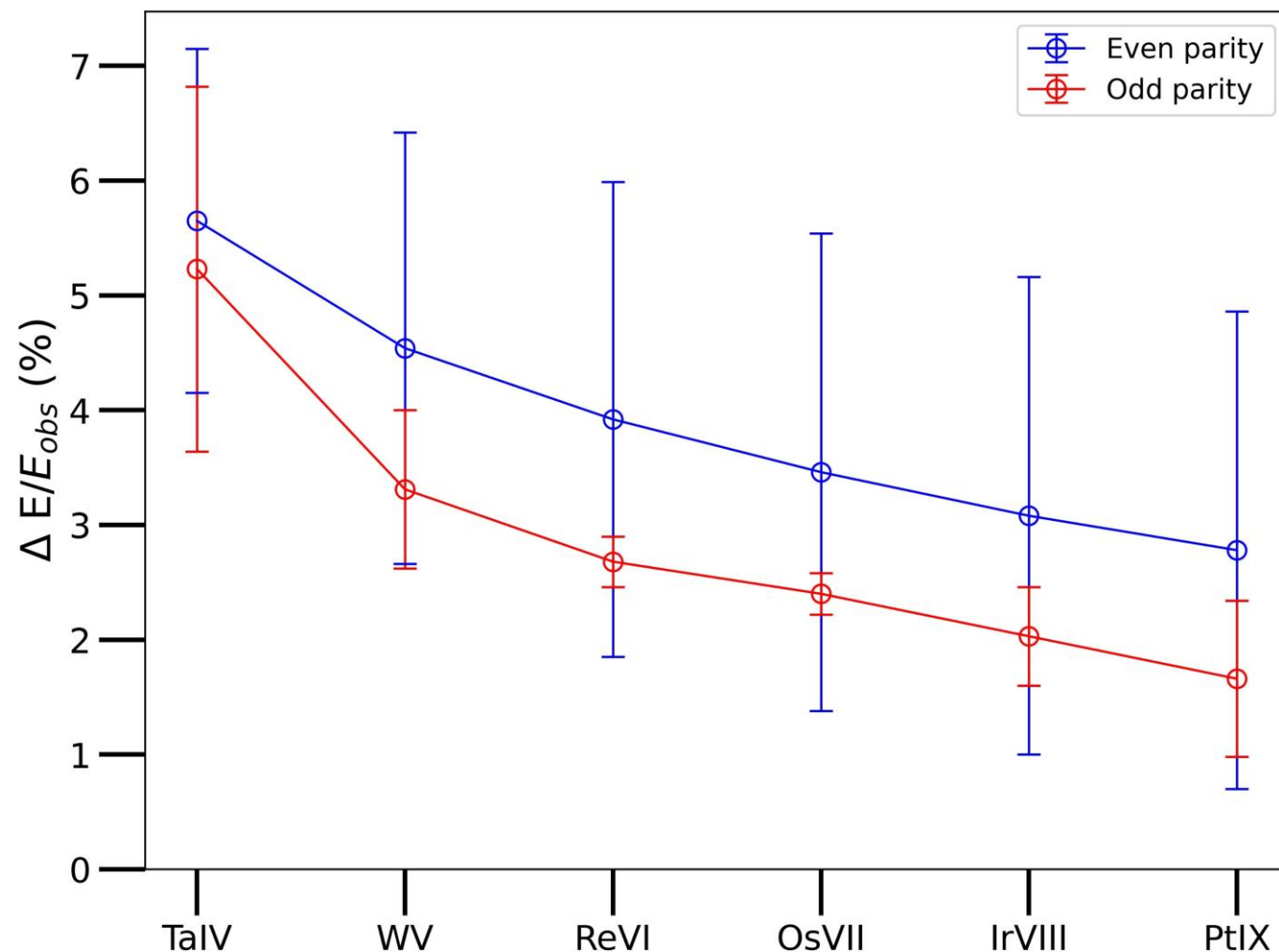
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Odd parity		Number of CSFs
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Even parity (CC3) : SD (MR {4f,5p}) + SrD (MR {5s}) $\rightarrow$ (9s,8p,8d,7f)

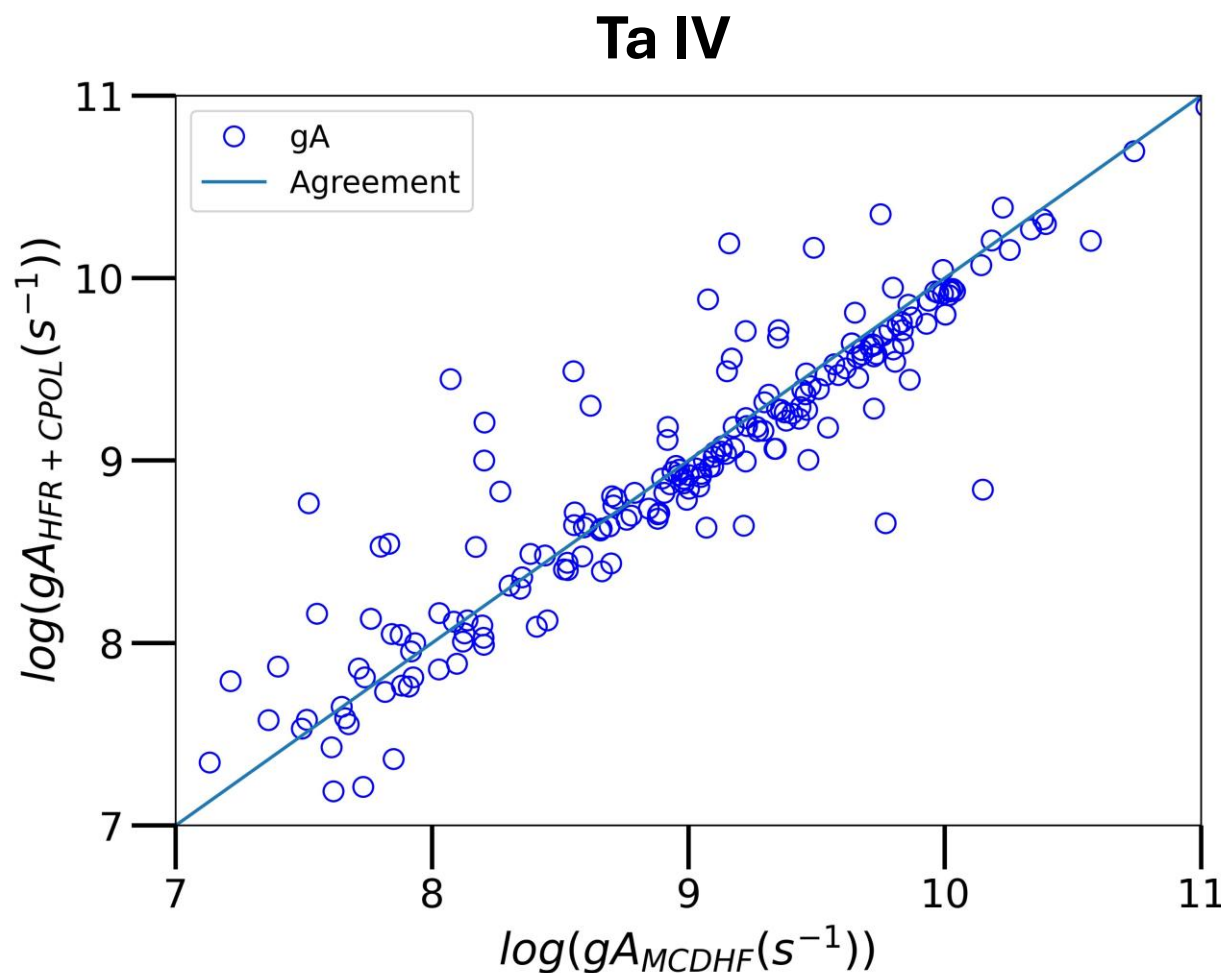
Odd parity (CC3) : SD (MR {4f,5p}) + SrD (MR {5s}) $\rightarrow$ (9s,10p,8d,8f)

	$\Delta E/E_{obs}$	
	Even parity	Odd parity
TaIV	5.65%	5.23%
WV	4.54%	3.31%
ReVI	3.92%	2.68%
OsVII	3.46%	2.40%
IrVIII	3.08%	2.03%
PtIX	2.78%	1.66%

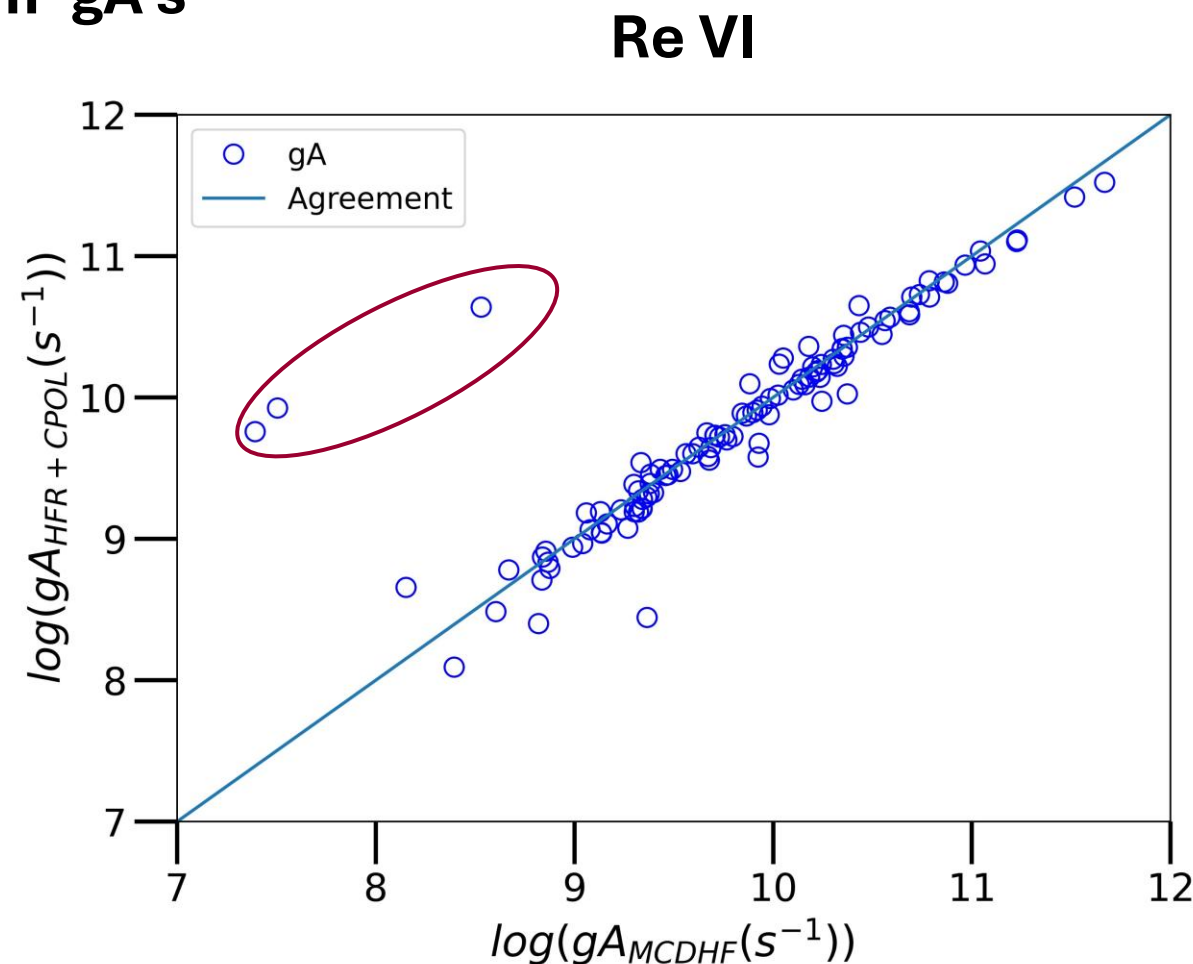
$5d^2 \ ^3P_0, 5d^2 \ ^1S_0 : \Delta E/E_{obs} > 8\%$



Comparisons between HFR+CPOL and MCDHF gA's



Mean relative difference $\Delta gA / \max(gA_{HFR+CPOL}; gA_{MCDHF}) = 32.3\%$
 Mean ratio $gA_{HFR+CPOL} / gA_{MCDHF} = 2.16 \pm 2.30$

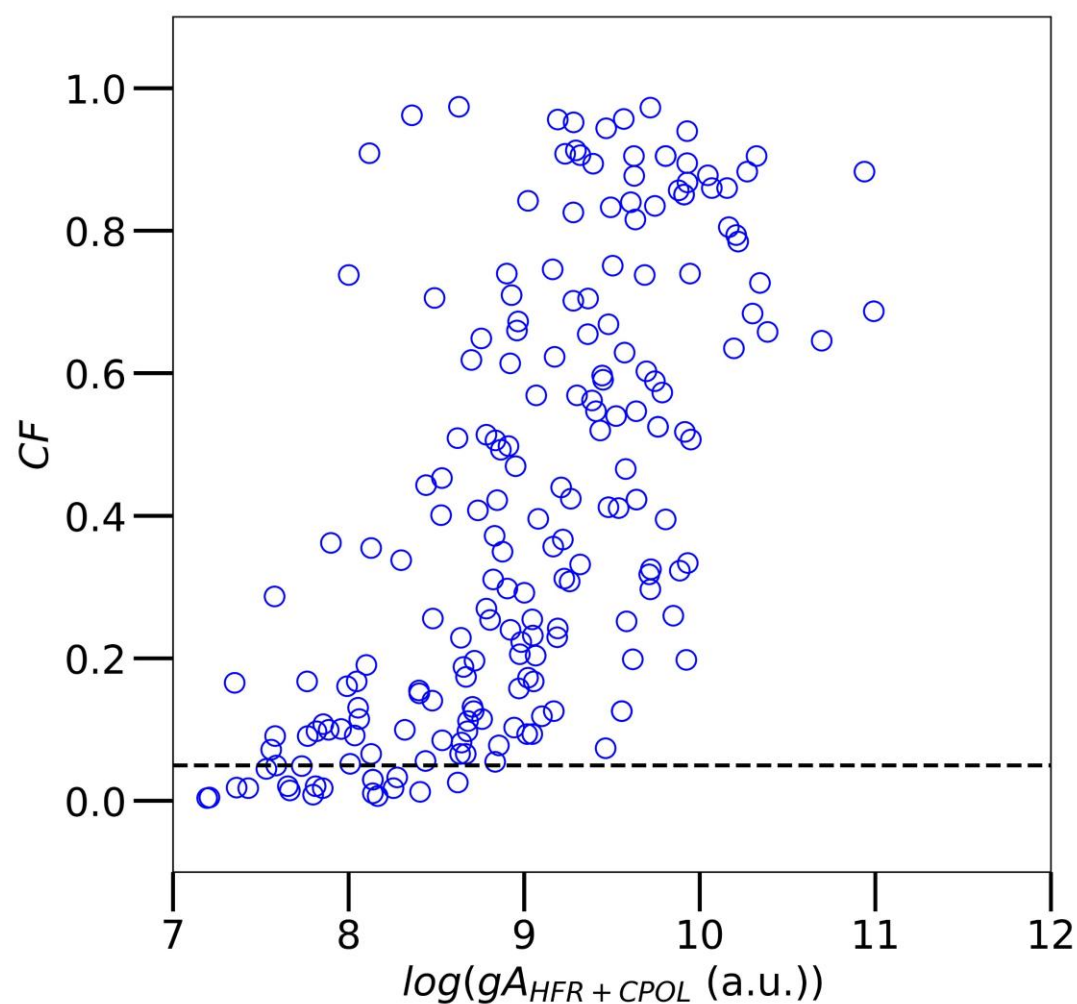


$\Delta gA / \max(gA_{HFR+CPOL}; gA_{MCDHF}) = 18.9\%$
 $gA_{HFR+CPOL} / gA_{MCDHF} = 6.52 \pm 10.81$
 (three high ratios : 129, 232 and 263)

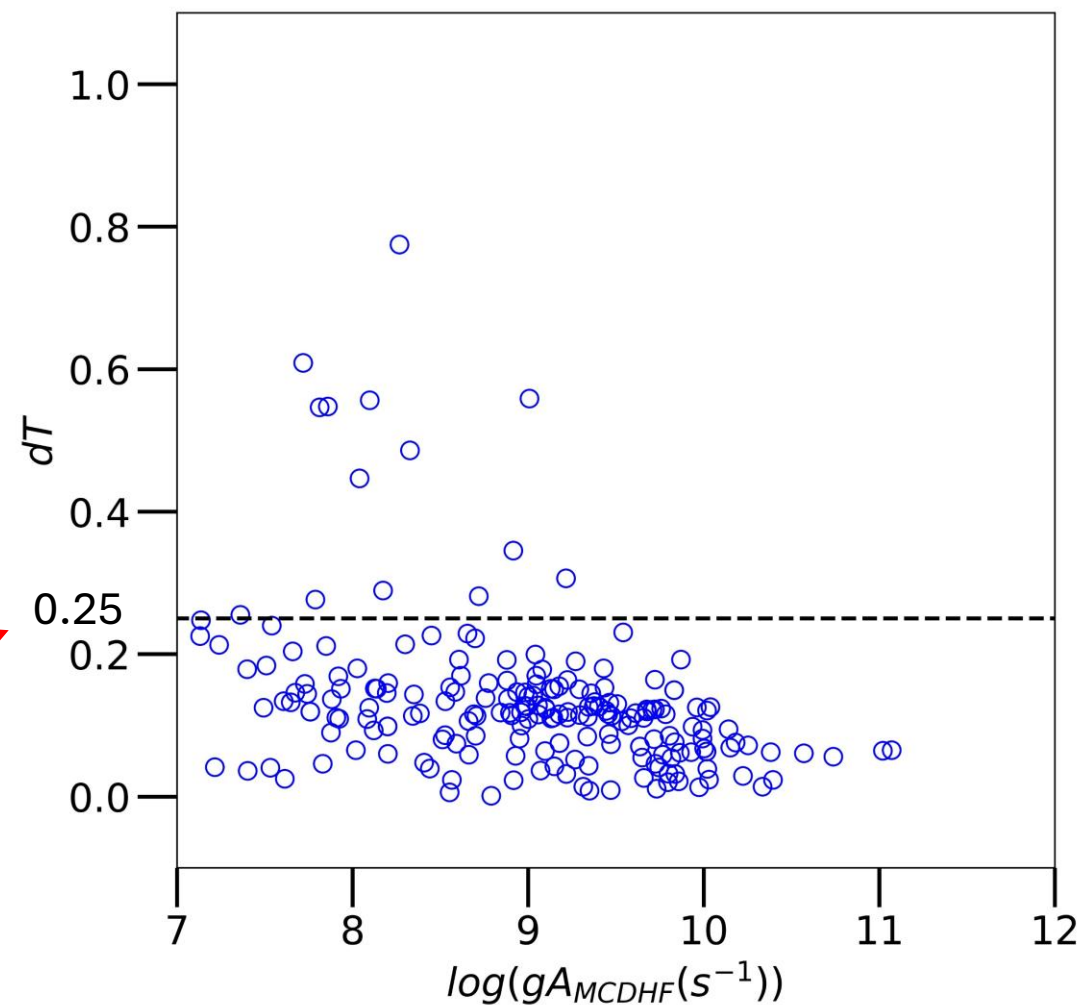
All gA (and gf) computed in both methods are not reliable

Ta IV

Cancellation Factor $CF_{ij} = \left(\frac{|\sum_b \sum_c y_j^b y_i^c \langle \psi_c | \mathbf{P}^{(1)} | \psi_b \rangle|}{\sum_b \sum_c |y_j^b y_i^c \langle \psi_c | \mathbf{P}^{(1)} | \psi_b \rangle|} \right)^2$ [Cowan, R. D., 1981] Uncertainty of MCDHF transition rates : $dT = \frac{|A_B - A_C|}{\max(A_B, A_C)}$



0.05
[Cowan, R. D., 1981]

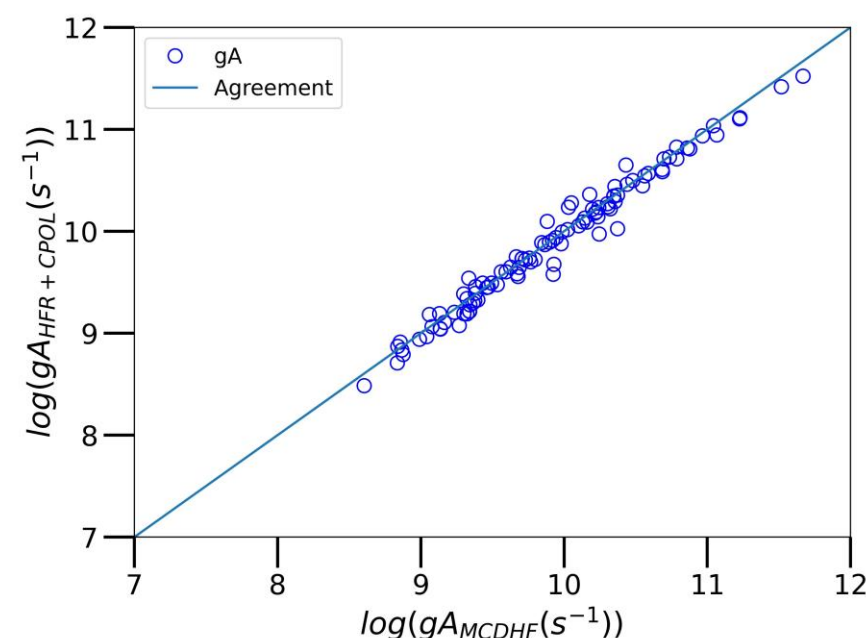
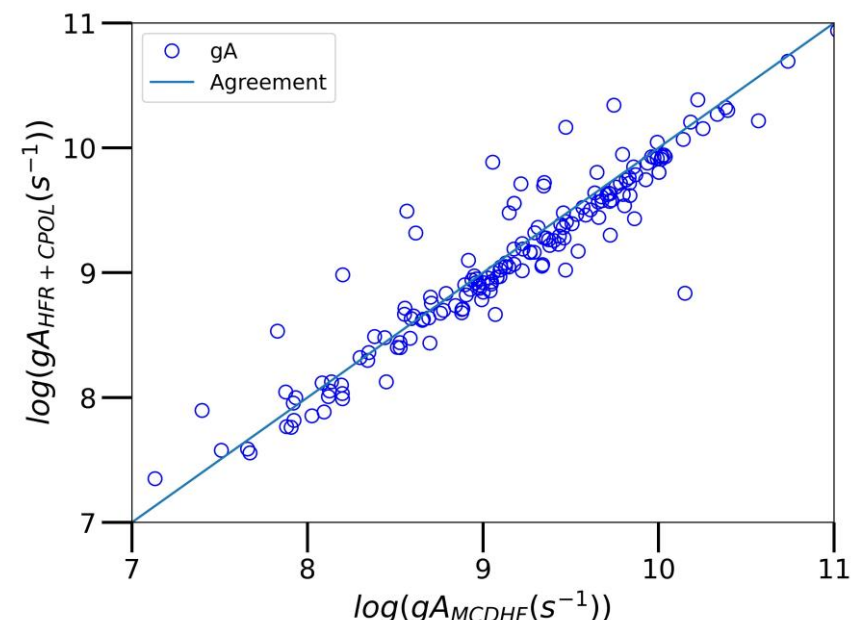


Ta IV

- It remains 170 transitions (among 213) with $CF > 0.05$ and $dT < 0.25$
- The uncertainty on the HFR+CPOL and MCDHF results estimated on average 26%
 $\rightarrow \frac{\Delta gA}{\max}$ where $\Delta gA = |gA_{HFR+CPOL} - gA_{MCDHF}|$
 and $\max = \max(gA_{HFR+CPOL}; gA_{MCDHF})$
- Mean ratio $gA_{HFR+CPOL}/gA_{MCDHF} = 1.12 \pm 0.54$

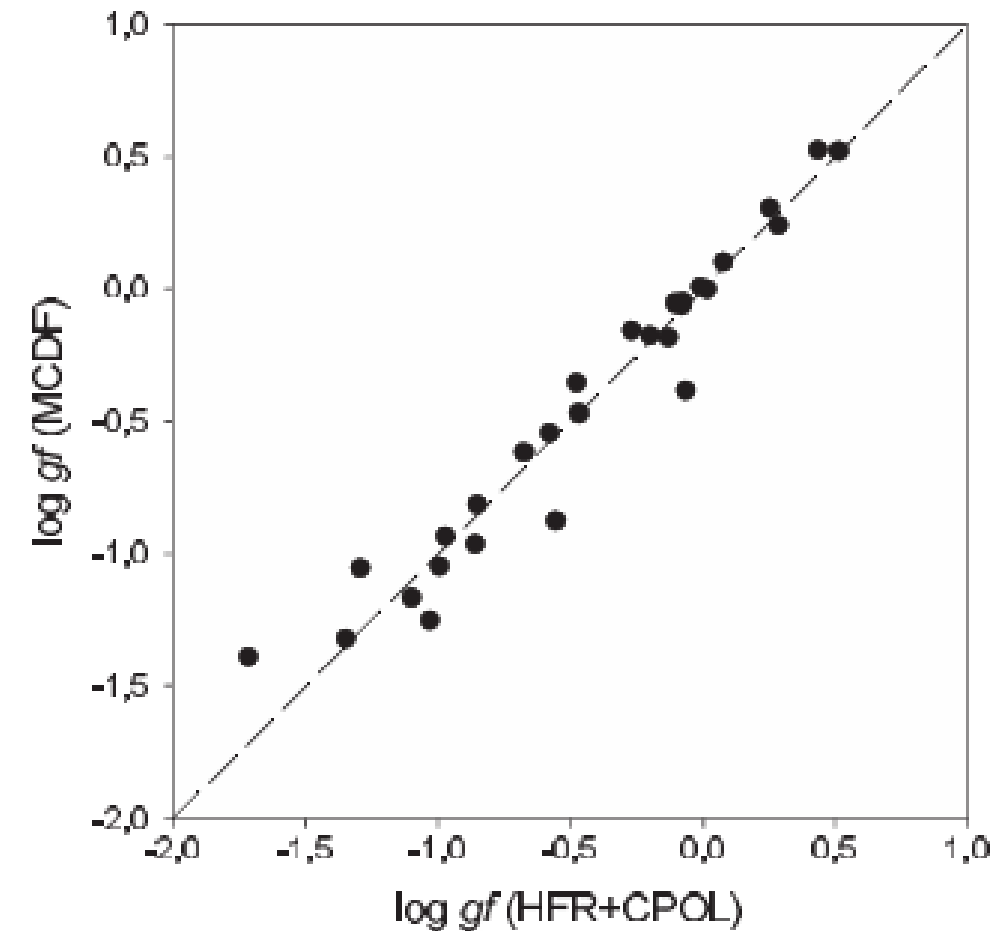
Re VI

- It remains 101 transitions (among 112) with $CF > 0.05$ and $dT < 0.25$
- The uncertainty estimated on average 15% with a mean ratio 0.96 ± 0.16



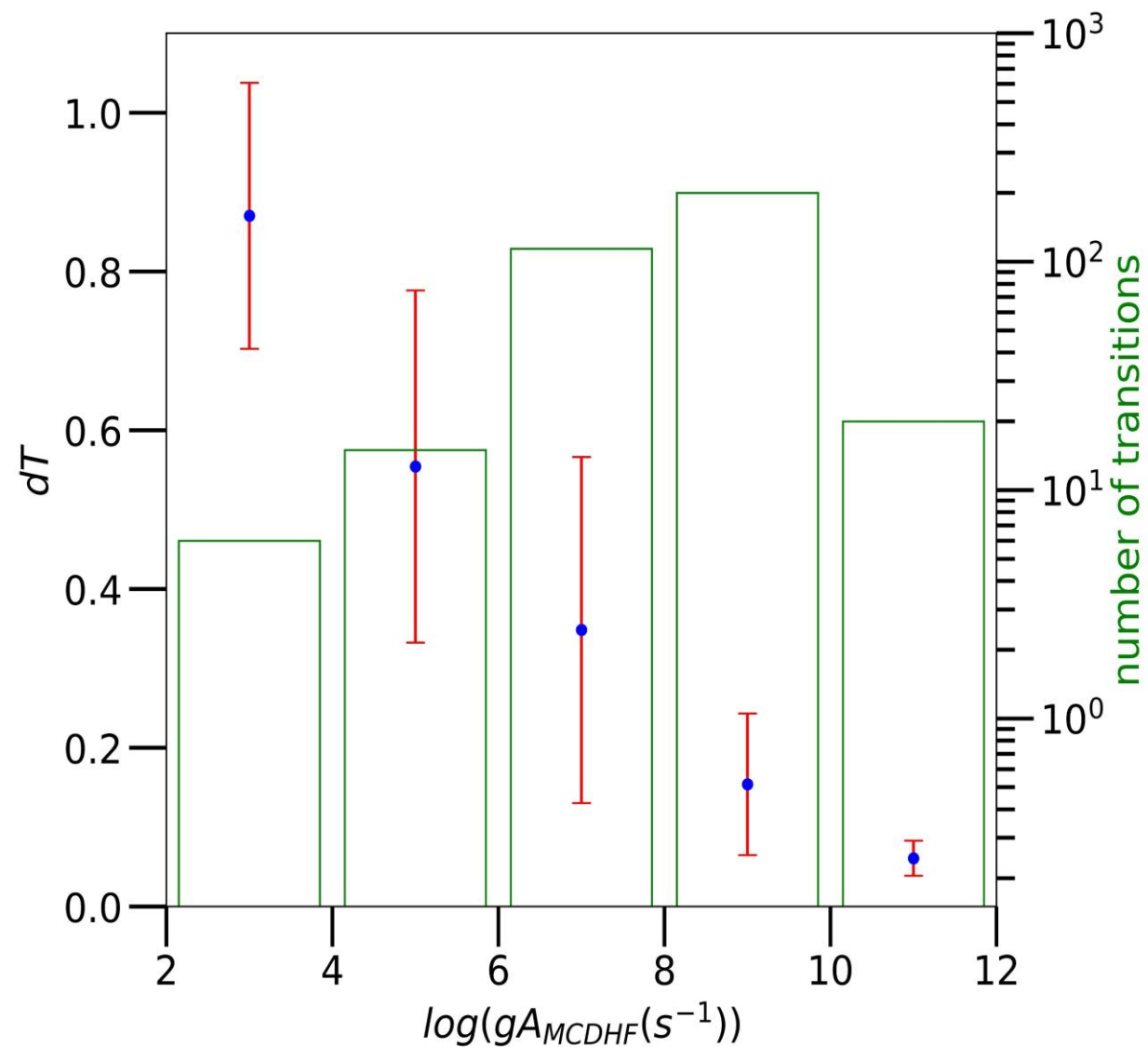
W V

- 193 dipole electric lines between 390 – 2190 Å [Kramida, A. et al, 2024]
- For 5d²-5d6p and 5d6s-5d6p transitions :
- $\Delta gf = |gf_{HFR+CPOL} - gf_{MCDHF}| = 21\% (<30\%)$ [Yoca, S.E et al, 2012]



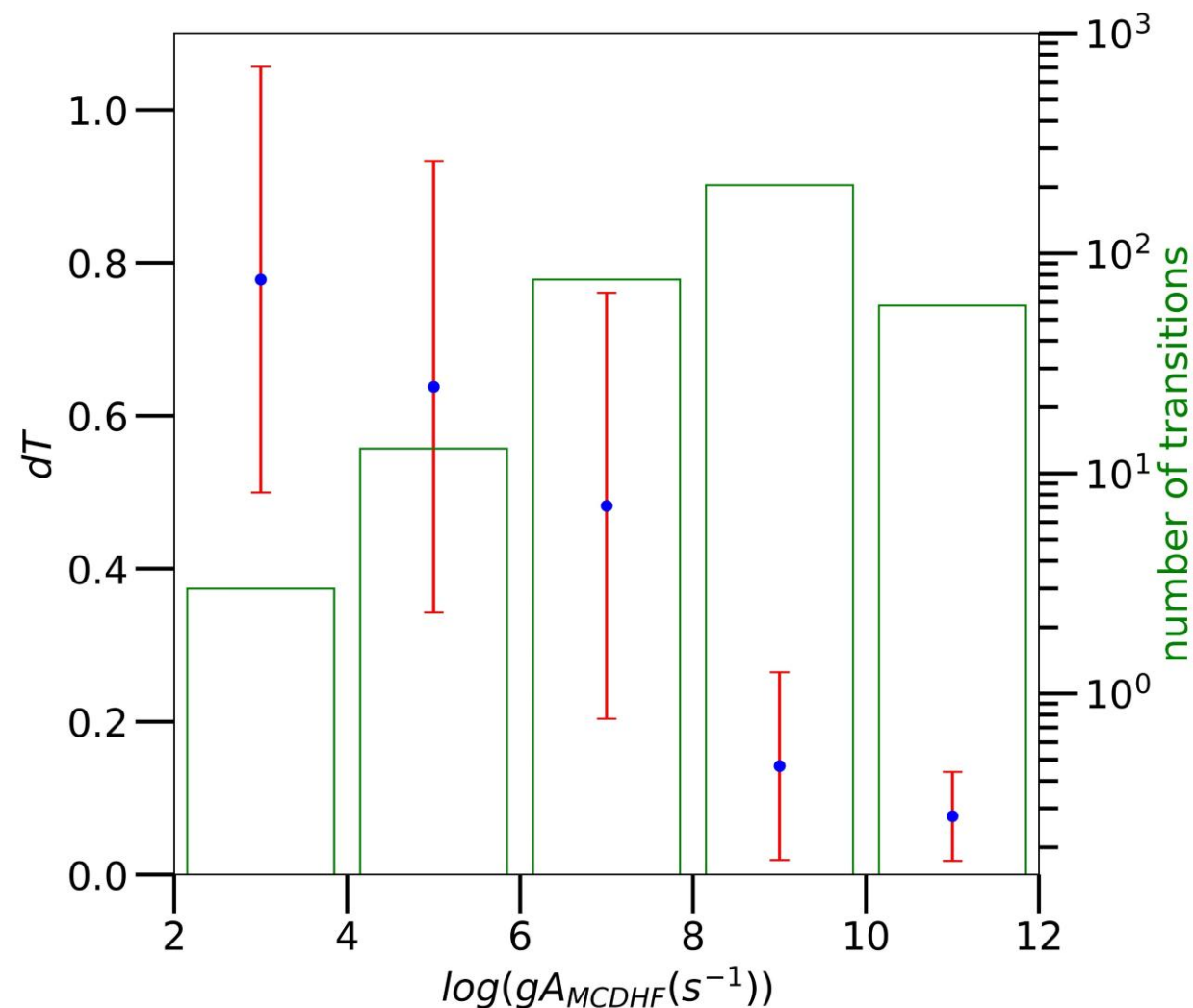
Ta IV

- All 355 transitions computed in the MCDHF method
- ⇒ 251 transitions have $dT < 0.25$
- ⇒ The more intense the transitions, the lower the dT value



Re VI

- All 355 transitions computed in the MCDHF method
- ⇒ 253 transitions have $dT < 0.25$
- ⇒ The more intense the transitions, the lower the dT value



- Origin of Tantalum up to Platinum in the fusion plasma (tungsten transmutation)
- Compute radiative parameters with HFR+CPOL method + least squares adjustment
- Increasing the ionic charge, configurations with an open 4f,5p orbitals begin to overlap the valence configurations → Introduce explicitly these configurations ?!
- Good agreement between HFR+CPOL and MCDHF method and comparisons (170 transitions for TaIV and 101 for ReVI) allow to estimate uncertainty (TaIV:26%; ReVI:15%)
 - set of new atomic data for TaIV → PtIX for plasma diagnostics
- Same procedure for higher ionic charge state of tungsten transmutation products (Ta, Re, Os, Ir, Pt)
- However, few atomic data published

THANK YOU !

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