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THE UNIVERSITY OF MELBOURNE

In memoriam Ian Grant

Chanh Tran, Diamond Light Source, UK Tony Kirk, Minh Dao, Paul di Pascale, Jack Stephens, Ramesh Rijal Truong Nguyen, Jonathan Dean, Hamish Melia, Scott Thompson

Experimental: Zwi Barnea, Martin de Jonge, Stephen Best, Ryan Trevorah, Martin Schalken, Ruwini Ekanayake, Geoff Cousland, Marcus John, Daniel Sier, Nich Tran, Alexis Illig, Mark N Kinnane, Justin A Kimpton, Lucas F Smale, D Paterson, A Payne

- Diamond: Sofia Diaz-Moreno, J Fred W Mosselmans, Matteo Armani, Shusaki, Nitya
- James Hester [ANSTO], Dudley Creagh, Joel Brugger, Barbara Etschmann

Theory: Jay Bourke, Lucas Smale, Chris Witte, Andrew Hayward, John Lowe, Joni Pham, Feng Wang, Finn Jenssens, Paarangat Pushkarna, Yves Joly [CNRS]



The ultimate inner shell transition metal spectra

Copper is the most investigated X-ray source and target, the most studied with X-rays, one of the metals of the ancients. One of the most studied materials in history. Whereas hydrogen has defined and led to quantum mechanics, relativistic quantum mechanics and quantum electro-dynamics, and silicon has been developed as the core of the chip and nanotech industry, copper is one of the most stable and useful of transition metals. So, we should understand it. Really well! So, we should be able to tell a Story of Copper. This presentation will be such a story. However, the ending is complex, the physics is complex, and indeed we do not understand the X-ray spectra of copper yet, nor of any other material less well studied. There is a great and bright future for research and science for copper. We will present past and future challenges with copper and with inner shell X-ray spectra.

My first interactions with **Ian Grant** were during my doctorate with **Josh Silver** 1985-1990. Some decades of this work on copper were with Ian, with **publications in 2009 through 2023 14 publications co-authored with Ian**. With **four research students: A L C Hayward, J A Lowe, T V B Nguyen, T L H Pham**

On five main topics:

Cu Kα, Kβ Ti Kα Oxygen IR, UV Self-energy: Welton Approximation for QED, the LCG method [Lowe Chantler Grant] Convergence of the Breit interaction









The most investigated X-ray source and target, the most studied with X-rays, one of the metals of the ancients. One of the most studied materials in history.

Hydrogen: quantum mechanics, relativistic quantum mechanics and quantum electrodynamics; **silicon:** core of the chip and nanotech industry; **copper:** one of the most stable and useful of transition metals.

The ending is complex, the physics is complex, and we do not understand the X-ray spectra of copper yet, nor of **any other material** less well studied. There is a great and bright future for research and science for copper. We will present past and future challenges with copper and with X-ray spectra.

Copper Z=29, Latin cuprum "from the island of Cyprus." Man's oldest metal, >10,000 years. Copper pendant in northern Iraq dated to 8,700 BC. [Gold, Electrum] High conductivity, thermal and electrical. Coinage metal, tools.

Copper age: end of Neolithic, New Stone Age. Beginning of Metallurgy: in Mesopotamia as copper was used to make molds 4000 BC Brass (Cu, 45% Zinc): 5000 BC Bronze age (Cu, tin Sn, +) : 3500 BC - 1200 BC



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Chemistry: Oxidation state 0 (metals, alloys); +1 (Cuprous oxides etc.); +2 (Cupric oxides, etc.) IR spectra, UV-Vis spectra, NMR [Valence levels, band structure] Visible, optical properties ... [band structure, relativistic QM, d electron eigenvalues]

Physics: Group 11A Note Mendeleev [1869], Moseley [1913] K, L + new elements since Edges: de Broglie [1913] Ground state [Ar]3d¹⁰4s¹ ²S_{1/2} Quantum Mechanics [Heisenberg, Schrödinger, Dirac, Feynman, 1930s]

- Copper anodes and X-ray spectra are well-defined
- The quality of experimental Copper Ka characteristic spectra is unmatched in the literature
- Important calibration tool
- Yet to be adequately theoretically explained!

Copper isotopes, nuclear processes, gamma-spectra ... see later talks from Canada X-ray Diffraction, Neutron Diffraction [Crystal, space group, atomic sites]

- Space group: Fm-3m.
- Space group number: 225.
- Structure: ccp (cubic close-packed)
- Cell parameters: a: 361.49 pm. b: 361.49 pm. c: 361.49 pm. α: 90.000° β: 90.000° γ: 90.000°





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K-edge: 8980.476(20) eV K α_2^0 8027.8416(26) eV K α_1^0 8047.8227(26) eV K β_1^0 8905.413(38) eV LI, LII, LII, M - edges: [PS you might use the little orange book, but the numbers are often incorrect] So what do we know about the Relativistic Quantum Mechanics of Copper, and Why do we care?



1	1.00794 1312.0 2.20 H Hydrogen	2					T	he	Per	iodi	ic To	aple	e of	the	Ele	eme	ents	4.002602 2 Helium
2	6.941 0.98 520.2 0.98 +1 Lithium 19/251	9.012182 899.5 1.57 4 Beryllium 19.29							alkal	ine metals metals ition metals	nonmeta halogen	us ils ises	10.811 5 Boron 197 297 201	12.0107 1086.5 2.55 Carbon 15° 25' 2p ²	14.0067 7 1402.3 3.04 Nitrogen 15º 25º 2pº	15.9994 1313.9 3.44 Oxygen 19 ² 29 ² 29 ⁴	18.998403 9 1681.0 3.98 -1 F Fluorine 19 25* 25*	20.1797 10 2080.7 Neon 16' 26' 26'
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4	39.0983 19 418.8 0.82 19 K Potassium J49 41	40.078 20 589.8 1.00 20 Calcium	44,95591 21 533.1 1.36 21 Scandium [Ar] 3d' 4s ²	47.867 658.8 1.54 Titanium [Ar] 3d ⁷ 4s ²	$\bigvee_{\substack{650.9 \\ 650.9 \\ Vanadium \\ [ke] 3d^9 4s^2}}^{50.9415} 23$	51.9962 24 Chromium Jul 3d ⁹ 4a ¹	54.93804 25 Manganese [Ar] 3d*4s ²	55.845 26 762.5 1.83 26 Fe Iron [Ar] 34* 447	58.93319 27	58.6934 737.1 1.88 28 Nickel Mrj 38 ⁸ 4s ²	$\bigcup_{\substack{\text{Copper}\\ [Ar] 3d^{16} 4a^3}}^{63.546} 29$	^{65.38} 206.4 Zinc (Ar) 3d ¹⁰ 4s ²	69.723 31 578.8 1.81 31 Gallium [ke] 3d ¹⁹ 4a ² 4p ¹	72.64 762.0 2.01 32 Germanium [kr] 3d ¹⁰ 4s ² 4p ²	74.92160 33 947.0 2.18 Arsenic [Ar] 34°6 48° 49°	78.96 2.55 Selenium [4e] 3d ¹⁴ 4s ² 4p ⁴	79.904 1139.9 2.96 Bromine Heij 3d ¹⁹ 44 ² 4p ⁸	83.798 36 1350.8 3.00 36 Krypton Mej 3d ¹⁰ 4p ⁴
5	85.4678 37 403.0 0.82 37 Rubidium JKI 55'	87.62 549.5 0.95 38 Strontium	88.90585 39 600.0 1.22 **********************************	91.224 40 640.1 1.33 40 Zirconium [K] 4d ² 59 ²	92.90638 41 652.1 1.60 Niobium IKI 44* 551	95.96 42 Molybdenum 10140 561	(98) 702.0 1.90 43 Tc Technetium K 4d ⁹ 5s ⁷	101.07 710.2 2.20 Ruthenium [K0] 4d ^r 5s ¹	Rhodium 102.9055 45	Palladium	107.8682 47	112.441 867.8 1.69 48 Cadmium [K] 4d ¹⁰ 55 ⁷	114.818 558.3 1.78 49 Indium Joj 4d° 5s° 5p°	118.710 50 Sn Tin Koj 4d** 5a* 5a*	121.760 51 834.0 2.05 51 Sb Antimony (Ki) 4d ¹⁶ 52 5p ²	127.60 52 Tellurium [Ki] 4d ¹⁰ 5s ² 5p ⁴	126.9044 53	131.293 54 Xenon Kenon
6	132.9054 55 375.7 0.79 Caesium	137.327 502.9 0.89 56 Barium Kei 649	174.9668 71 523.5 1.27 *3 LU Lutefium [Xa) 414 5d1 6s2	178.49 658.5 1.30 Hafnium [Xe] 46 ¹⁴ 5d ² 6d ²	180.9478 73 761.0 1.50 Tantalum Kej 4P* 5dP 6d2	183.84 770.0 2.36 Tungsten (%) 414 544 642	186.207 760.0 1.90 75 Ree Rhenium	190.23 840.0 2.20 Osmium [Xe] 41** 54* 54*	192.217 77 880.0 2.20 77 Iridium Iridium Jej 46* 56* 69*	195.084 870.0 2.28 Pt Platinum (xe) 41* 54* 64	196.9665 79 890.1 2.54 Gold (%) 41** 5d** 6s*	200.59 1007.1 2.00 80 Hercory [Re] 45*5d**6d2	204.3833 81 589.4 1.62 TI Thallium [44] 46* 5d* 5d* 6d* 6d*	207.2 715.6 2.33 82 Pb Leod (%) 41°4 5d1° 65° 65° 65°	208.9804 83 7030 2.02 Bismuth (Xa) 414 546 60 602 602	(210) 812.1 2.00 84 Polonium [xe] ef* 5d* 6e² 6p*	(210) 890.0 2.20 85 Astotine [se] 40* 5d* 6d* 6d* 6d*	(220) 1037.0 86 Radon [He] 4F* 5d* 6s* 6s*
7	(223) 380.0 0.70 87 Francium [Rn] 75 ¹	(226) 509.3 0.90 88 Radium Redium Rej 75?	(262) 103 470.0 ** Lawrencium (Ra) 5P ⁴ 79 ¹ 7p ¹	(261) 104 SB0.0 Rf Rutherfordium (Pn) 5H ⁴ 6d ² 7s ²	(262) 105	(266) 106 Sg Seaborgium	(264) 107 Bh Bohrium	(277) 108 Hassium	(268) 109	(271) 110 Ds Darmstadium	⁽²⁷²⁾ 111 Rg Roentgenium	(285) 112 Copernicium	(284) 113 Uuutrium	⁽²⁸⁹⁾ 114 Uuunguadidm	(288) 115 Uuunpenfium	(292) 116 Uununhexium	117 UUS ^{Ununseptium}	⁽²⁹⁴⁾ 118 Uununactium
$ \frac{1}{10^{100}} = \frac{1}{94.5} + \frac{1}{10^{10}} + \frac{1}{94.5} + \frac{1}{10^{10}} + \frac{1}{5} + \frac{1}{94.5} + \frac{1}{10^{10}} + \frac{1}{1$								⁴ 70 ³⁹ 102 ⁹										

Classic [X-ray] spectra of copper: X-ray Absorption Spectra.

K edge. 1s 8.979 keV L edges. Ll $2s_{1/2}$ Lll $2p_{1/2}$ Llll $2p_{3/2}$ 1096 - 932 eV Soft X-ray regime e.g. Bianconi 1988 Xu 2013 M edges. MI $3s_{1/2}$ MII $3p_{1/2}$ MIII $3p_{3/2}$ 122.5 eV - 75 eV. Soft X-ray or VUV regime. MIV $3d_{3/2}$ MV $3d_{5/2}$ N - conduction band for copper metal $4s_{1/2}$

XAS Compilations Hubbell[1995] [sparse] Chantler theory [1995, 2000]

[double crystal monochromators] [curved crystal, Johann]

Characteristic Spectra

e.g. Moseley 1914 Siegbahn1931 Parratt1936 BremerSorum1979 Sorum1987 Sauder1997 Deutsch+ 1995-2004 Mendenhall+ 2017 Compilations Bearden-Burr[1967] Desclaux[1973] Deslattes [2003] Note C T Chantler, L F Smale, L T Hudson, ITC C, 4th edition, Section 4.2.2 X-ray Energies [2025]

Cu K $\alpha_{2^{0}}$ 8027.8416(26) eV K $\alpha_{1^{0}}$ 8047.8227(26) eV K $\beta_{1^{0}}$ 8905.413(38) eV Cu K $\alpha_{1,2}$ 1s-2p Cu K $\beta_{1,3}$ 1s-3p Cu K $\alpha_{3,4}$ Cu K $\beta_{2,5}$ Cu K α hypersatellite Cu L $\alpha_{1,2}$ 2s-3p Soft X-ray regime Cu L $\beta_{1,3}$ 2s-4p

First X-ray spectra of copper 1913 - 1948+



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Optical and Quantum Electronics 31: 495-505, 1999. © 1999 Kluwer Academic Publishers, Printed in the Netherlands, ELSEVIER

Physics Letters A 286 (2001) 338-346

www.elsevier.com/locate/pla

in atomic physics, resolving discrepancies in

the X-ray regime

C. T. CHANTLER, Z. BARNEA, C. O. TRAN, J. B D. PATERSON School of Physics, University of Melbourne, Parkville, Victoria 3052, Austr X-RAY SPECTROMETRY X-Ray Spectrom. 2000; 29: 449-458

Monitoring fluctuations at a s beamline using matched ion c 1. modelling, data collection, simple measures of association

C. T. Chantler,¹* C. Q. Tran,¹ D. Paterson,¹ Z. Barnea¹ a

Monitoring fluctuations at a sy beamline using matched ion cl 2. isolation of component noise sources, and application to attenuation measurements showing increased precision by two orders of magnitude GAMON Radiation Physics and Chemistry 61 (2001) 347-350

www.elsevier.com/locate/radph

Direct observation of scattering contributions in X-ray attenuation measurements, and evidence for Rayleigh scattering from copper samples rather than thermal-diffuse or Bragg–Laue scattering

C.T. Chantler^{a,*}, C.O. Tran^a, D. Paterson^a, Z. Barnea^a, D.J. Cookson^{b,c}

Precision X-ray optics for fundamental interactions^{X-ray} extended-range technique for precision measurement of the X-ray mass attenuation coefficient and Im(f) for copper using synchrotron radiation

> C.T. Chantler^{a,*}, C.O. Tran^a, D. Paterson^a, D. Cookson^b, Z. Barnea^a PHYSICAL REVIEW A, VOLUME 64, 062506

Measurement of the x-ray mass attenuation coefficient of copper using 8.85–20 keV synchrotron radiation

C. T. Chantler,¹ C. Q. Tran,¹ Z. Barnea,¹ D. Paterson,¹ D. J. Cookson,² and D. X. Balaic¹ ¹School of Physics, University of Melbourne, Victoria 3010, Australia ²ANSTO, Private Mail Bag 1, Menai, NSW 2234, Australia and Chem-Mat-CARS-CAT (Sector 15, Bldg 434D), Argonne National Laboratory, 9700 S. Cass Avenue, Argonne, Illinois 60439 (Received 27 June 2001; published 19 November 2001)

This work presents the x-ray extended range technique for measuring x-ray mass attenuation coefficients. This technique includes the use of multiple foil attenuators at each energy investigated, allowing independent tests of detector linearity and of the harmonic contributions to the monochromated synchrotron beam. Measurements over a wide energy range allow the uncertainty of local foil thickness to be minimized by the calibration of thin sample measurements to those of thick samples. The use of an extended criterion for sample thickness selection allows direct determination of dominant systematics, with an improvement of accuracies compared to previous measurements by up to factors of 20. Resulting accuracies for attenuation coefficients of copper (8.84 to 20 keV) are 0.27–0.5 %, with reproducibility of 0.02%. We also extract the imaginary component of the form factor from the data with the same accuracy. Results are compared to theoretical calcula-





Transmission and Fluorescence X-ray Spectroscopy <u>X-ray Extended Range Technique:</u> Determine [µ/_p] accurately & quantify systematics





NO fitting parameter: <u>accurate</u> experiment and <u>accurate</u> theory <u>CAN</u> agree! FDMX

Question: Is it POSSIBLE or USEFUL to have any consistent experiment & beam-line independent method to 0.2% accuracy?



JL Glover, CT Chantler, Z Barnea, NA Rae, CQ Tran, DC Creagh, D Paterson, BB Dhal, Phys. Rev. A78 (2008) 052902

Copper:





Measurement with Uncertainty



- CQ Tran, CT Chantler, Z Barnea, Physical Review Letts 90 (2003) 257401-1-4
- JD Bourke, CT Chantler, Measurements of Electron Inelastic Mean Free Paths in Materials, Phys. Rev. Letters 104 (2010) 206601-1-4
 Jun 2025 CT Chantler CompAS Lund

Developing theory for XAFS: Copper - plot of $[\mu/\rho]$ vs E



Figure 6: Comparison of experimental (red circles) Cu K-edge XANES [20] with calculations from FEFF90 (green) and FEFF84 (blue). In order to compare the details of the XANES calculation with experiment we have also included additional experimental data (black +) [21, 22], which has been scaled to match absolute measurement.

J. J. Kas, J. J. Rehr, J. L. Glover, C. T. Chantler, Comparison of Theoretical and Experimental Cu and Mo K-edge XAS, NIM A619 (2010) 28-32

Dynamic Bonding from XAFS:

Copper plot of the Fourier transform $\chi(r)$ vs r



Figure 7: Best first shell fits of FEFF8.4 (blue) and FEFF90 (black) compared to experimental (red) EXAFS signal in R space. The fit range was 1.700 - 3.523 Å

J. J. Kas, J. J. Rehr, J. L. Glover, C. T. Chantler, Comparison of Theoretical and Experimental Cu and Mo K-edge XAS, NIM A619 (2010) 28-32

The first Encyclopaedia on XAFS: International Tables for Crystallography Volume I: XAS, Eds Chantler Bunker Boccherini 2024 > 500 cites

Measurement of plasmon-coupling



• Established theories appeared to overestimate the IMFP below 120 eV

 JD Bourke, CT Chantler, Measurements of Electron Inelastic Mean Free Paths in Materials, Phys. Rev. Letts 104 (2010) 206601-1-4
Jun 2025 CT Chantler CompAS Lund

MEASUREMENT OF PLASMON-COUPLING

NEW Plasmon coupling theory (2014) – calculate γ using the loss spectrum itself

$$\begin{split} \lambda(E)_N^{-1} &= \frac{\hbar}{a_o \pi E} \int_0^{\frac{E-E_F}{\hbar}} \int_{q_-}^{q_+} \int_0^{\infty} \frac{2}{\pi} \frac{\omega'}{q} \\ &\times \operatorname{Im} \left[\frac{-1}{\epsilon_{\text{data}}(0,\omega')} \right] \qquad \qquad \gamma_i(q)_N = \hbar \left. \frac{d\omega_q}{dq} \right|_{\omega_q,q} \lambda(E)_N^{-1} \Theta(N-\delta) \\ &\times \operatorname{Im} \left[\frac{-1}{\epsilon_M(q,\omega,\gamma_i(q)_{N-1};\omega_p = \omega_i)} \right] d\omega' dq d\omega \end{split}$$

- First physical, uniquely constrained optical data model since the Penn algorithm [1987]
- Self consistent from successive iterations broadening comes from coupling between excitation channels

Measurement of plasmon-coupling

- Broadening contributes a substantial reduction in IMFP from a fully lossless Lindhard model
- The reduction has a strong impact from excitations below the plasma frequency
 Image: A strong impact from excitations
- Agreement with experiment is improved greatly across all energies
- New predictions for low energy electron transport in any matter. Major differences below 200 - 300 eV, LEED, EELS, Monte Carlo, detector design



Figure 4. Electron inelastic mean free paths of copper calculated using theoretical optical loss data (solid blue curve) and measured optical loss data from Hagemann et al. (dashed red curve) compared with recent measurements using high-accuracy XAFS spectroscopy (black).⁷ Also shown is a result inclusive of plasmon broadening quantified via a previous analysis of copper IMFPs (dotted blue curve),¹⁴ along with a maximum-variational confidence interval.

Bourke, Chantler J Phys Chem A 118 (2014) 909

C T Chantler, J D Bourke, Low-energy electron properties: Electron inelastic mean free path, energy loss function and the dielectric function. Recent measurements the plasmon-coupling theory. Ultramicroscopy 201 Mar (2019) 38-48

[C Q Tran et al.] the new field of Complex Atomic Fine Structure measurement





Jun 2025 CT Chantler CompAS Lund

1. X-ray Absorption and Fluorescence Spectra. Experiment.

- 2. X-ray Absorption and Fluorescence Spectra. Theory.
- 3. X-ray Characteristic Spectra. Experiment.
- 4. X-ray Characteristic Spectra. Theory.

First characteristic spectra of copper [X-ray] [K, L] 1914 - 1948+







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Why study Characteristic radiation?

- Characteristic radiation provides a monochromatic, high intensity source of known energy, used in:
- Astrophysics e.g. for identifying the processes of star formation and properties of black holes. (Winkler, *The Astro. Journal* 1981), (Nguyen et al., *MNRAS* 2014), (Wilkins et al., *Nature* 2021)
- Mining Industry Monitoring and process of metal ores (Tickner et al. 2015, *Rad. Phys. Chem.* 2015)
- Environmental applications measuring microplastics and marine pollutants (Turner et al., *Talanta* 2016), sorting of electronic waste (Bonifazi et al., *Recycling* 2021).
- Medical research advancement of medical radioisotopes (Lee et al. Comp. & Math. Meth. in Med. 2012)
- O Calibration standard in X-ray fluorescence and many fields transferable standards for Cu Kα is critical.
- Characteristic radiation is a major international field of research, allowing us to test and understand fundamental physical processes:
 - Shake off process
 - Auger and Coster-Kronig transitions
 - Radiative Auger emission
- Research into Characteristic radiation also allows:
- Accurate test for multi-configuration Dirac-Fock theory. (Pham et al. J. Phys. B 2016)
- High accuracy tests of QED Anomalous He-like Ti spectra (Payne et al. J. Phys. B 2014)
- Determination of peak energies to the ppm level (Dean et al., *J Phys. B* 2019)
- Measuring hole widths, binding energies, fluorescence yields and intensity ratios
- Ab initio calculations of complex, many-body systems (Nguyen et al., Phys. Rev. A 2022)
- Comparison of computational methods and software.





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A Story of Copper CHARACTERISTIC SPECTRA: "STMPLF" HE UNIVERSITY OF $i\frac{\partial}{\partial t}\Psi = \sum_{i}^{N} \left(\beta c^{2} + c\alpha \cdot p\right)\psi + \sum_{i < i}^{N} \left(\frac{1}{|r_{i} - r_{j}|}\right)$ //ELBOURNE $(\overline{\mathcal{O}}, \overline{\mathcal{O}}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO}, \overline{\mathcalO},$ $i \frac{\partial \alpha_i \cdot \alpha_j}{\partial r_i - r_j} \sum_{i=1}^{n} \frac{\partial \alpha_i \cdot \alpha_j}{\partial r_j}$ $+c^{2}(\alpha_{i}\cdot\nabla_{i})(\alpha_{j}\cdot\nabla_{j})\frac{\cos\left(\frac{|\omega_{i}-\omega_{j}||r_{i}-r_{j}|}{c}\right)-1}{|\omega_{i}-\omega_{j}|^{2}|r_{i}-r_{j}|}$

Calculating many-electron atom wavefunctions MCDHF

Electron correlation \rightarrow multiconfiguration Dirac Hartree Fock method (MCDHF).

 $\text{Relativistic} \rightarrow \text{Dirac equation:} \quad i\hbar \frac{\partial \phi}{\partial t} = \left(c\boldsymbol{\alpha} \cdot \boldsymbol{p} + \beta mc^2 + \frac{Ze^2}{4\pi\epsilon_0 r} \right) \phi \quad \rightarrow \quad \phi_{E\kappa m} = \frac{1}{r} \left| \begin{array}{c} P_{E\kappa}(r)\chi_{\kappa m}(\theta,\varphi) \\ iQ_{E\kappa}(r)\chi_{-\kappa m}(\theta,\varphi) \end{array} \right|$

Our atomic state wavefunction is approximated as a linear combination of configuration state functions (CSFs).

 $\Psi \approx \sum_{j}^{n} c_{j} \phi_{j}$

CSFs are obtained by taking the Slater determinants of excited state electron configurations. eg., the helium atom wavefunction is approximated:

$$\Psi \approx c_1 \phi(1s^2) + c_2 \phi(1s^1 2s^1) + c_3 \phi(2s^2) + c_4 \phi(1s^1 2p^1) + c_5 \phi(2s^1 2p^1) + c_6 \phi(2p^2) \dots$$

The number of possible CSFs is infinite. Therefore we must limit their number by choosing a cut-off subshell.

The subshells for which excitations are allowed (into and out of) form what we call the Active Set. e.g.:

Single configuration: $\Psi(He) \approx c_1 \phi(1s^2)$

Active set up to 2s: $\Psi(He) \approx c_1 \phi(1s^2) + c_2 \phi(1s^1 2s^1) + c_3 \phi(2s^2)$

Active set up to 2p: $\Psi(He) \approx c_1 \phi(1s^2) + c_2 \phi(1s^12s^1) + c_3 \phi(2s^2) + c_4 \phi(1s^12p^1) + c_5 \phi(2s^12p^1) + c_6 \phi(2p^2)$



week ending

18 SEPTEMBER 2009

PHYSICAL REVIEW LETTERS

PRL 103, 123002 (2009)

Theoretical Determination of Characteristic X-Ray Lines and the Copper $K\alpha$ Spectrum

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L.P. Grant

Mathematical Institute, Oxford University, Oxford, United Kingdom (Received 9 May 2009; published 16 September 2009)

Core excitations above the K edge result in $K\alpha$ characteristic x-ray emission. Understanding these spectra is crucial for high accuracies in investigations into QED, near-edge x-ray structure and advanced crystallography. We address unresolved quantitative discrepancies between experiment and theory for copper. These discrepancies arise from an incomplete treatment of electronic interactions. By finding solutions to relativistic multiconfigurational Dirac-Fock equations accounting for correlation and exchange corrections, we obtain an accurate reproduction of the peak energies, excellent agreement of theory with experiment for the line shapes, good convergence between gauges, and account for the $K\alpha$ doublet ratio of 0.522 ± 0.003 :1.



PHYSICAL REVIEW A 82, 052505 (2010)

Multiconfiguration Dirac-Fock calculations in open-shell atoms: Convergence methods and satellite spectra of the copper *K*α photoemission spectrum

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I. P. Grant Mathematical Institute, Oxford University, Oxford, United Kingdom (Received 18 August 2010; published 8 November 2010)

The copper $K\alpha$ photoemission spectra is one of the most widely studied. Recent Dirac-Fock calculations have produced transition energies in good agreement with experiment, though they have relied on approximations that may not be transferable to other complex atoms in which uncertainties in theoretical results are dominated by poor convergence. Through a detailed examination of convergence issues in the copper spectrum, we consider the accuracy obtainable with the multiconfiguration Dirac-Fock (MCDF) method, provide the first determination of fine structure contributions to the spectrum, and demonstrate reliable techniques for modeling spectator states with vacancies in the 3p, 3d, and 4s shells.

PHYSICAL REVIEW A 83, 060501(R) (2011)

Ab initio determination of satellite intensities in transition-metal photoemission spectroscopy using a multiconfiguration framework

J. A. Lowe and C. T. Chantler^{*} School of Physics, University of Melbourne, Australia

I. P. Grant Mathematical Institute, Oxford University, Oxford, United Kingdom (Received 8 March 2011; revised manuscript received 11 June 2011; published 30 June 2011)

Following atomic photoionization, the abrupt change in potential can lead to secondary ionization of an outer-shell electron in a phenomenon known as shake-off, a process which gives rise to the asymmetric $K\alpha$ profile and satellite lines. Investigation of chemical effects and relativistic quantum mechanics requires a theoretical determination of these satellite intensities; however, existing theoretical predictions are inconsistent with experimental results by up to an order of magnitude. Previously theoretical modeling required up to 12 fitting parameters to account for transition widths, energy corrections, spectator intensities, and spectator broadening. Using a multiconfiguration atomic model to account for electron-electron correlation, we provide here the first *ab initio* calculations of shake-off probabilities which are in agreement with experimental results (except for copper), an important step toward a complete theoretical profile.



PHYSICAL REVIEW A 85, 032513 (2012)

Anomalous satellite intensity discrepancy in copper x-ray lines

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I. P. Grant

Mathematical Institute, Oxford University, Oxford, United Kingdom (Received 25 October 2011; published 16 March 2012)

The copper $K\alpha$ spectrum contains asymmetries and satellite features due to secondary electron emission during ionization. Theoretical attempts at determining the intensities of these features are highly discrepant from experimental results. This discrepancy has been the subject of much discussion. In the present work we show that widely applied fitting procedures produce satellite intensities which depend strongly on assumptions regarding parametrization. We also show that recent high-accuracy satellite calculations can provide a copper $K\alpha$ spectrum in good agreement with experiment, thus resolving the discrepancy. This represents a major step toward a complete *ab initio* x-ray spectrum, a goal with significant implications for astrophysics, plasma physics, and tests of quantum electrodynamics.





Copper Ka Photoemission Spectrum



THE UNIVERSITY OF MELBOURNE



A Story of Titanium



Physics Letters A 374 (2010) 4756-4760

A new approach to relativistic multi-configuration quantum mechanics in titanium

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ARTICLE INFO

ABSTRACT

Article history Received 2 August 2010 Received in revised form 30 August 2010 Accepted 27 September 2010 Available online 29 September 2010 Communicated by P.R. Holland

Keywords: Dirac_Fock Atomic structure Photoionization MCDF Diagram lines

Multiply-ionized atoms in plasmas and astronomical systems are predominantly of intermediate atomic numbers with open electron shells. The spectra seen in laboratory plasmas and astrophysical plasmas are dominated by characteristic $K\alpha_{1,2}$ photoemission lines. Modelling these transitions requires advanced relativistic frameworks to begin to formulate solutions. We present a new approach to relativistic multiconfiguration determination of $K\alpha_{1,2}$ diagram and satellite energies in titanium to a high level of convergence, allowing accurate fitting of satellite contributions and the first agreement with profile to negligible residuals. These developments also apply to exciting frontiers including temporal variation of

fundamental constants, theoretical chemistry and laboratory astrophysics. © 2010 Elsevier B.V. All rights reserved.





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J. Phys. B: At. Mol. Opt. Phys. 46 (2013) 015002 (7pp)

High-accuracy reconstruction of titanium x-ray emission spectra, including relative intensities, asymmetry and satellites, and ab initio determination of shake magnitudes for transition metals

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Abstract

High resolution x-ray spectroscopy has revealed a complex structure in the spectrum of core-ionized elements. To date, theoretical reproductions must be fitted to experimental results using fitting parameters to account for transition widths, energy corrections, spectator intensities and spectator broadening-up to 12 or more parameters depending on complexity. We provide here the first accurate reconstruction of the $K\alpha$ spectra in titanium using only instrumental broadening widths as free parameters. We also determine structural systematics in observed shake processes in transition metals for the first time.

A Story of Solar and Astrophysics

THE ASTROPHYSICAL JOURNAL, 769:84 (5pp), 2013 May 20 © 2013. The American Astronomical Society. All rights reserved. Printed in the U.S.A. doi:10.1088/0004-637X/769/1/84

RELATIVISTIC CALCULATION OF TRANSITION PROBABILITIES FOR 557.7 nm AND 297.2 nm EMISSION LINES IN OXYGEN

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ABSTRACT

The 557.7 nm green line and the 297.2 nm ultraviolet line in oxygen have been studied extensively due to their importance in astrophysics and atmospheric science. Despite the enormous effort devoted to these two prominent transition lines over 30 years, and in fact going back to 1934, the ratio of their transition probabilities remains a subject of major discrepancies amongst various theoretical calculations for many decades. Moreover, theoretical results are inconsistent with available laboratory results, as well as recent spacecraft measurements of Earth's airglow. This work presents new relativistic theoretical calculations of the transition probabilities of these two photoemission lines from neutral oxygen using the multi-configuration Dirac–Hartree–Fock method. Our calculations were performed in both length and velocity gauges in order to check for accuracy and consistency, with agreement to 8%. Whilst remaining a challenging computation, these results directly bear upon interpretations of plasma processes and ionization regimes in the universe.

Key words: atmospheric effects – atomic processes – line: identification – methods: analytical – radiation mechanisms: general – techniques: spectroscopic

MNRAS 440, 3439-3443 (2014)

doi:10.1093/mnras/stu511

Advanced ab initio relativistic calculations of transition probabilities for some O I and O III emission lines

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Accepted 2014 March 11. Received 2014 March 7; in original form 2013 January 21

ABSTRACT

This work presents new ab initio relativistic calculations using the multiconfiguration Dirac–Hartree–Fock method of some O I and O III transition lines detected in B-type and Wolf–Rayet stars. Our results are the first able to be presented in both the length and velocity gauges, with excellent gauge convergence. Compared to previous experimental and theoretical uncertainties of up to 50 per cent, our accuracies appear to be in the range of 0.33–5.60 per cent, with gauge convergence up to 0.6 per cent. Similar impressive convergence of the calculated energies is also shown. Two sets of theoretical computations are compared with earlier tabulated measurements. Excellent agreement is obtained with one set of transitions but an interesting and consistent discrepancy exists between the current work and the prior literature, deserving of future experimental studies.





A Story of LCG Welton Self-energy



Radiation Physics and Chemistry 85 (2013) 118-123

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HIGHLIGHTS

▶ We develop a self-energy screening approximation suitable for multi-electron atoms.

- ▶ This approximation is tested in a number of few- and many-electron systems.
- ▶ We obtain superior agreement with experiment compared with existing approximations.
- ► An implementation of this approximation is provided for use with GRASP2K.

ARTICLE INFO

Article history Received 31 October 2012 Accepted 3 January 2013 Available online 11 January 2013 Keywords: OFD Self-energy GRASP Screening Atomic structure

ABSTRACT

Atomic structure calculations have reached levels of accuracy which require evaluation of manyelectron OED contributions. Since exact analytic solutions do not exist, a number of heuristics have been used to approximate the screening of additional electrons. Herein we present an implementation for the widely used GRASP atomic-structure code based on Welton's concept of the electron selfenergy. We show that this implementation provides far superior agreement compared with a range of other theoretical predictions, and that the discrepancy between the present implementation and that previously used is of comparable magnitude to other sources of error in high-accuracy atomic calculations. This improvement is essential for ongoing studies of complex atomic systems. © 2013 Elsevier Ltd. All rights reserved.

Electron self-energy corrections using the Welton concept for atomic structure calculations

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ARTICLE INFO

Keywords:

OED

Self energy

Hydrogenic screening

Atomic structure

ABSTRACT

The high level of accuracy achieved by atomic experiments in recent time has shone a spotlight on the need for a similarly high level of accuracy in atomic structure calculations, and in particular, OED prediction, A method of electron self-energy correction originally derived from the Welton idea by Lowe et al. (2013) (LCG-Welton method) has now been fully incorporated into the popular atomic structural package, GRASP2K, which we have introduced in this paper. A series of benchmark tests and results are presented, which enables the comparison of the implementations of different versions of GRASP2K, and the implementations on different platforms or operating systems. Test results presented in this paper demonstrate that these new implementations maintain the overall consistency and stability of the program across various platforms, while at the same time improve the accuracy of final energies. Our calculations for hydrogenic Ly $\alpha_{1,2}$ transitions show excellent agreement with experiment, to within less than 0.5 eV. On helium-like systems, our calculations show an improvement from the previous GRASP2K screening method. The new results from electronic self-energy contribution using the LCG-Welton method is more consistent with current standards in the literature, where they now fit within experimental variability of up to 0.1 eV. An option for users to adjust the gauge factor in the electric component of the transition rate has also been added to facilitate further investigation of this particular topic.





A Story of Breit in GRASP

PHYSICAL REVIEW A 90, 062504 (2014)

Convergence of the Breit interaction in self-consistent and configuration-interaction approaches

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I. P. Grant Mathematical Institute, Oxford University, Oxford OX1 2JD, United Kingdom (Received 15 October 2014; published 8 December 2014)

Much discussion in relativistic atomic physics and quantum optics has related to the interaction of gauge and perturbation of the Hamiltonian or Dirac operator. It has been commented that Lorentz and gauge independence requires different forms of the perturbation operator in shifting from one gauge to another. Equally, it has been commented that gauge convergence is not possible without different operator forms in different bases and without the operator being embedded within the self-consistent kernel. We explore the logic and self-consistency of these arguments, applied to the well-known Breit operator in an area of continuing discussion. We find that convergence is now possible to a remarkable degree including a Breit interaction operator in a form consistent with the gauge for length and velocity relativistic forms of the multipole operator, implemented at the configuration-interaction level. Excellent convergence is obtained for Breit interaction energies, interaction mixing coefficients, interaction transition probabilities and eigenenergies and transition probabilities in complex open shells (transition metal K α transitions and shake satellites), and forbidden transitions.



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J. Phys. B: At. Mol. Opt. Phys. 49 (2016) 035601 (14pp)

Characterization of the copper $K\beta$ x-ray emission profile: an *ab initio* multiconfiguration Dirac–Hartree–Fock approach with Bayesian constraints

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Abstract

We investigate the $K\beta$ characteristic radiation and the complex asymmetric structure of photoemission lines of copper, which provides a benchmark for theoretical and experimental studies of x-ray calibration series in transition metals. *Ab initio* multi-configuration Dirac– Hartree–Fock (MCDHF) calculations have been performed to study the complex open-shell many-electron problem in copper. The biorthogonalization technique permits determination of transition intensities and Einstein *A* coefficients. The results from our MCDHF calculations demonstrate excellent convergence in transition energies and intensities, as well as gauge invariance to 0.6%. Shake processes caused by single and double spectator vacancies from 3d, 3p, 3s and 4s subshells have also been investigated extensively. MCDHF has been performed to calculate energies and relative intensities of 3d, 3d², 3p, 3s and 4s satellites, resulting in the total number of configuration states exceeding 100 000 and more than 1500 transition components. Our theoretical calculations of shake-off probabilities using the multi-configuration method in the sudden limit have a high degree of internal consistency with the best available experimental data for copper $K\beta$. This supports the validity of relativistic atomic theory and sets a new benchmark even for poorly resolved characteristic spectra using current techniques of analysis.



Hamish A. Melia, Christopher T. Chantler,* Lucas F. Smale and Alexis J. Illig

The characteristic radiation of copper $Ka_{1,2,3,4}$

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A characterization of the Cu $K\alpha_{1,2}$ spectrum is presented, including the 2p satellite line, $K\alpha_{3,4}$, the details of which are robust enough to be transferable to other experiments. This is a step in the renewed attempts to resolve inconsistencies in characteristic X-ray spectra between theory, experiment and alternative experimental geometries. The spectrum was measured using a rotating anode, monolithic Si channel-cut double-crystal monochromator and backgammon detector. Three alternative approaches fitted five Voigt profiles to the data: a residual analysis approach; a peak-by-peak fit; and a simultaneous constrained method. The robustness of the fit is displayed across three spectra obtained with different instrumental broadening. Spectra were not well fitted by transfer of any of three prior characterizations from the literature. Integrated intensities, line widths and centroids are compared with previous empirical fits. The novel experimental setup provides insight into the portability of spectral characterizations of X-ray spectra. From the parameterization, an estimated 3d shake probability of 18% and a 2p shake probability of 0.5% are reported.



Theory of copper $K\alpha$ and $K\beta$ diagram lines, satellite spectra, and *ab initio* determination of single and double shake probabilities



School of Physics, The University of Melbourne, Melbourne Australia

ARTICLE INFO

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Keywords: Copper Kα

ABSTRACT

High-accuracy MCDHF Cu $K\alpha$ and Cu $K\beta$ diagram spectra and major satellites are presented. Spectral eigenvalues reach a theoretical expansion convergence of 0.03 eV or 0.00025%, with gauge and amplitude convergence to 0.7% for Cu $K\alpha$ diagram spectra. Dominant theoretical spectral eigenvalues are shown to be within 0.032 eV of experimental data. *Ab initio* shake-off contributions are presented in terms of total probability from each subshell and as separated single and double shake probabilities. Discrepancies between theory and experiment in copper X-ray spectra arise from profile asymmetries, relativistic terms and shake spectra not yet accounted for. These results and this improved agreement provide key fundamental parameters for processes towards XFEL evolution and resolving key discrepancies between theory and experiment.



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ARTICLE INFO

A B S T R A C T

Keywords: Auger electron spectroscopy X-ray physics X-ray spectroscopy Relativistic quantum mechanics Ab initio theoretical calculations for Auger electron kinetic energies are presented for twenty prominent Auger peaks for the 3d transition metals ($21 \le Z \le 30$). These are the twenty Auger peaks listed for these elemental solids in the National Institute of Standards and Technology (NIST) X-ray photoelectron spectroscopy database. Adding to these values, over one hundred Auger electron kinetic eigenenergies are calculated for tianium. Many of these transition lines are not established in current literature due to their relatively small yields and overlapping widths. These data can be of importance for determination of previously unaccounted Auger electron peaks with titanium and will be useful for X-ray fluorescence studies into the radiative Auger effect. The consistency between our values and the empirical data is an improvement compared with previous binding energy approaches and the methodology is convergent. The methods presented can be extended to other elements for future investigations of Auger electron kinetic energies which is particularly useful where were current experimental values do not exist. Furthermore, these calculations provide evidence in the success of multiconfigurational Dirac-Hartree-Fock approaches in complex quantum mechanics.

PHYSICAL REVIEW A 107, 012809 (2023)

Cu $K\alpha_{3,4}$ satellite spectrum with *ab initio* Auger-rate calculations

H. A. Melia[®], J. W. Dean[®], T. V. B. Nguyen, and C. T. Chantler[®] The University of Melbourne, Parkville, Victoria 3010, Australia

(Received 11 October 2022; revised 7 December 2022; accepted 13 December 2022; published 13 January 2023)

This work investigates the capability of the multiconfiguration Dirac-Hartree-Fock (MCDHF) method in predicting the Cu $K\alpha_{3,4}$ spectrum. *Ab initio* energy eigenvalues, relative intensities, and radiative widths are calculated for the Cu 2p and 2s satellite transitions. By fitting to the most accurate experimental Cu $K\alpha_{3,4}$ spectra available, we show that our 2p satellite energy eigenvalues agree with experiment to within 0.35 eV and that our 2p shake probability agrees with the 2p fitted intensity to within 0.05%. Our fits suggest a I(2s):I(2p) satellite intensity ratio (as a percentage of the total $K\alpha$ spectrum) of 0.03(1):0.76(1) Theoretical predictions of this ratio can be examined using shake probabilities. We calculate the probability ratio of 0.194:0.742. Using MCDHF, the rates at which hole states, created through the shake processes, depopulate via Auger transitions are determined. These results explain the apparent discrepancy between experimental satellite intensities and shake probabilities, and characterize the Cu $K\alpha_{3,4}$ spectrum with a satellite intensity ratio of 0.04(1):0.76, consistent with the experiment





Mendenhall, Marcus H., et al. Journal of Physics B: Atomic, Molecular and Optical Physics 50.11 (2017): 115004. Deutsch, M., et al. Physical Review A 51.1 (1995) 283



• [J W Dean, The Auger effect and its applications. H A Melia, Characteristic Spectral Profiles: Copper Kα and Kβ]

J W Dean C T Chantler, L F Smale, H A Melia, High Accuracy Characterisation for the Absolute Energy of Scandium Kα, J Phys B 52 (2019) 165002 – 1 -12; H A Melia, C T Chantler, L F Smale, A J Illig, The Characteristic Radiation of Copper Kβ including radiative Auger processes, J Phys B 53 (19) (2020) 195002-1-12; J W Dean C T Chantler, L F Smale, H A Melia, An Absolute Energy Measurement of Scandium Kβ to 2 parts per million, J Phys B53 (2020) 205004 – 1 - 10.

- Mendenhall+: Most accurate characteristic spectrum to date; compared with Deutsch, Fritsch+, the best for 20 years ...
- a) Cu Kα_{1,2}: red line Holzer+ 1997, blue line Mendenhall+ 2017 b) Parametrisations around Kα₂. The vertical lines indicate the maximum of each profile characterisation c) Parametrisations around Kα₁. d) Parametrisations of Kα₃₄ Mendenhall+ 2017 (blue); Fritsch+ 1998 (red) e) The difference between the two characterisations (black line) is structured and larger than the s.e. (dark blue) indicating locations of disagreement between the parametrisations





- New experiment can see complex Auger and satellite processes for the first time
- Cu Kβd The origin of the low energy contribution (Kβ') is the least clear feature in the spectrum. The right panel shows 5 Lorentzians including RAS. χ₂r reduced to 1.58 by including RAS, indicating the significance for obtaining transferable parameterisations.

• H A Melia, C T Chantler, L F Smale, A J Illig, The Characteristic Radiation of Copper K β including radiative Auger processes, J Phys B 53 (19) (2020) 195002-1-12

G. Holzer, M. Fritsch, M. Deutsch, J. Hartwig, and E. Forster, Ko1;2 and K β 1;3 x-ray emission lines of the 3d transition metals, Physical Review A56 4554-4568, 1997







Theory of copper $K\alpha$ and $K\beta$ diagram lines, satellite spectra, and *ab initio* determination of single and double shake probabilities



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ABSTRACT

High-accuracy MCDHF Cu $K\alpha$ and Cu $K\beta$ diagram spectra and major satellites are presented. Spectral eigenvalues reach a theoretical expansion convergence of 0.03 eV or 0.00025%, with gauge and amplitude convergence to 0.7% for Cu $K\alpha$ diagram spectra. Dominant theoretical spectral eigenvalues are shown to be within 0.032 eV of experimental data. *Ab initio* shake-off contributions are presented in terms of total probability from each subshell and as separated single and double shake probabilities. Discrepancies between theory and experiment in copper X-ray spectra arise from profile asymmetries, relativistic terms and shake spectra not yet accounted for. These results and this improved agreement provide key fundamental parameters for processes towards XFEL evolution and resolving key discrepancies between theory and experiment.



J.W. Dean^a, C.T. Chantler^{a,*}, B. Ganly^b ^a The University of Melbourne, Parkville, 3010, Australia ^b CSIRO, Glen Osmond, 5064, Australia

ARTICLE INFO

A B S T R A C T

Keywords: Auger electron spectroscopy X-ray physics X-ray spectroscopy Relativistic quantum mechanics Ab initio theoretical calculations for Auger electron kinetic energies are presented for twenty prominent Auger peaks for the 3d transition metals ($21 \le Z \le 30$). These are the twenty Auger peaks listed for these elemental solids in the National Institute of Standards and Technology (NIST) X-ray photoelectron spectroscopy database. Adding to these values, over one hundred Auger electron kinetic eigenenergies are calculated for tianium. Many of these transition lines are not established in current literature due to their relatively small yields and overlapping widths. These data can be of importance for determination of previously unaccounted Auger electron peaks with titanium and will be useful for X-ray fluorescence studies into the radiative Auger effect. The consistency between our values and the empirical data is an improvement compared with previous binding energy approaches and the methodology is convergent. The methods presented can be extended to other elements for future investigations of Auger electron kinetic energies which is particularly useful where were current experimental values do not exist. Furthermore, these calculations provide evidence in the success of multiconfigurational Dirac-Hartree-Fock approaches in complex quantum mechanics.

PHYSICAL REVIEW A 107, 012809 (2023)

Cu $K\alpha_{3,4}$ satellite spectrum with *ab initio* Auger-rate calculations

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(Received 11 October 2022; revised 7 December 2022; accepted 13 December 2022; published 13 January 2023)

This work investigates the capability of the multiconfiguration Dirac-Hartree-Fock (MCDHF) method in predicting the Cu $K\alpha_{3,4}$ spectrum. *Ab initio* energy eigenvalues, relative intensities, and radiative widths are calculated for the Cu 2p and 2s satellite transitions. By fitting to the most accurate experimental Cu $K\alpha_{3,4}$ spectra available, we show that our 2p satellite energy eigenvalues agree with experiment to within 0.35 eV and that our 2p shake probability agrees with the 2p fitted intensity to within 0.05%. Our fits suggest a I(2s):I(2p) satellite intensity ratio (as a percentage of the total $K\alpha$ spectrum) of 0.03(1):0.76(1) Theoretical predictions of this ratio can be examined using shake probabilities. We calculate the probability ratio of 0.194:0.742. Using MCDHF, the rates at which hole states, created through the shake processes, depopulate via Auger transitions are determined. These results explain the apparent discrepancy between experimental satellite intensities and shake probabilities, and characterize the Cu $K\alpha_{3,4}$ spectrum with a satellite intensity ratio of 0.04(1):0.76, consistent with the experiment



A Story of Copper Ab Initio Shake Probability

- The shake probability can be determined for each satellite spectra
- Shake probability is used to scale the satellite intensity
- The copper K α satellite intensity, Nguyen et al.



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T V B Nguyen et al. Phys. Rev. A 105 (2022) p022811

Kα Satellite Spectra







- By including the Auger effect, the model is improved significantly
- The 4s satellite (Orange) is significantly smaller and broader
- This is the best fit for the experimental features

Jun 2025 CT Chantler CompAS Lund

- The most complete theoretical model for the copper $K\alpha_{1,2}$ profile to date
- Importance of the shake processes to modelling characteristic radiation, providing the first evidence of the 3p and 3s satellite spectra
- First **ab initio lifetime calculations** for the upper states of the copper $K\alpha_{1,2}$ transition
- Development of new lifetime convergence metric
- First Auger rate calculations for the copper $K\alpha_{1,2}$ profile
- Strong evidence for the necessity of the Auger Effect in modelling the copper Kα_{1,2} transition through corrections to the FWHM and satellite intensities



THE UNIVERSITY OF



A Story of Copper Characteristic radiation and Cu Ka_{3.4}

Diagram transitions:



Mendenhall, M. H., Henins, A., Hudson, L. T., Szabo, C. I., Windover, D., & Cline, J. P. (2017). High-precision measurement of the x-ray Cu Ka spectrum. JPhysB: Atomic, Molecular and Optical Physics, 50(11), 115004.

Modelling with only the 2*p* satellite

We can perform 3 Fits using the results of our 3 active sets.

Expanded to:	χ^2_r	Shift	Contribution to $K\alpha$ (%)
5s	29.055	1.59(3)	0.919(4)
4s	4.777	-0.34(2)	0.856(3)
4f	2.910	-0.73(2)	0.843(3)
5s	3.051	-0.25(2)	0.791(3)

- We have converged to within 0.25 eV of experimental peak energies (0.003 %).
- As we expand the active set our results improve.
- Residuals between the peak point to new physics





Jun 2025 CT Chantler CompAS Lund

Modelling including non-radiative decay channels

- Non-radiative rates using MCDHF RATIP (Fritzche S. *et al.* (2012) Comp. Phys. Comm.)
- Considering the depopulation of the 2s initial states through nonradiative rates AND the shake probabilities we determine a new *ab initio* satellite intensity ratio: I(2p)/I(2s) = 26



ccuy		5		10 3	
Initial Ho	le Final Hole(s)	Type	Name	Rate (eV/\hbar)	when
[1s]	[2p]	Radiative	$K\alpha$ diagram	0.256	
[1s2s]	[2p2s]	Radiative	$K\alpha \ 2s \ sat.$	0.102	ESCAN
	[1s2p3s]	Auger	$L_1L_{2,3}M_1$	1.342	AUDE
	[1s2p3s]	Auger	$L_1L_{2,3}M_{2,3}$	2.210	'ERSITY O
	[1s2p3d]	Auger	$L_1L_{2,3}M_{4,5}$	3.021	
	[1s2p4s]	Auger	$\mathrm{L_{1}L_{2,3}N_{1}}$	0.741	JUKNI
	[1s3s3s]	Auger	$L_1M_1M_1$	0.012	Contraction in
	[1s3s3p]	Auger	$L_1M_1M_{2,3} \\$	0.142	SCALES ?
	[1s3s3d]	Auger	$L_1M_1M_{4,5}$	0.547	1 A.M.
	[1s3p3p]	Auger	$L_1M_{2,3}M_{2,3}$	1.561	
	[1s3p3d]	Auger	$L_1M_{2,3}M_{4,5}$	2.833	
	[1s3d3d]	Auger	$L_1M_{4,5}M_{4,5}$	1.695	ALC: NOT
	$[1s3\ell 4s]$	Auger	$\Sigma \ \mathrm{LMN}_1$	< 0.01	
]	$\begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \end{array} \end{array} \begin{array}{c} \begin{array}{c} \begin{array}{c} \end{array} \end{array} \begin{array}{c} \begin{array}{c} \begin{array}{c} \end{array} \end{array} \end{array} \begin{array}{c} \begin{array}{c} \end{array} \end{array} \end{array} \left(\begin{array}{c} \end{array} \end{array} \left(\begin{array}{c} \end{array} \right) \end{array} \left(\end{array} \right) \end{array} \left(\begin{array}{c} \end{array} \right) \end{array} \left(\end{array} \right) \end{array} \left(\end{array} \right) \end{array} \left(\begin{array}{c} \end{array} \right) \end{array} \left(\end{array} \right) \end{array} \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \left(\end{array} \right) \left(\end{array} \left) \left(\end{array} \right) \left(\end{array} \left) \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \left) \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \right) \left(\end{array} \left) \left(\end{array} \left) \left(\end{array} \right) \left(\end{array} \left) \left(\end{array} \left) \left(\end{array} \left) \left(\end{array} \left) \left(\end{array} \right) \left(\end{array} \left) \left(\end{array} $	ger rate: 1	4.104		
[1s2p]	[2p2p]	Radiative	$K\alpha \ 2p \text{ sat.}$	0.257	
	[1s3s3s]	Auger	$\mathrm{L}_{2,3}\mathrm{M}_{1}\mathrm{M}_{1}$	0.019	Mar Alexand
	[1s3s3p]	Auger	$L_{2,3}M_1M_{2,3}$	0.119	
	[1s3s3d]	Auger	$L_{2,3}M_1M_{4,5}$	0.508	
	[1s3p3p]	Auger	$L_{2,3}M_{2,3}M_{2,3}$	1.342	
	[1s3p3d]	Auger	$L_{2,3}M_{2,3}M_{4,5}$	1.240	
	[1s3d3d]	Auger	$L_{2,3}M_{4,5}M_{4,5}$	1.891	
	$[1s3\ell 4s]$	Auger	$\Sigma \text{ LMN}_1$	< 0.01	
	2p total Auge	r rate: 5.2	99		

A new process ... not 1s - 2p in the characteristic spectrum



• New theory can now see the 100-1000 spectral components of characteristic X-ray radiation for X-ray spectroscopy and fundamental processes

- Multiconfiguration Dirac-Hartree-Fock (MCDHF). Diagram spectra expanded to 5s with simultaneous convergence of 28000 configuration state functions (CSFs), K α , and to 6g with simultaneous convergence of 91000 CSFs, K β , eigenvalue convergence to ±0.03 eV or 0.00025%, 10x improved upon past work.
- Biorthogonalisation, developments of the active space approach, analysis of markers for theoretical convergence of eigenvalues, and the question of self-consistency for K α and K β .
- Gauge convergence, eigenvalue convergence, A-coefficient convergence. Without the satellite spectra it is not possible to make use of the increased accuracy of the diagram computations.
- Cu Kα 3d⁸ double-shake satellite spectrum: 1506 unique eigenvalues (transitions); simultaneous convergence of 593 000 CSFs



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