



Searches for top squark pair production with the CMS detector

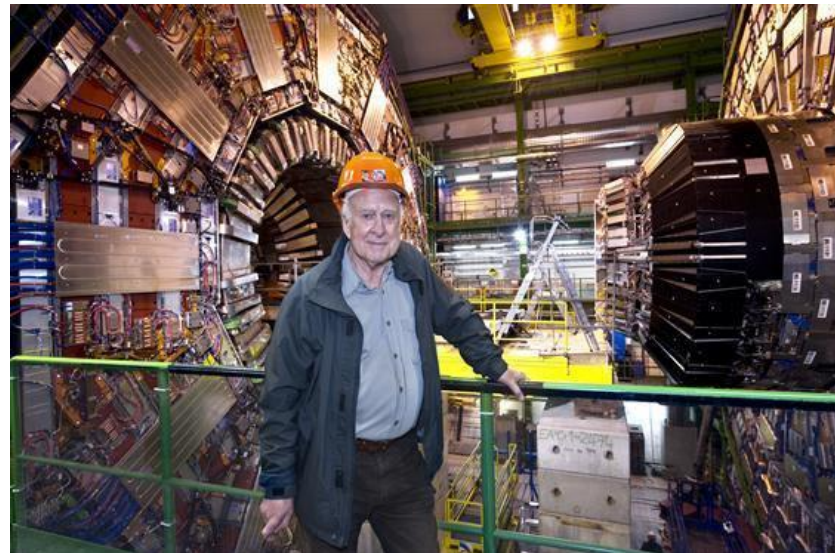
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CERN

November 15, 2016

- Why would we look for top squarks?
- How would top squarks decay and how can we detect them?
- Looking for top squarks in the single-lepton final state
- Peek at difficult mass spectra
- Conclusions and outlook

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# Motivation

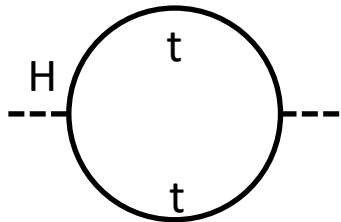


- In 2012 the Higgs was seen in the ATLAS and CMS detectors
  - Gives mass to other particles
- Led to detailed measurements of its properties, e.g.  $m_H \approx 125 \text{ GeV}$
- Irony: the mass of the Higgs not understood

# Corrections to Higgs mass

- Quantum corrections to the Higgs boson mass become uncomfortably large in the Standard Model (SM)
  - Due to top quark loops

$$\Delta M_H^2 \sim |\lambda_t|^2 \left[ -\Lambda_{UV}^2 + \frac{3}{2} M_t^2 \log\left(\frac{\Lambda_{UV}^2}{M_t^2}\right) \right]$$



$\Lambda_{UV}$  is SM cut-off,  
without new physics at  
Planck scale ( $10^{19}$  GeV)

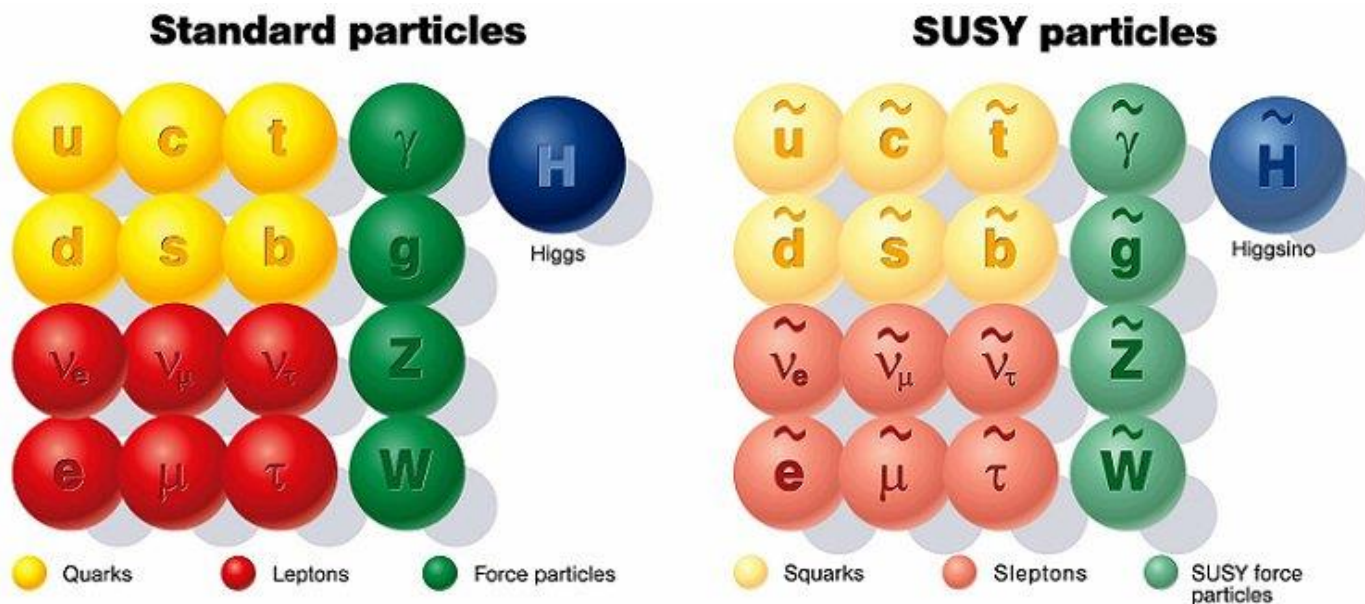
- $M_H^2 = M_{H,0}^2 + \Delta M_H^2$   
 $O(10^4) = O(10^{38}) - O(10^{38})$

– **Cancellation of  $10^{-34}$ ? Highly unlikely!**

- Alternative: new symmetry (supersymmetry)

# Supersymmetry (SUSY)

- New fundamental symmetry relating bosons with fermions. Thus, each SM particle has a SUSY partner
- Imposing R-parity conservation:
  - SUSY particles produced in pairs
  - Lightest SUSY particle (LSP) is stable neutral particle
    - **WIMP dark matter candidate!**



# Solving the problem

- Quantum corrections to the Higgs boson mass become uncomfortably large in the Standard Model (SM)
  - Adding top squarks

$$\Delta M_H^2 \sim |\lambda_t|^2 \left[ -\cancel{\Lambda_{UV}^2} + \frac{3}{2} M_t^2 \log\left(\frac{\Lambda_{UV}^2}{M_t^2}\right) + \cancel{\Lambda_{UV}^2} - \frac{3}{2} M_{\tilde{t}}^2 \log\left(\frac{\Lambda_{UV}^2}{M_{\tilde{t}}^2}\right) + \dots \right]$$

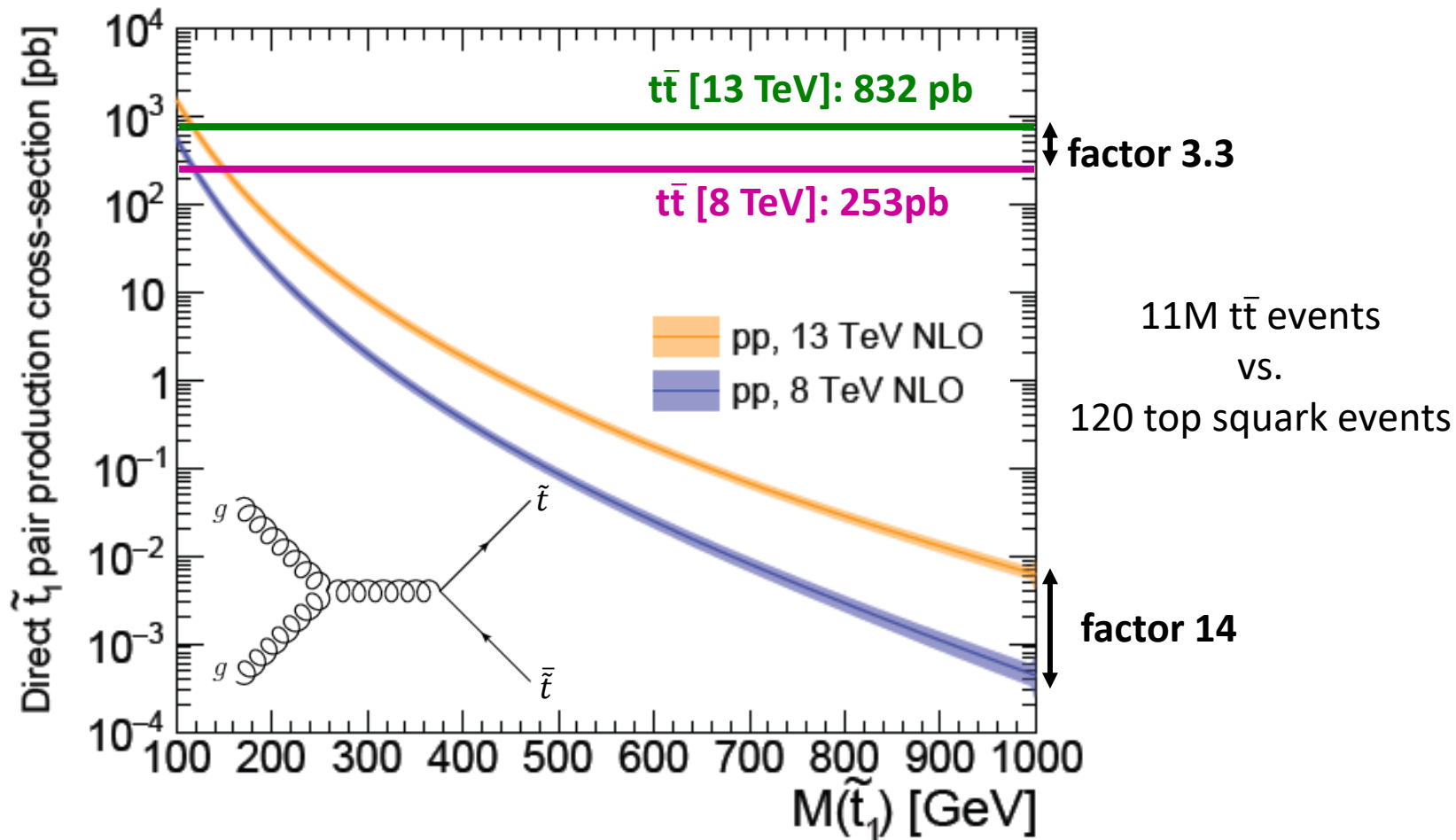
- Top squarks cancel the quadratic divergence**
- Left-over terms proportional to  $M_t^2 - M_{\tilde{t}}^2$ 
  - Light stops make a more “natural” solution

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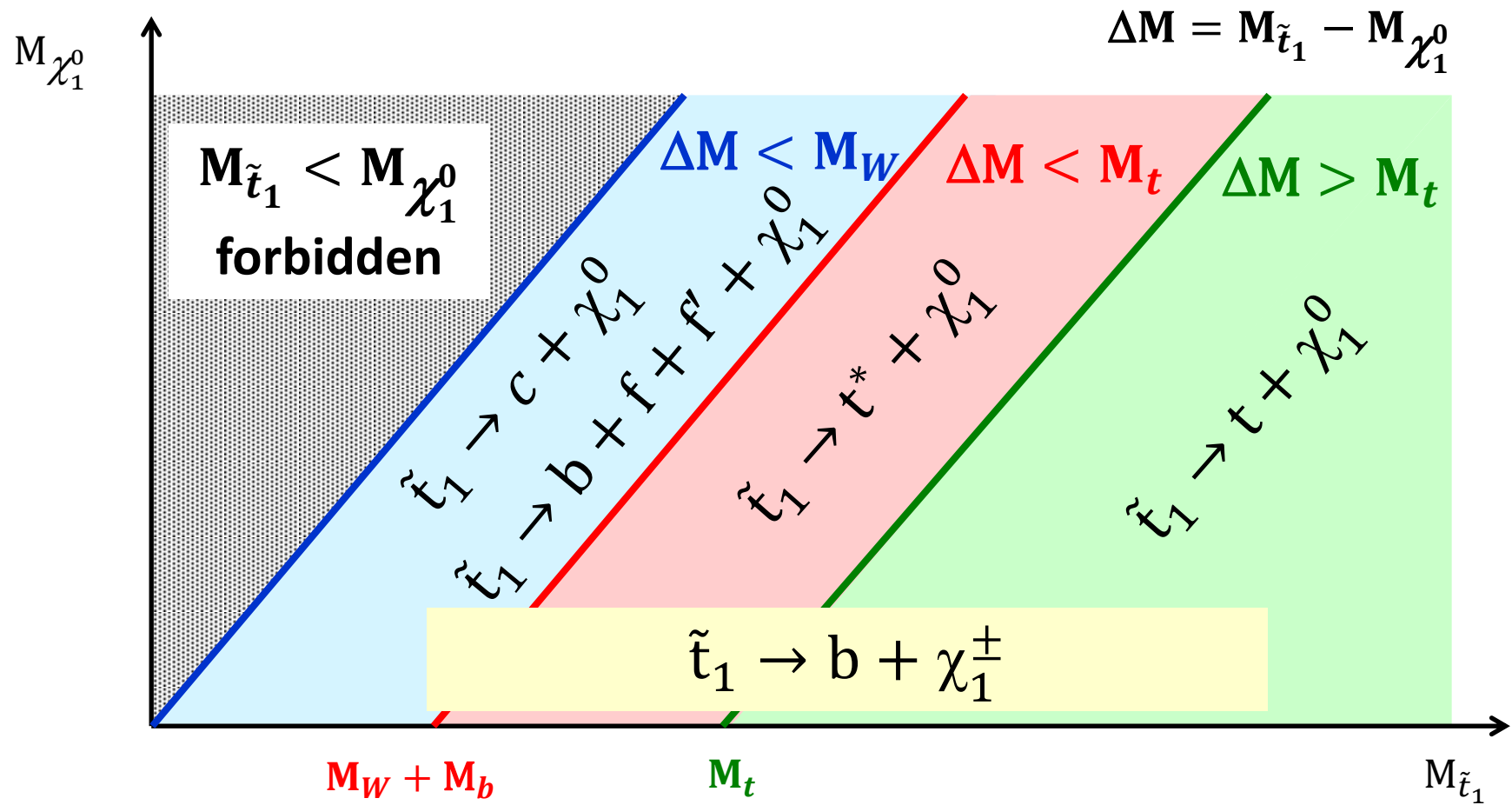
# Top squark production

- For high top squark masses larger cross section increase than for main  $t\bar{t}$  background



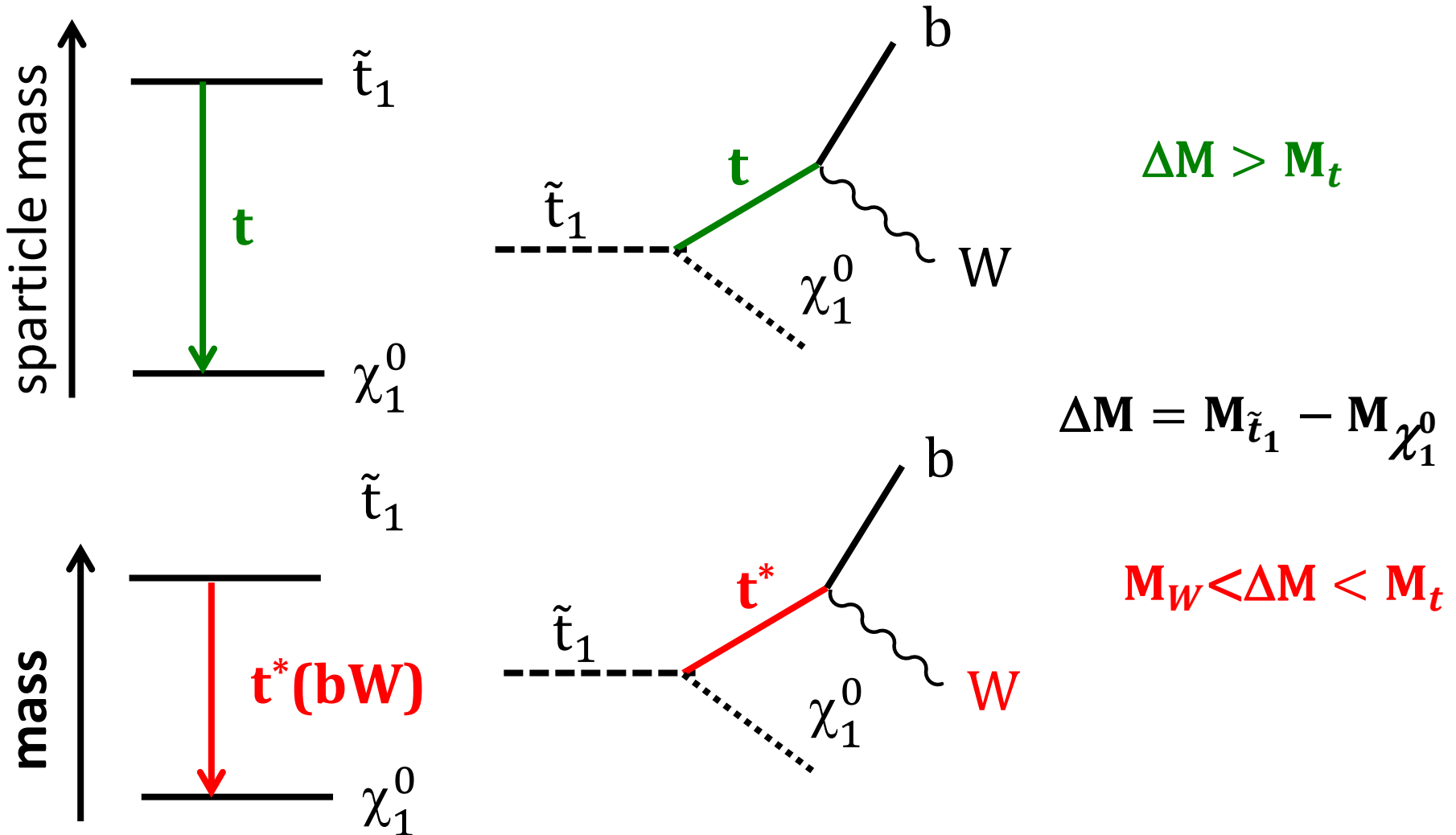
# Top squark decay

- Rich spectrum of decay modes depending on SUSY mass spectrum and couplings



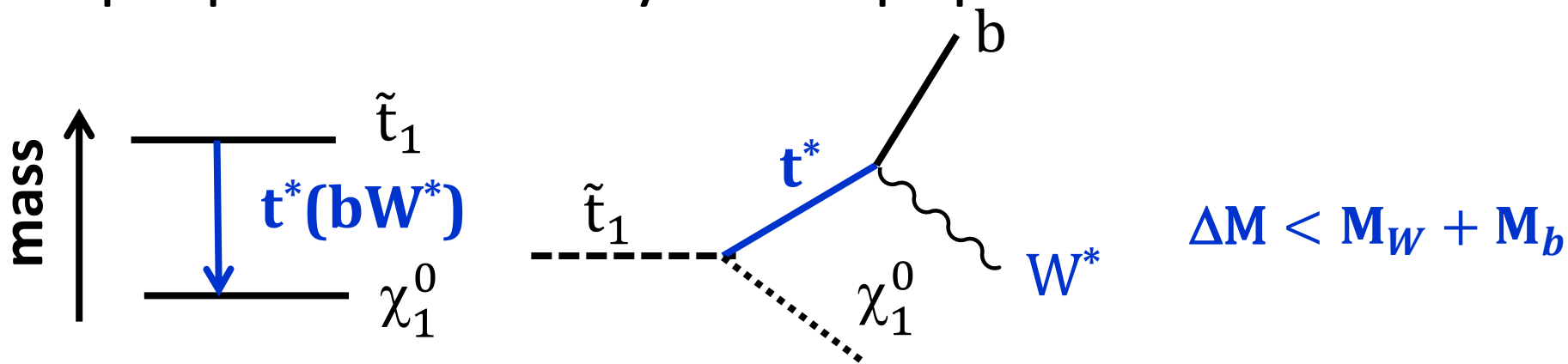
# Top neutralino decay

- Top squark can decay to a top quark and a LSP

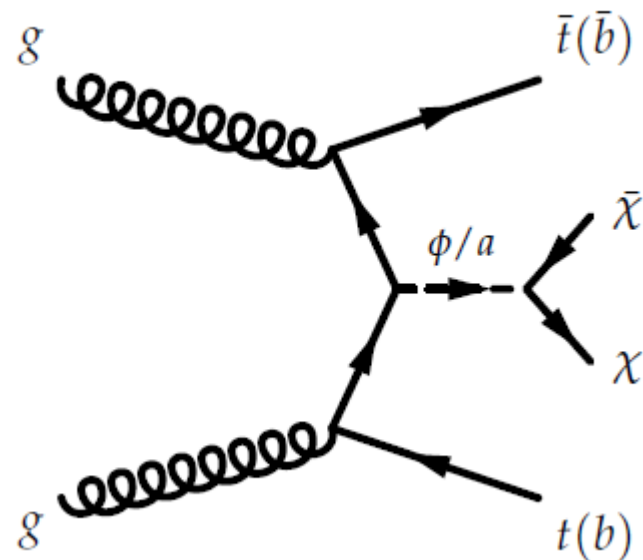


# Top neutralino decay

- Top squark can decay to a top quark and a LSP

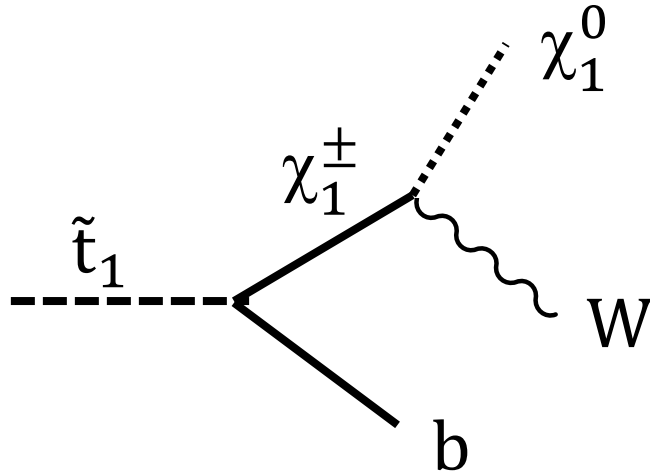


- Same final state can also target direct DM searches with low mass spin-0 mediator
  - In ATLAS same analysis used for both with optimized selection



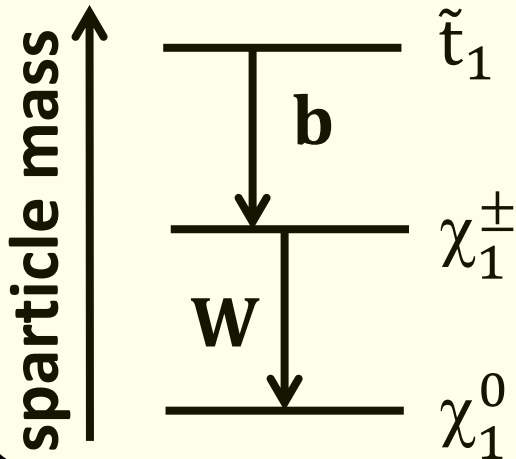
# Bottom chargino decay

- Top squark can decay to a bottom quark and a chargino

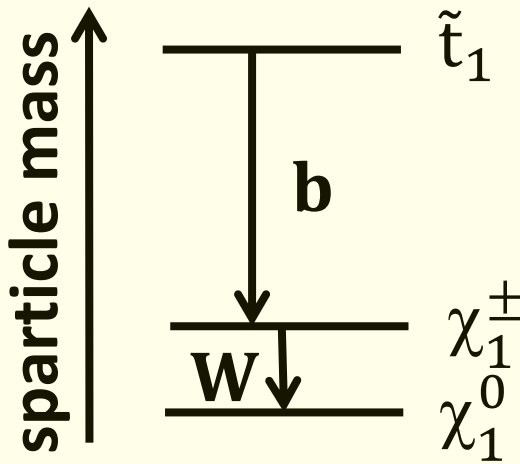


- Same particles in final state as top neutralino decay mode
  - BUT: sensitivity and best analysis depends a lot on mass spectrum

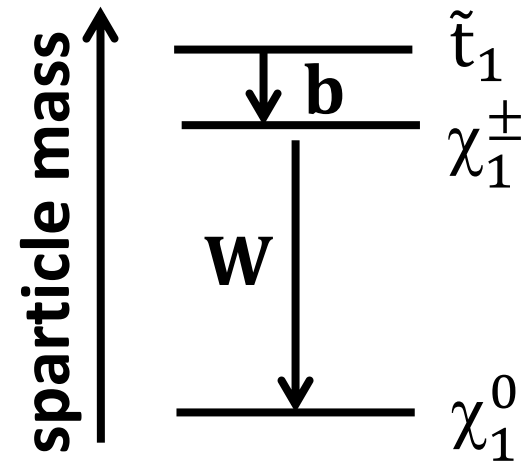
High- $p_T$   $b$  and  $W$



Very soft  $W$  bosons

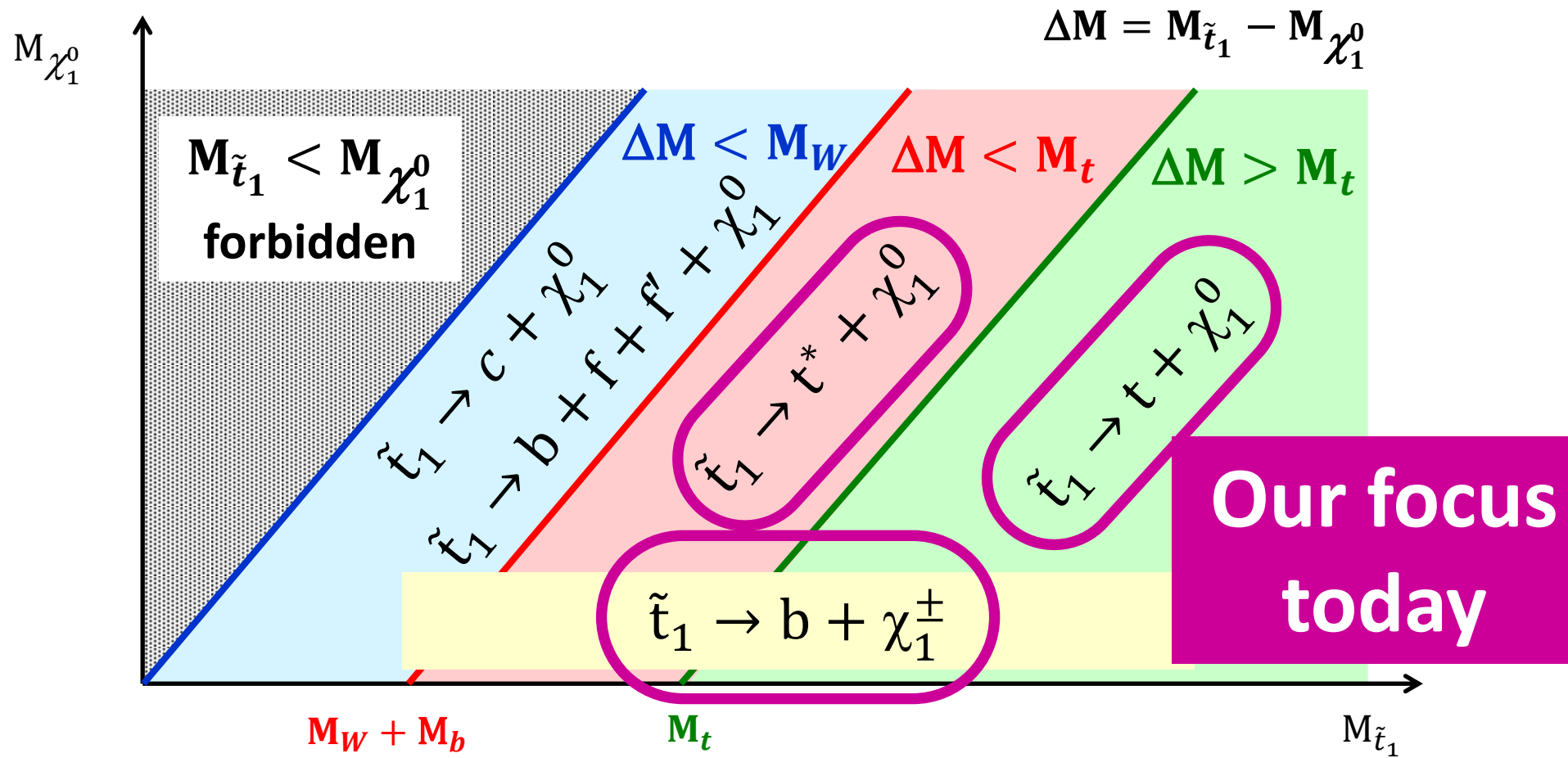


Very soft  $b$ -jets



# Top squark decay

- Rich spectrum of decay modes depending on SUSY mass spectrum and couplings

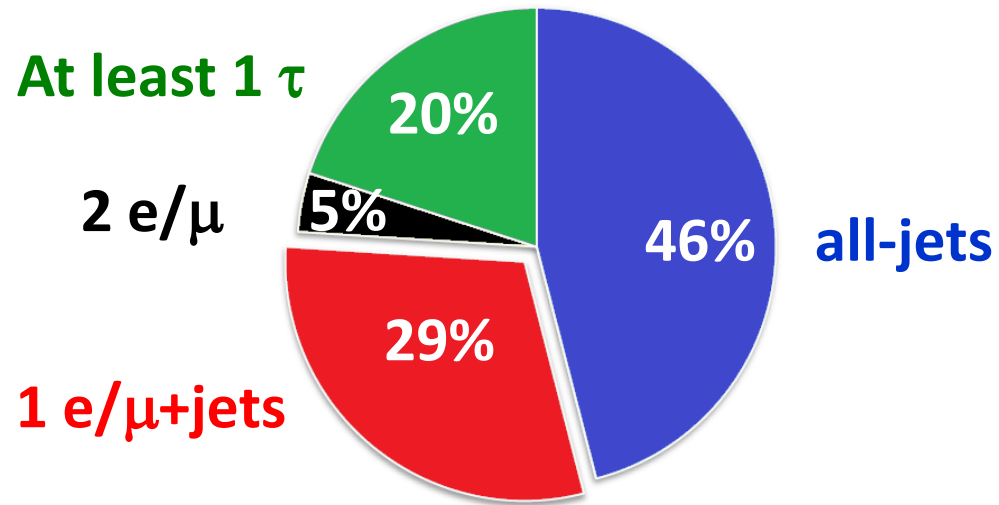


# Experimental signatures for pair production

- Final states (2 top squarks) have two on- or off-shell W bosons
  - Leads to 0, 1 or 2 leptons in the final state

**Without  $p_T$  or  $\eta$  requirements**  
**~30% in the 1  $e/\mu$  final state**

- Referred to as single-lepton



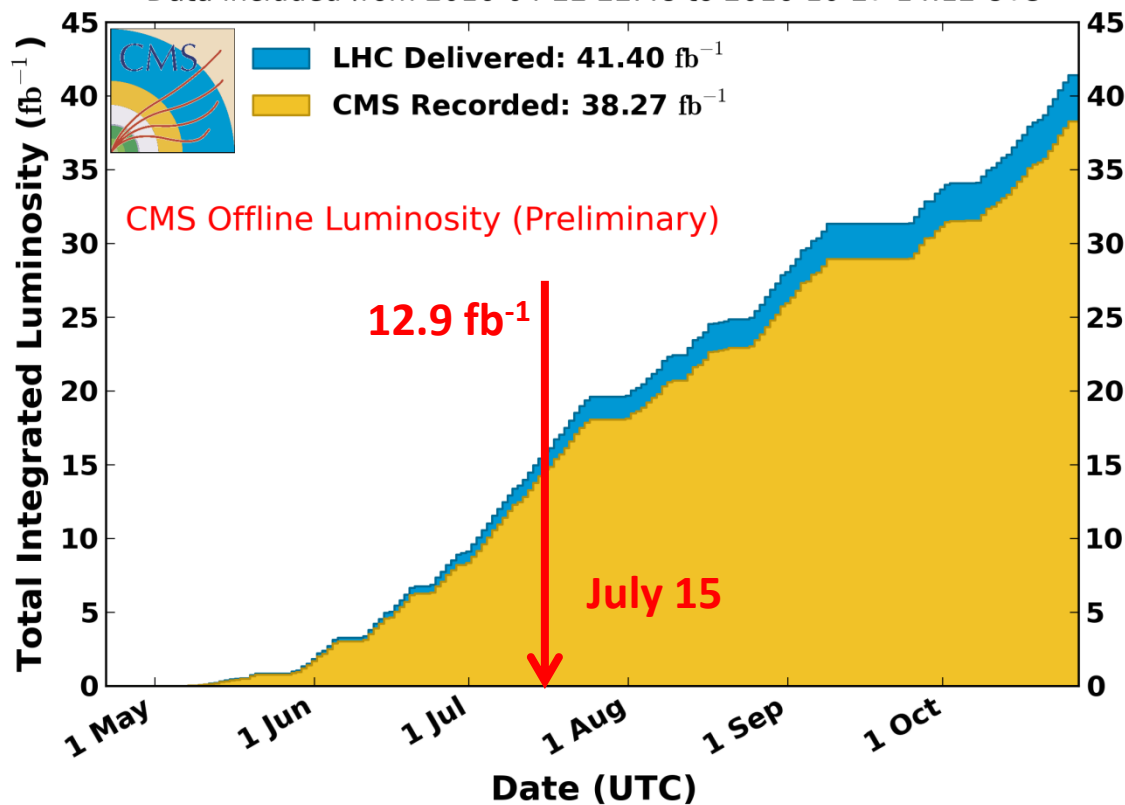
- $\tilde{t}_1 \rightarrow c + \chi_1^0$  a bit different
  - Target with mono-jet like search or use charm-tagging
- Alternative decays possible leading to lots of experimental signature:
  - Long-lived, R-parity violating, GSMB scenario's

# Dataset

- Corresponding to the dataset for summer conference:  $12.9 \text{ fb}^{-1}$
- Total 13 TeV dataset:  $\sim 38.3 \text{ fb}^{-1}$  (2016) +  $2.3 \text{ fb}^{-1}$  (2015)

## CMS Integrated Luminosity, pp, 2016, $\sqrt{s} = 13 \text{ TeV}$

Data included from 2016-04-22 22:48 to 2016-10-27 14:12 UTC





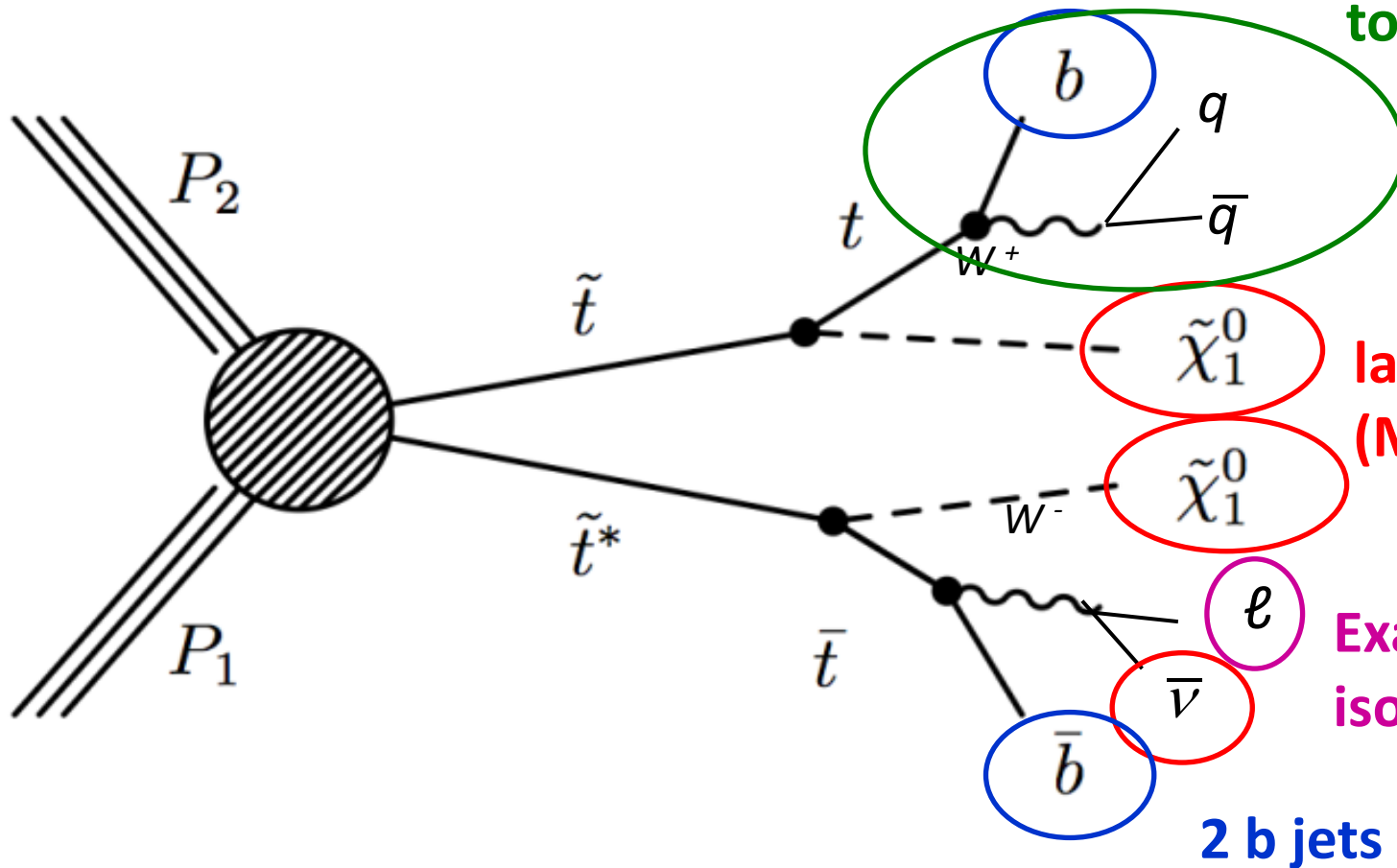
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- **Looking for top squarks in the single-lepton final state**
  - **Baseline selection**
  - Rejecting backgrounds
    - 1l backgrounds
    - 2l backgrounds
    - Rare Standard Model processes
  - Background prediction
  - Results and interpretation

# Single-lepton analysis

Both stops decay as  $\tilde{t} \rightarrow t\chi_1^0$

Hadronic top decay



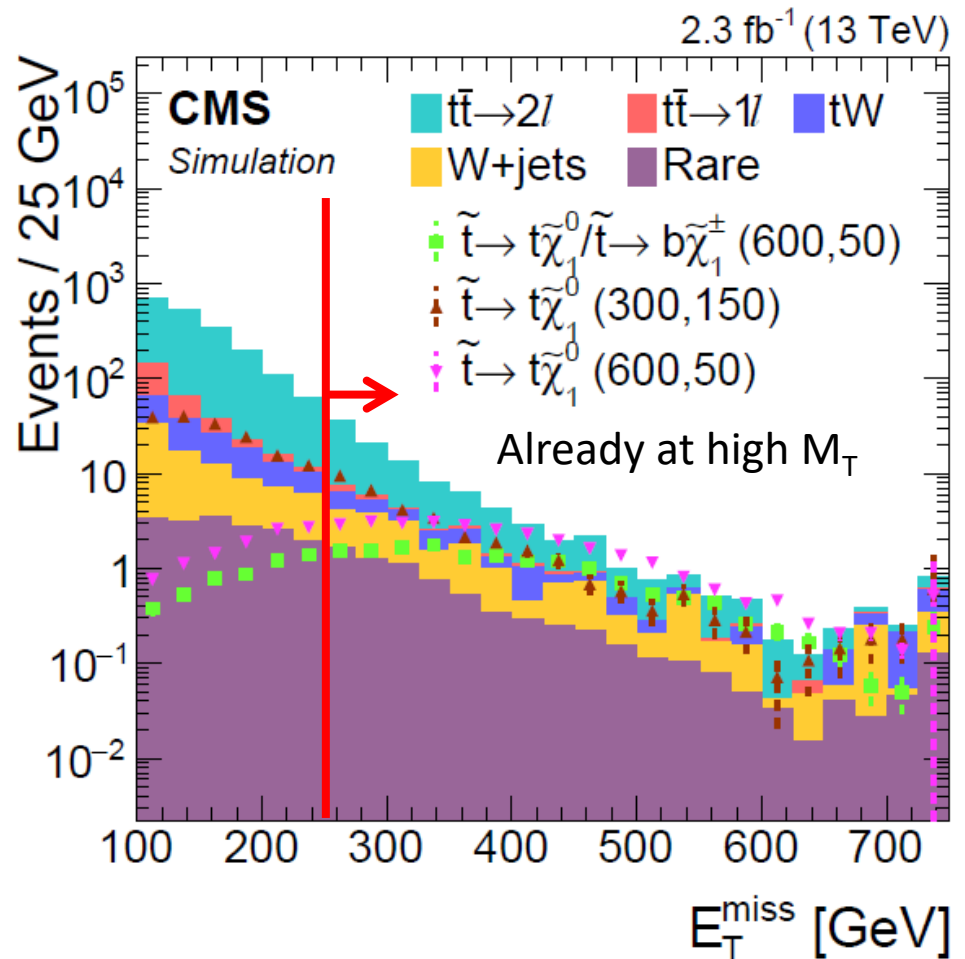
large MET  
( $M_T$ ,  $M_{T2}$ ,...)

Exactly 1  
isolated e/ $\mu$

2 b jets

# Baseline selection

- Selection (trigger-inspired):
  - Exactly 1 electron or muon with  $p_T > 20$  GeV
  - $MET > 250$  GeV
  - At least 1 b-tagged jet
- Main backgrounds:
  - $1\ell t\bar{t}$
  - W+jets
  - $2\ell t\bar{t}$
  - Single top
  - Dibosons,  $t\bar{t}V$



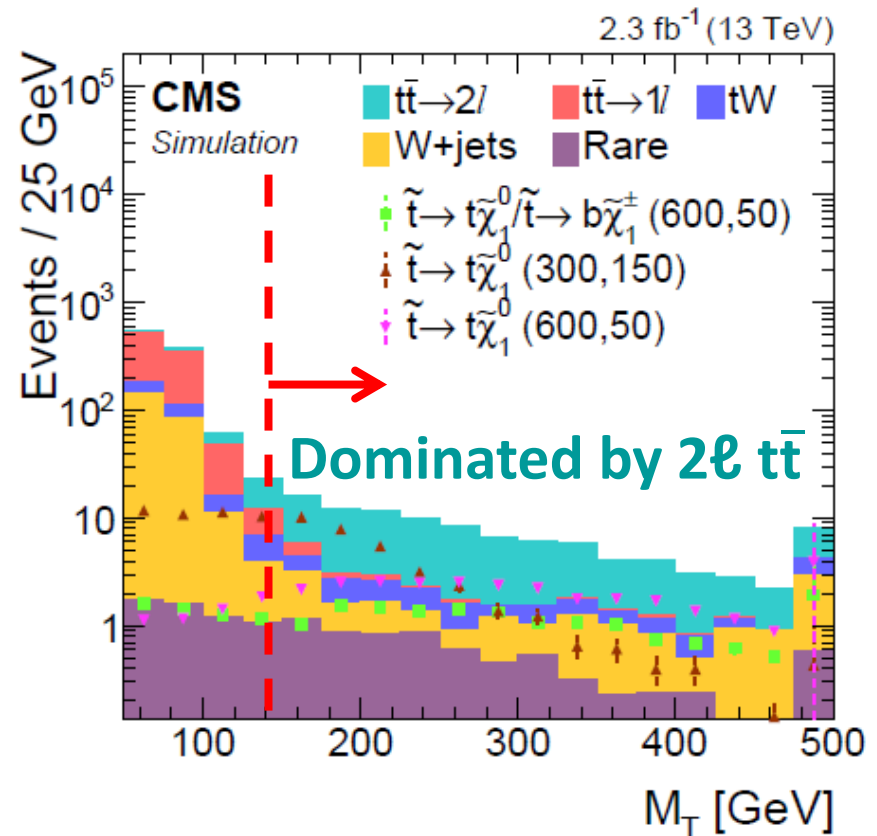
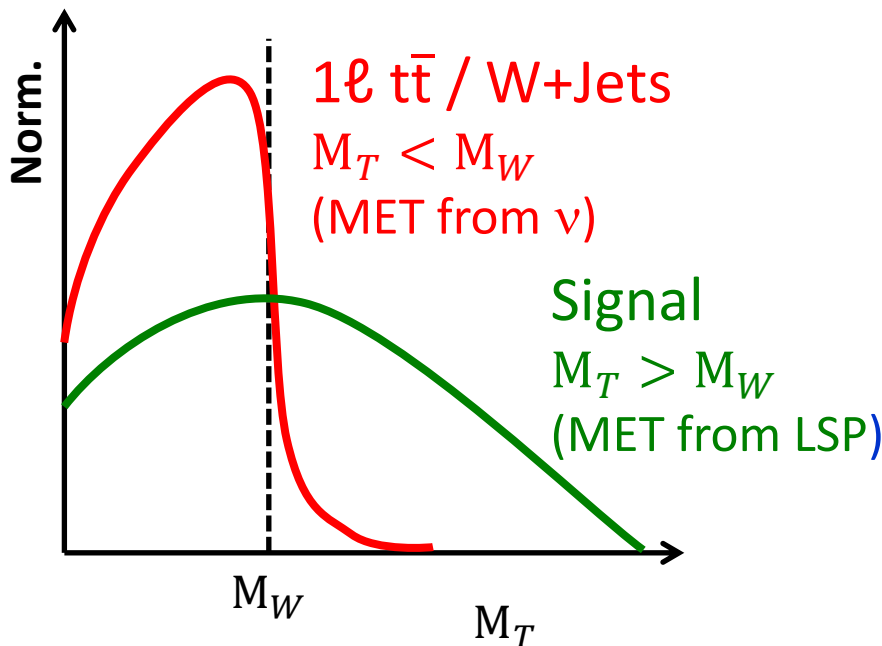
- **Looking for top squarks in the single-lepton final state**
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# Killing the single-lepton background

- Largest two backgrounds are  $t\bar{t}$  and W+Jets
  - No endpoint in transverse mass due to extra MET from LSPs

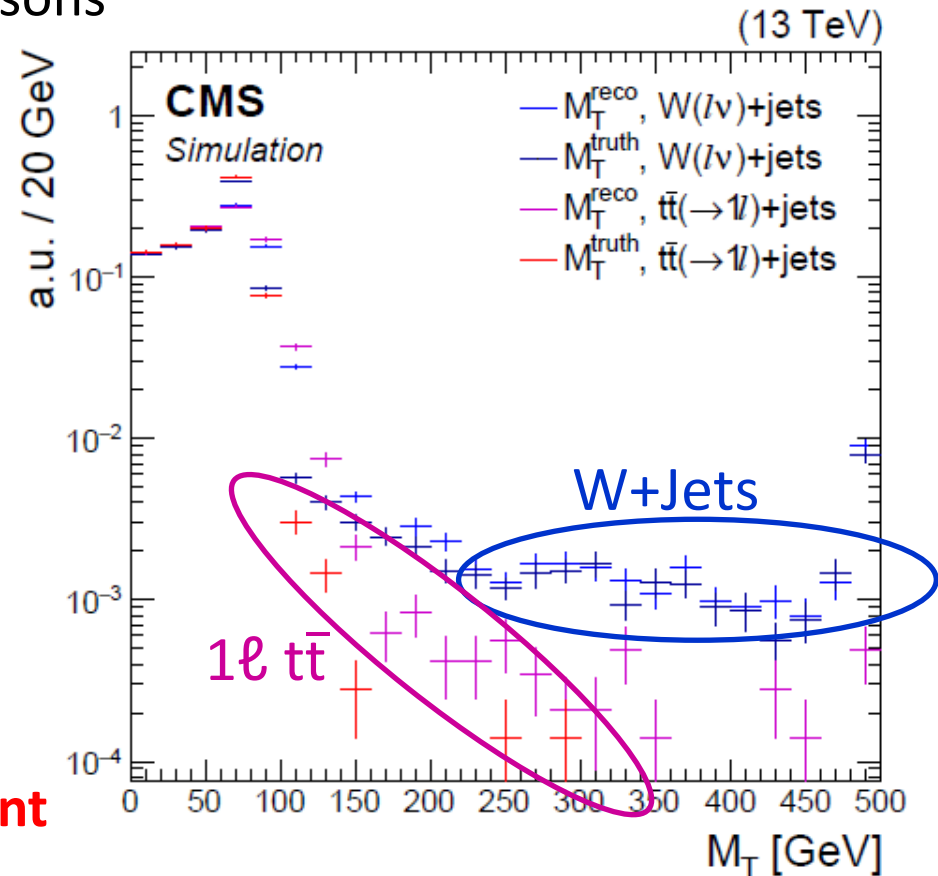
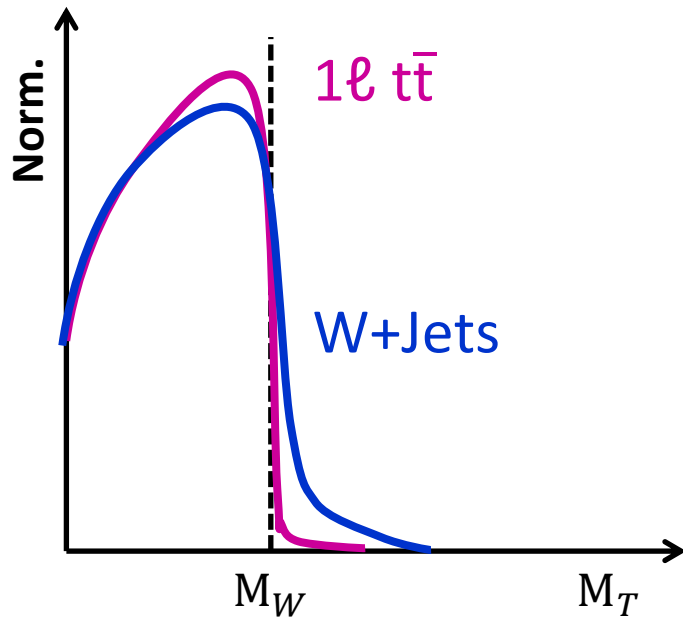
$$M_T^1(\ell, \nu)^2 = (E_T(\ell) + E_T(\nu))^2 - (\vec{p}_T(\ell) + \vec{p}_T(\nu))^2$$

$$\rightarrow 2E_T(\ell)E_T(\nu)(1 - \cos(\Delta\phi))$$



# Killing the single-lepton background

- Transverse mass point more efficient for  $1l t\bar{t}$ 
  - Top mass imposes kinematic constraint on W mass
    - Tail for  $t\bar{t}$  dominated by resolution effects, for W+Jets by off-shell W bosons



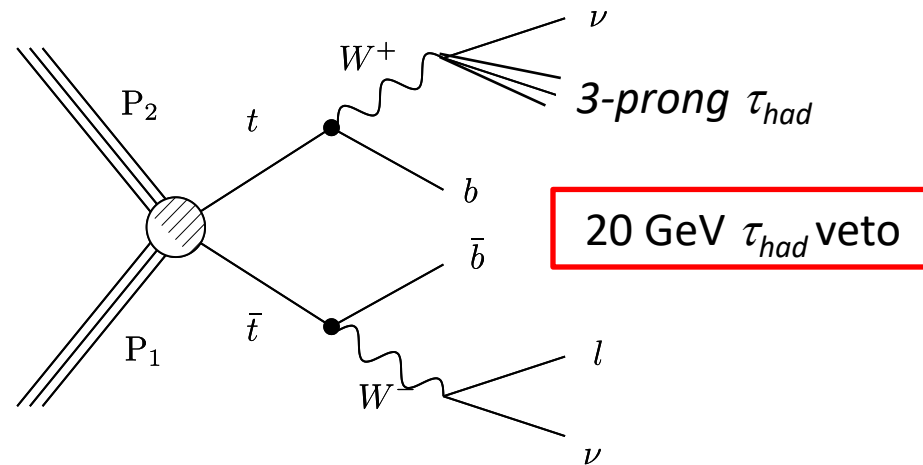
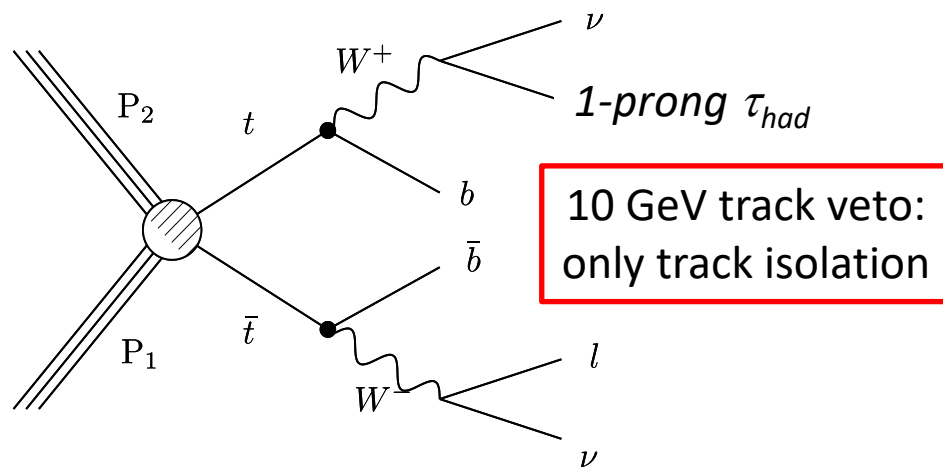
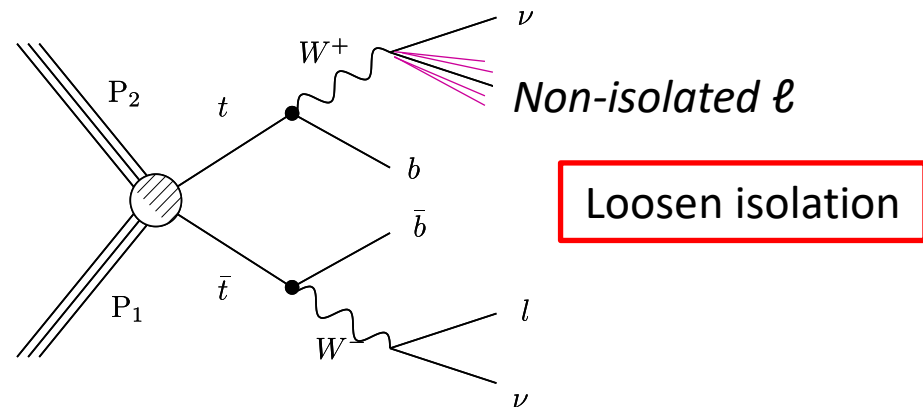
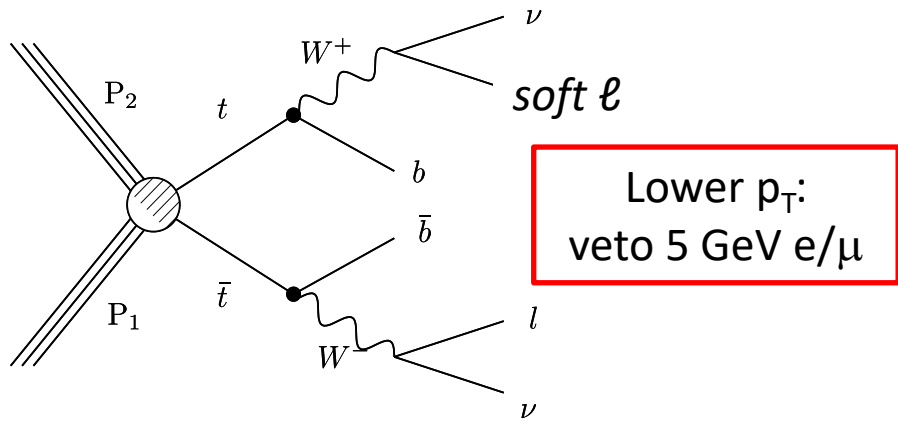
Reduce W+Jets with b-tag requirement

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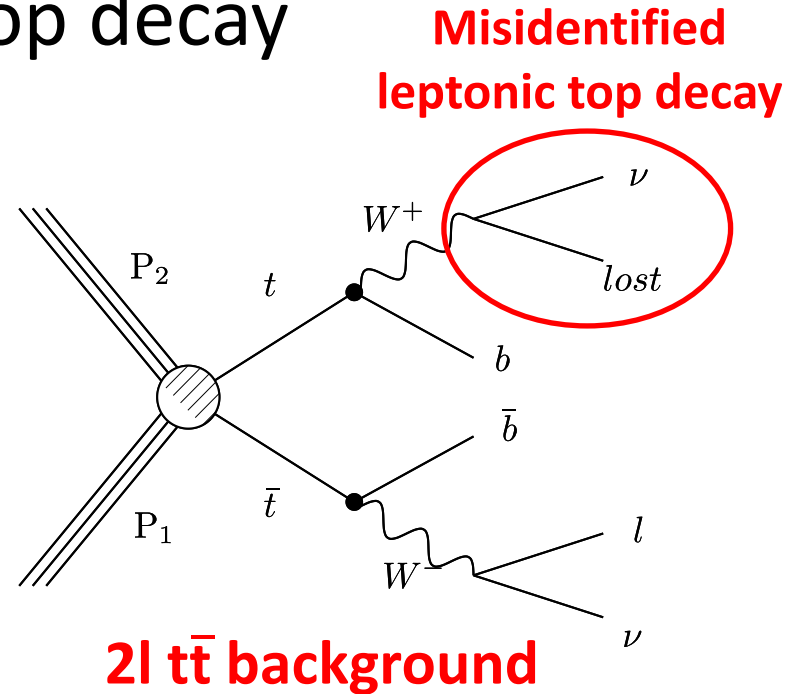
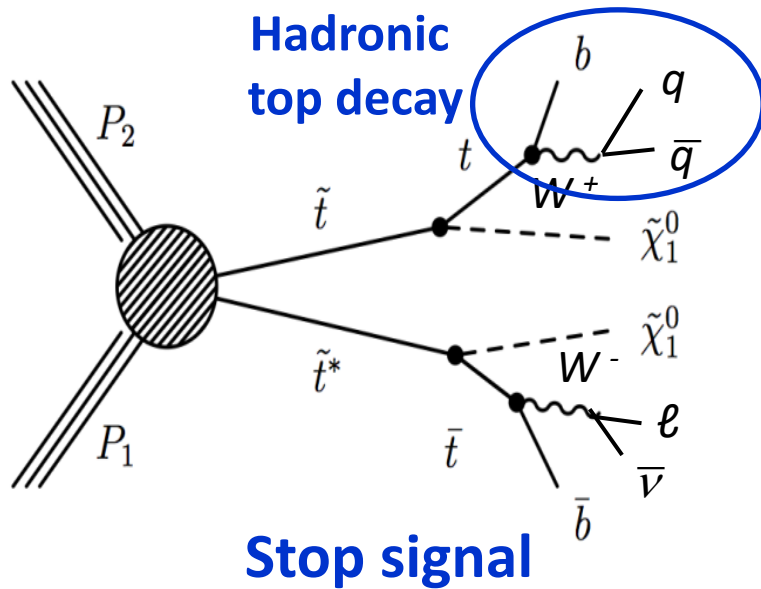
# Removing dilepton backgrounds

- After tight  $M_T$  cut dilepton backgrounds ( $t\bar{t}$  and  $tW$ ) dominate
  - Try to minimize this background by optimizing lepton vetoes



# Hadronic top decays

- Only signal has a hadronic top decay



- Require at least 4 jets
  - 3 for boosted scenario's when W jets will be merged
- Can be exploited further with (low- $p_T$ ) top tagger

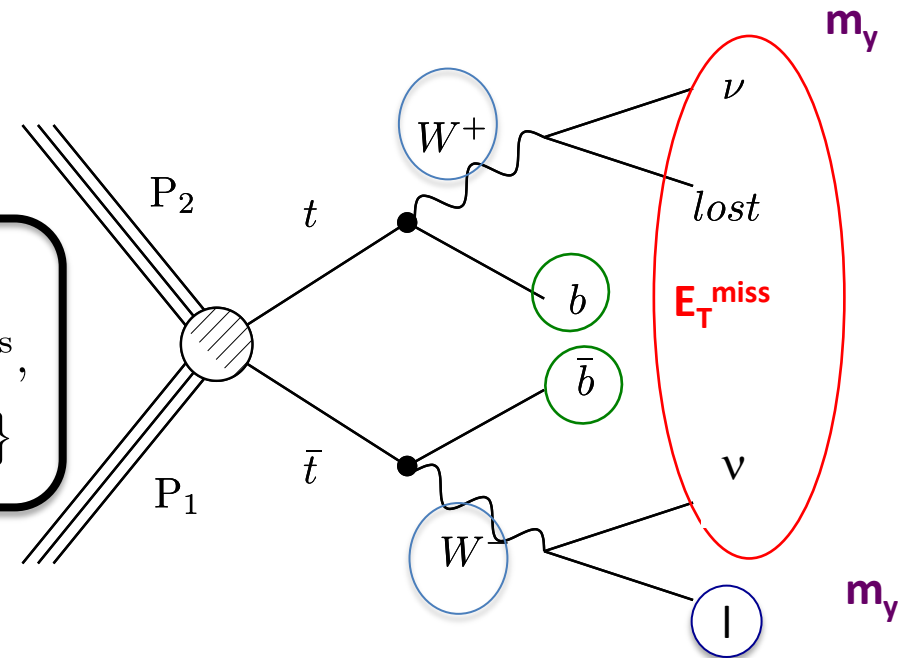
# Attacking the dilepton background

- After tight  $M_T$  cut dilepton backgrounds ( $t\bar{t}$  and  $tW$ ) dominate
  - Target lost lepton background with dedicated variable

$M_{T2}^W$ : minimum mother particle mass compatible with W mass and MET constraints for 2l  $t\bar{t}$  hypothesis

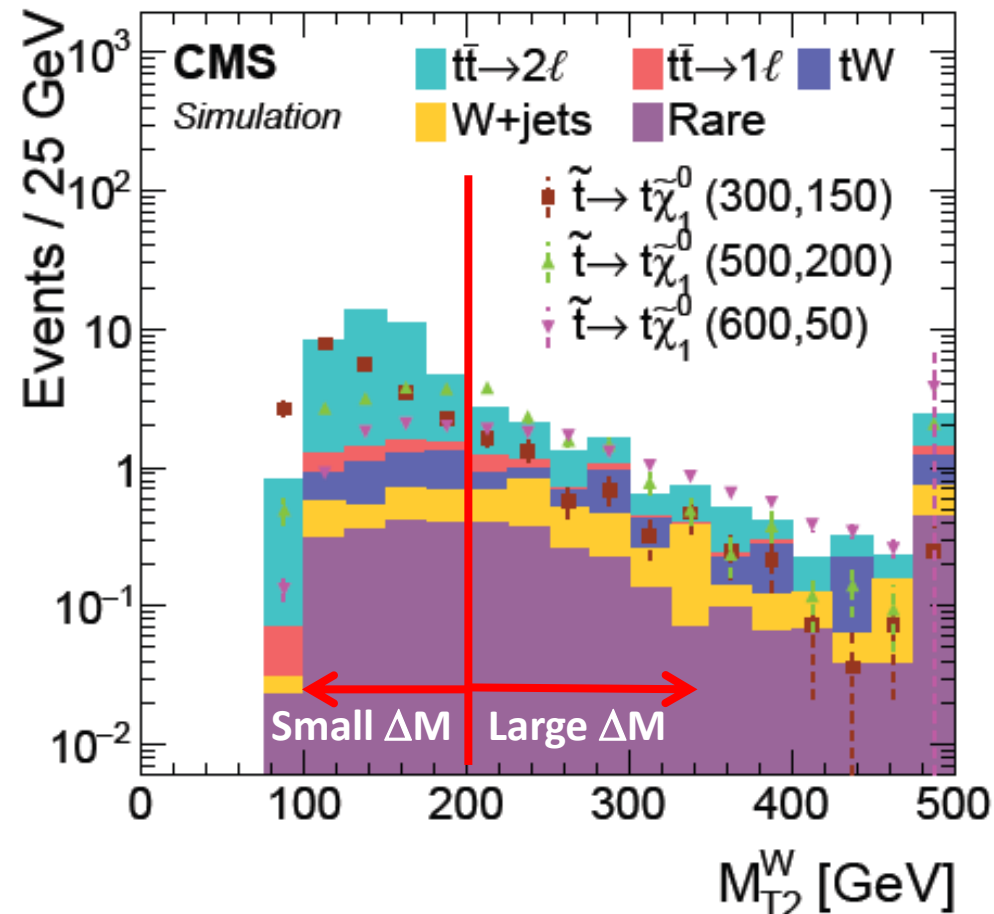
$$M_{T2}^W = \min \left\{ m_y \text{ consistent with } \left[ p_1^2 = 0, \right. \right. \\ \left. \left. (p_1 + p_\ell)^2 = p_2^2 = M_W, \vec{p}_T^1 + \vec{p}_T^2 = \vec{E}_T^{\text{miss}}, \right. \right. \\ \left. \left. (p_1 + p_\ell + p_{b_1})^2 = (p_2 + p_{b_2})^2 = m_y^2, \right] \right\}$$

arXiv:1203.4813



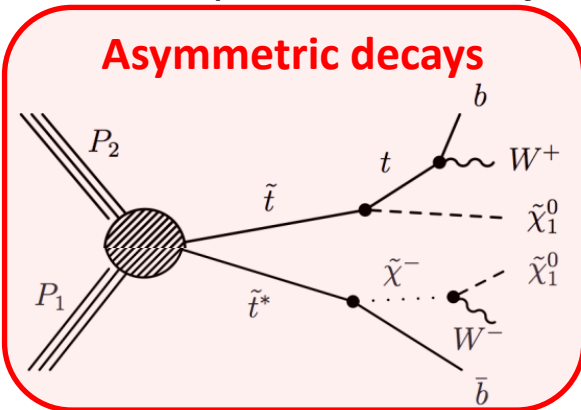
# Attacking the dilepton background

- After tight  $M_T$  cut dilepton backgrounds ( $t\bar{t}$  and  $tW$ ) dominate
  - Use  $M_{T2}^W$ : minimum mother particle mass compatible with W mass and MET constraints assuming  $2l\bar{t}\bar{t}$
- Natural endpoint around 200 GeV
- Compressed mass points at low  $M_{T2}^W$
- Tails for SM backgrounds due to different effects
  - $tW$  does not obey mass constraint
    - Only 1 real b tag
  - Only 17% of  $2l\bar{t}\bar{t}$  with high generator- $M_{T2}^W$ 
    - 9%: overlap between jets
      - Jet substructure?
    - 56%: one of the jets not in kinematic acceptance
      - Loosen  $p_T$  cut (30 GeV) on jets?
      - Allow for missing a b-jet
    - 18%: wrong jet choice
      - Investigated different algorithms to chose

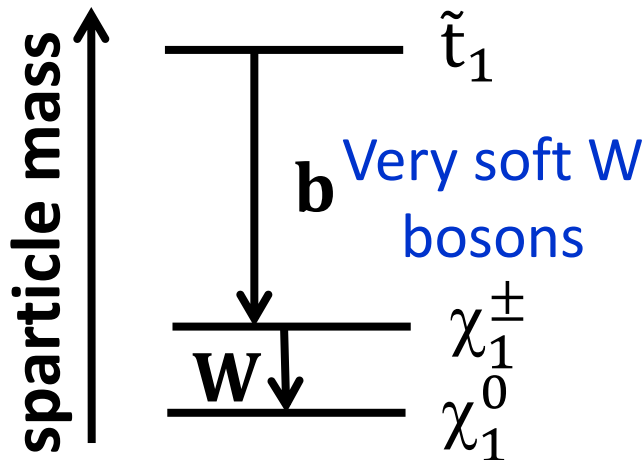


# Modified topness

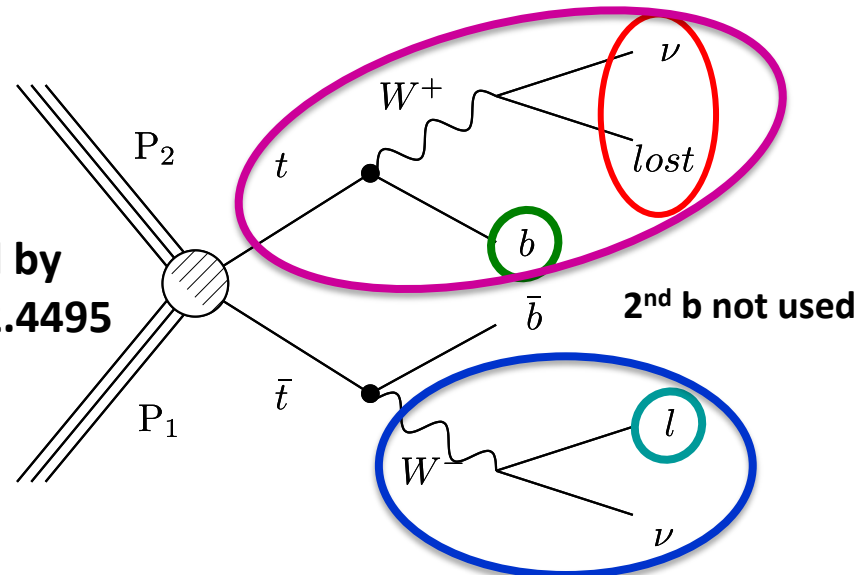
- Remove some of the background where one of the jets is not in the acceptance
  - Also helps for signal acceptance for asymmetric decays
  - Improved tW rejection



$$t_{\text{mod}} = \ln(\min S) \quad \text{with} \quad S(\vec{p}_W, p_{\nu, z}) = \underbrace{\frac{(m_W^2 - (p_\nu + p_\ell)^2)^2}{a_W^4}}_{\text{Constrained to } M_W} + \underbrace{\frac{(m_t^2 - (p_{b_2} + p_W)^2)^2}{a_t^4}}_{\text{Constrained to } M_{\text{top}}}$$



Inspired by  
arXiv:1212.4495



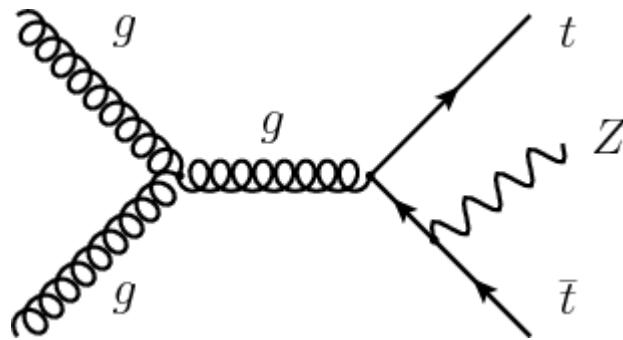
- **Looking for top squarks in the single-lepton final state**
  - Baseline selection
  - **Rejecting backgrounds**
    - 1l backgrounds
    - 2l backgrounds
    - **Rare Standard Model processes**
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# Rare SM backgrounds

- Rare SM backgrounds enter through  $Z \rightarrow \nu\nu$  decays:

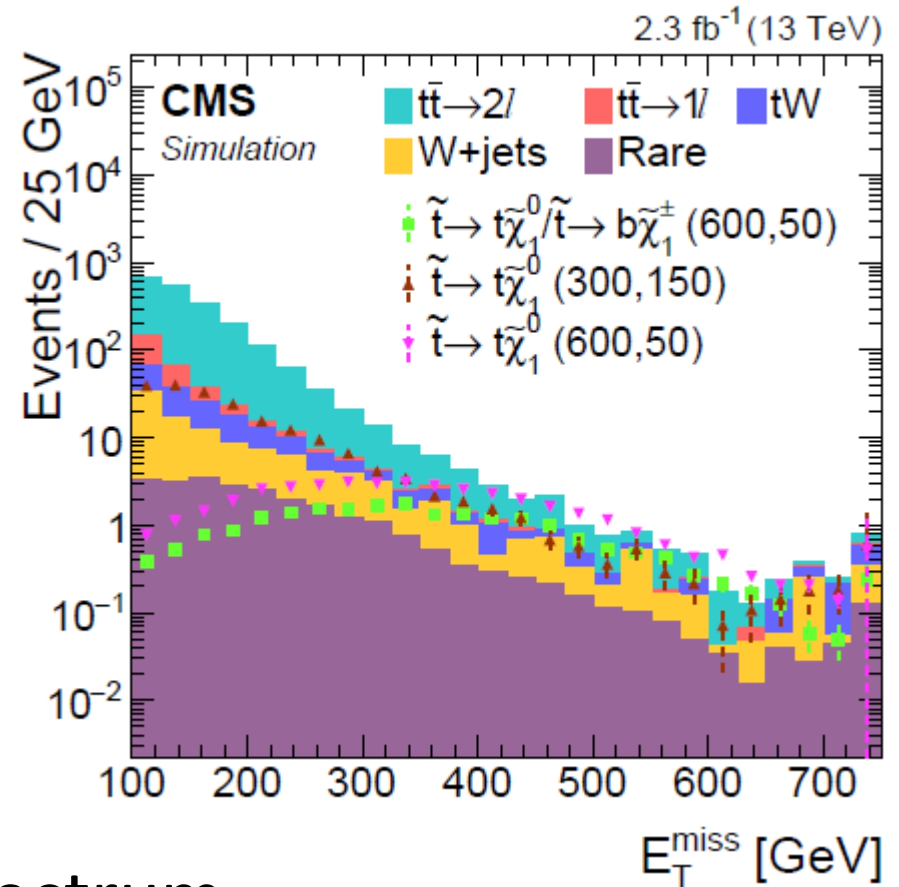
–  $t\bar{t}Z$ :

- Almost indistinguishable
- Real MET, 2 b quarks, hadronic top decay



– Also WZ, tqZ,...

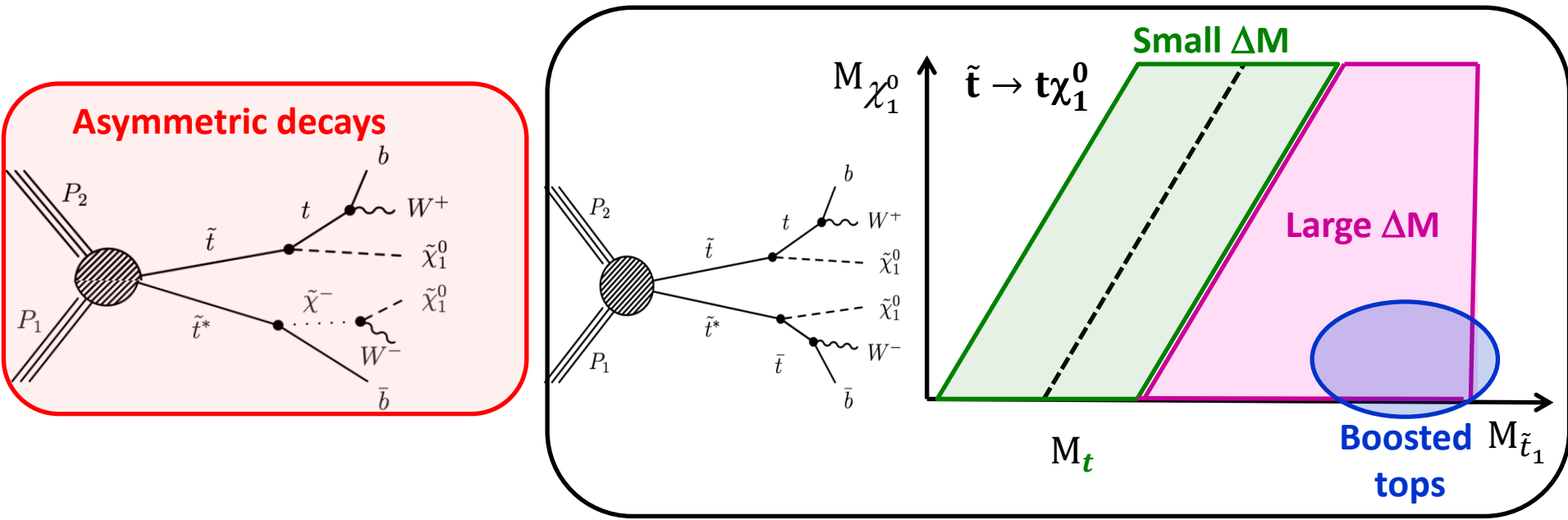
- Use differences in MET spectrum



# Selection summary

- Exclusive search regions designed to target specific regions
  - Extra gain due to statistical combination (low  $\Delta M$ )

	$N_{\text{jets}}$	$M_{T2}^W$ [GeV]	$t_{\text{mod}}$	$E_T^{\text{miss}}$ [GeV]		
<b>Asymmetric decays</b> ←	= 2		> 6.4	250–350	350–450	> 450
<b>Boosted tops</b> ←	= 3	> 200		250–350	350–450	450–550 > 550
<b>Small <math>\Delta M</math></b> ←	$\geq 4$	$\leq 200$		250–350	350–450	> 450
<b>Large <math>\Delta M</math></b> ←	$\geq 4$	> 200		250–350	350–450	450–550 550–650 > 650

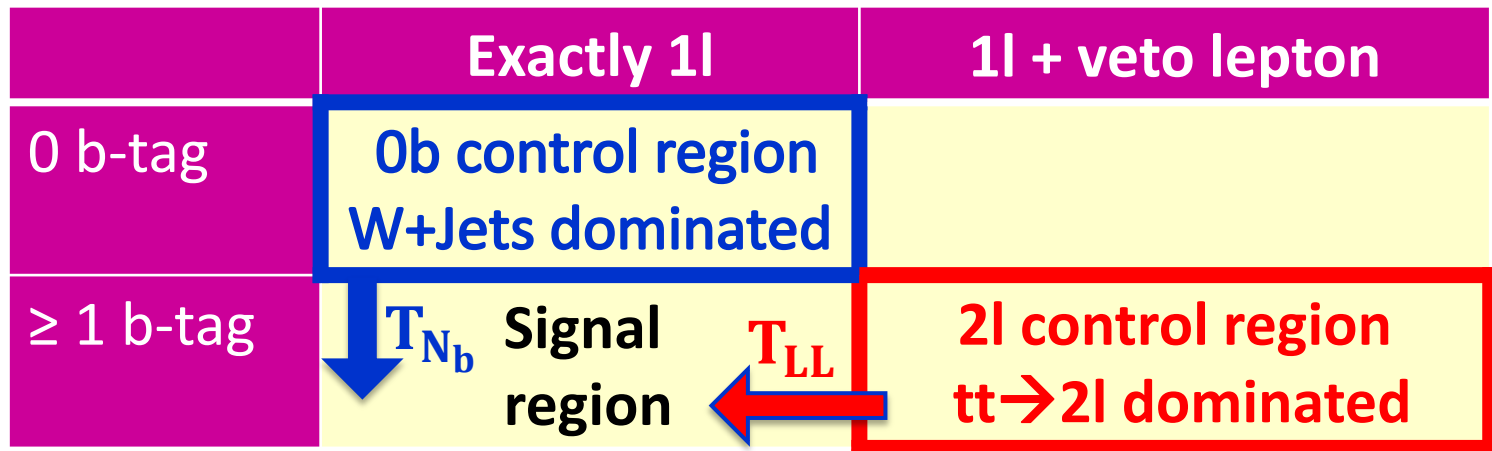




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# Background prediction

- Dominant background predicted in data-driven way
  - Use control regions very similar to signal region



- Smaller backgrounds taken from simulation after estimating experimental and theoretical uncertainties
  - Rare SM backgrounds
    - First checks done to get normalization from 3l CR and distributions from  $t\bar{t}$
  - 1l  $t\bar{t}$ : dominated by MET resolution
  - W+Jets in low  $M_{T2}^W$  regions:
    - even 0b control region dominated by lost lepton background

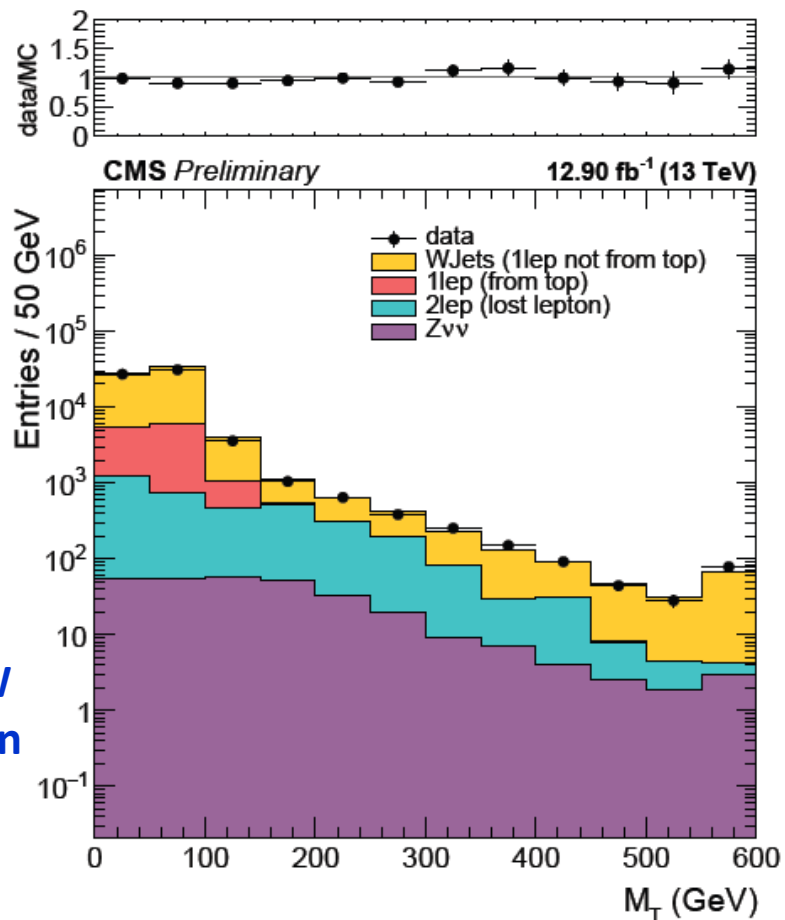
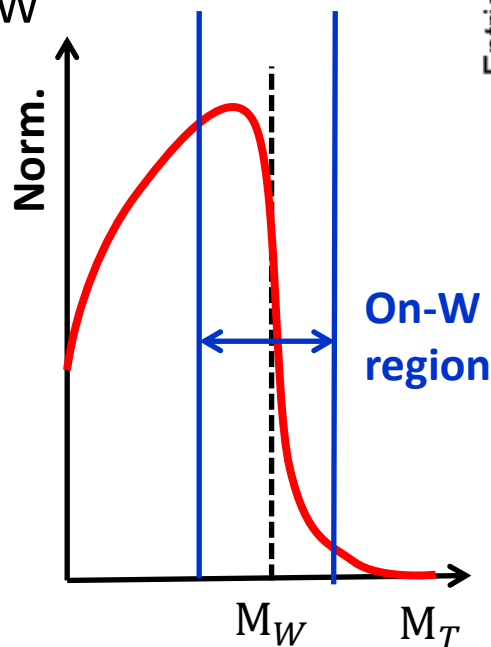
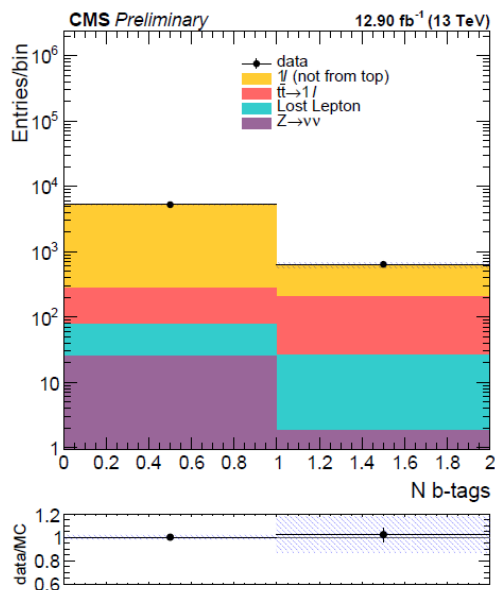
# W+Jets background

- Extrapolation in number of b-tagged jets

$$N_{W+jets}^{pred} = (N_{N_b=0}^{data} - N_{N_b=0}^{non-W+jets}) T_{N_b}$$

- Most important systematic uncertainty for b-tag extrapolation:

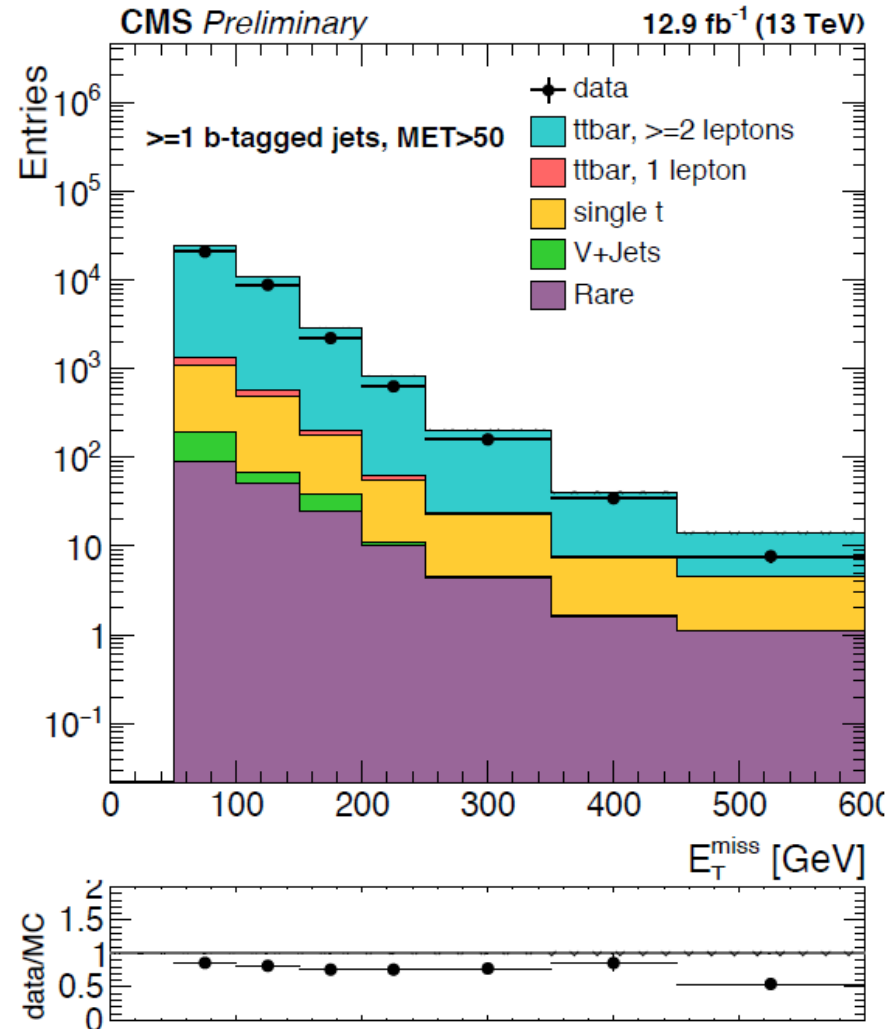
- b-tagging efficiency from data
- Uncertainty for W+b(b) fraction
  - Control region on-W



# Lost lepton background

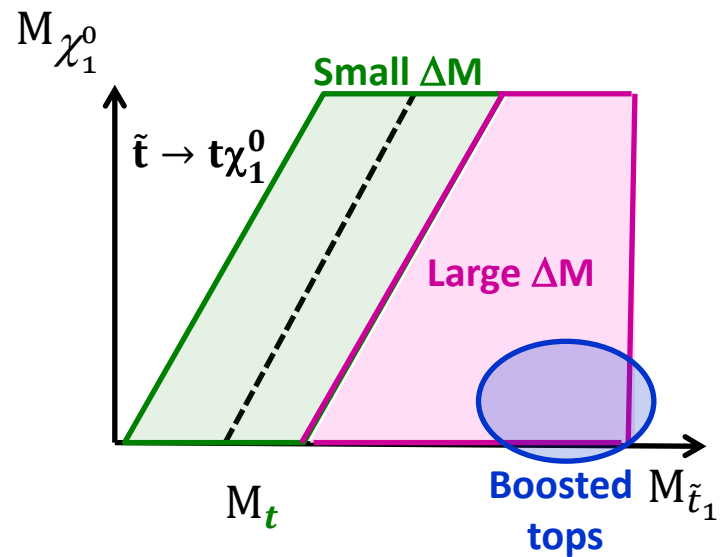
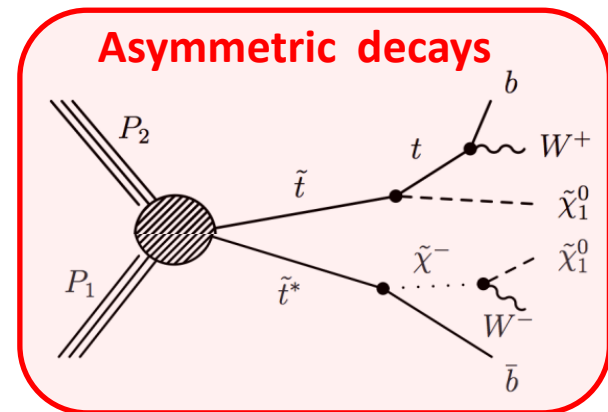
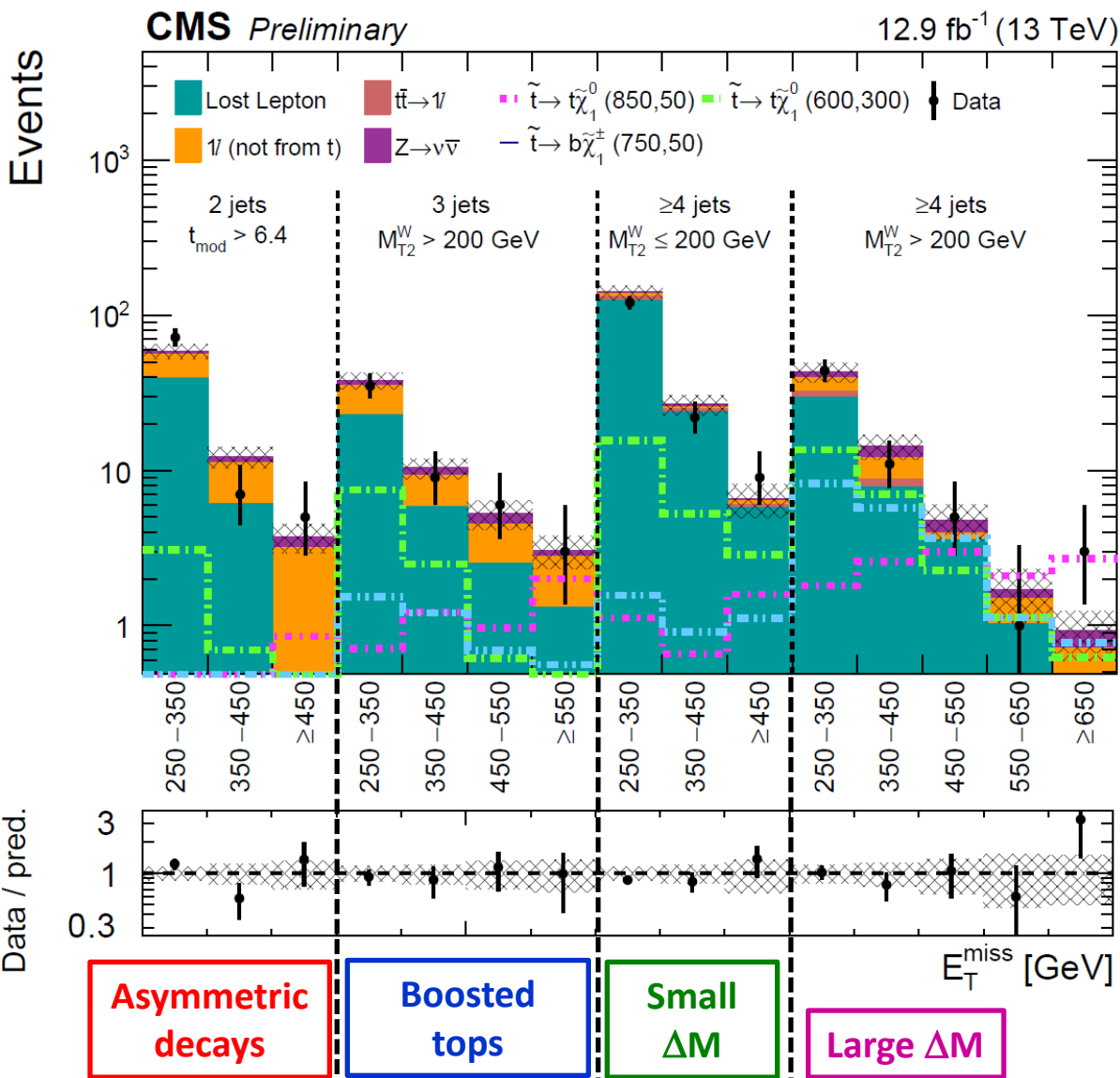
- Lost lepton background dominated by  $2\ell t\bar{t}$ , but also includes  $tW$ 
  - Enters signal region when one lepton out of **acceptance or not identified/isolated**
- Normalize the estimate to a dilepton control region in data
  - Require extra  $e/\mu$  ( $p_T > 10$  GeV)
  - Only extrapolate along the lost lepton category
    - Large statistical uncertainties, small systematic uncertainties

$$N_{lost\ell}^{Data, SR} = N_{\ell\ell}^{Data, CR} \times \frac{M_{lost\ell}^{MC, SR}}{M_{\ell\ell}^{MC, CR}}$$



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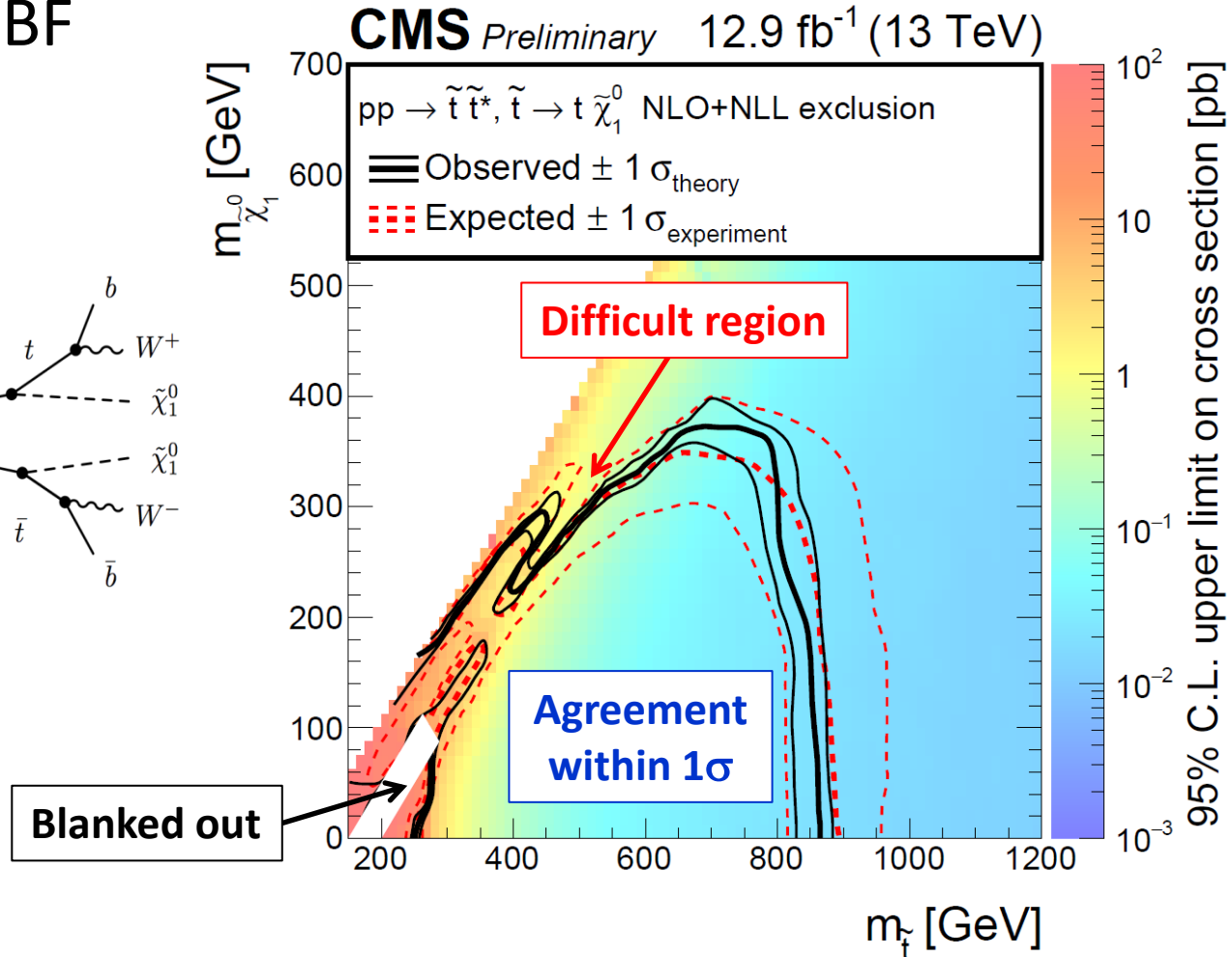
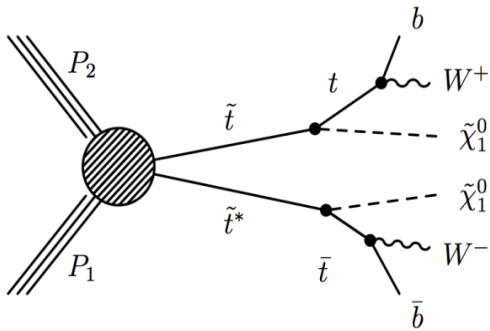
# No excess observed



# Interpretation

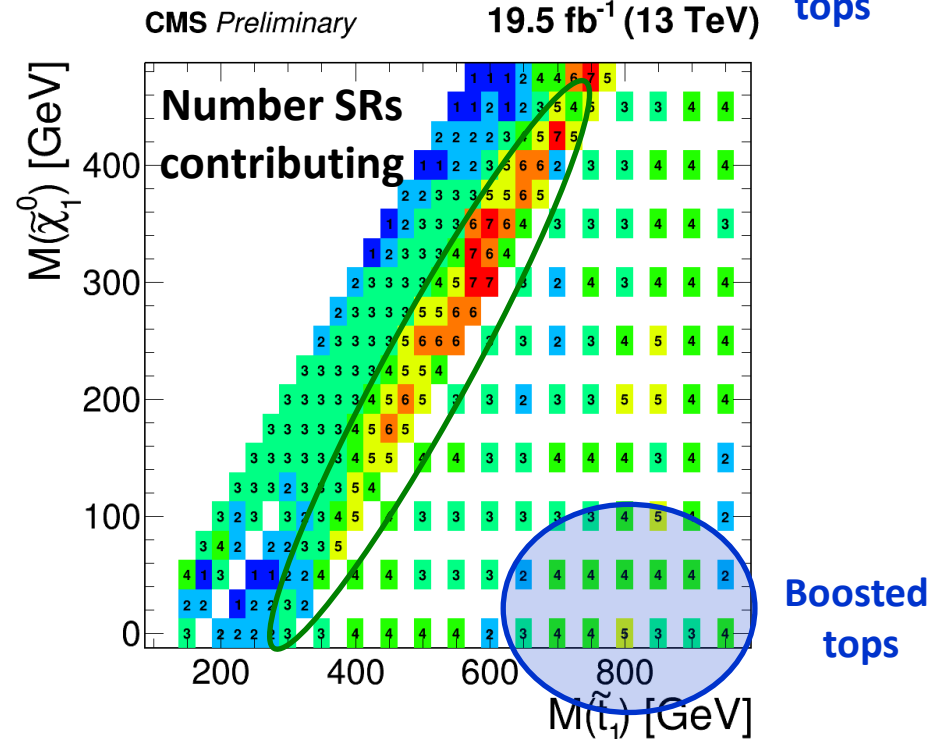
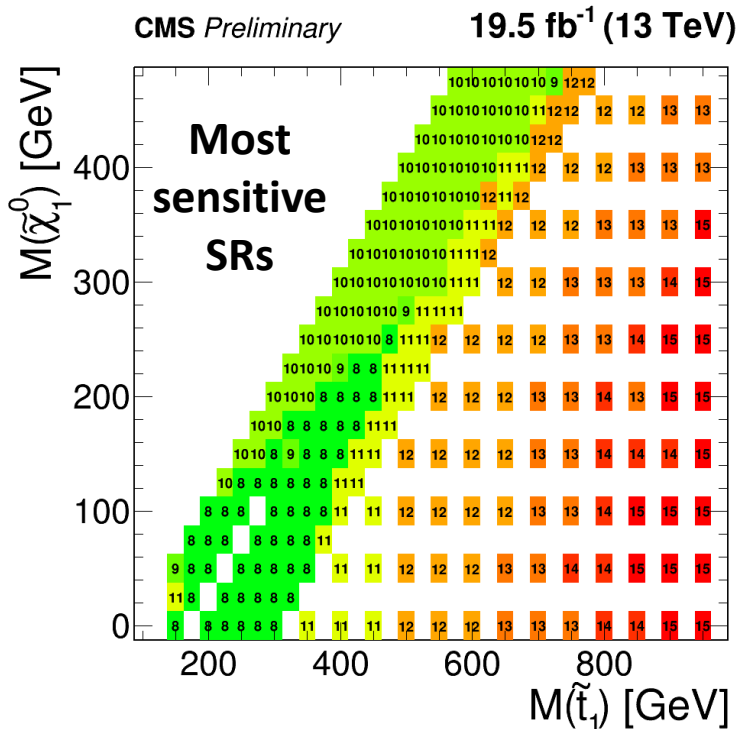
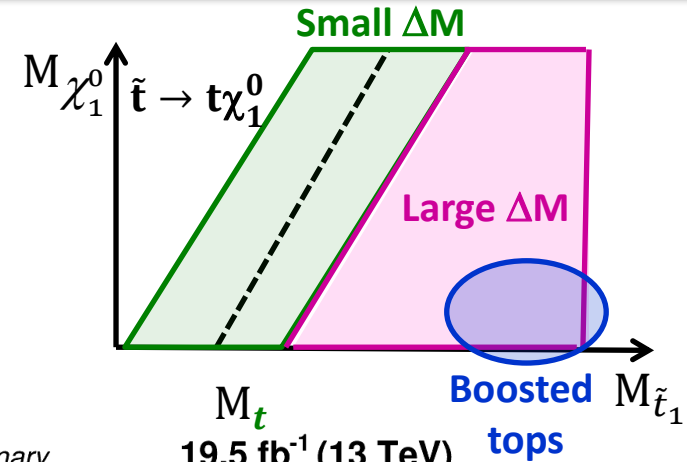
- Top squarks probed up to 870 GeV

– 100% BF



# Interpretation

- Most sensitive search regions divided between small and large  $\Delta M$
- Small mass splittings get sensitivity from a large group of search regions
- High mass tops have extra contributions from boosted SRs

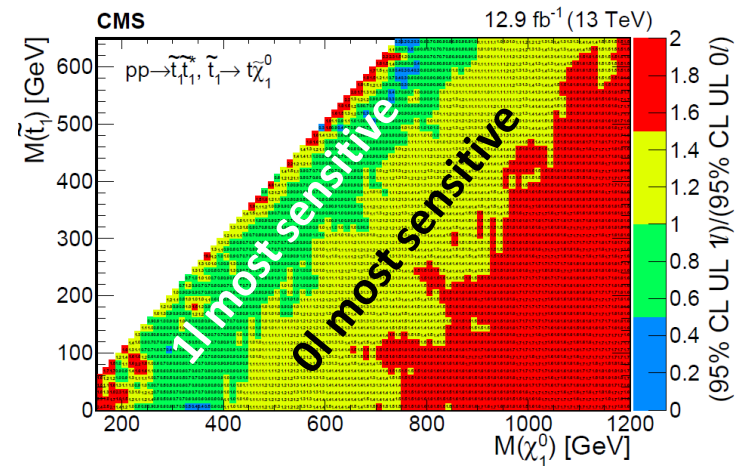
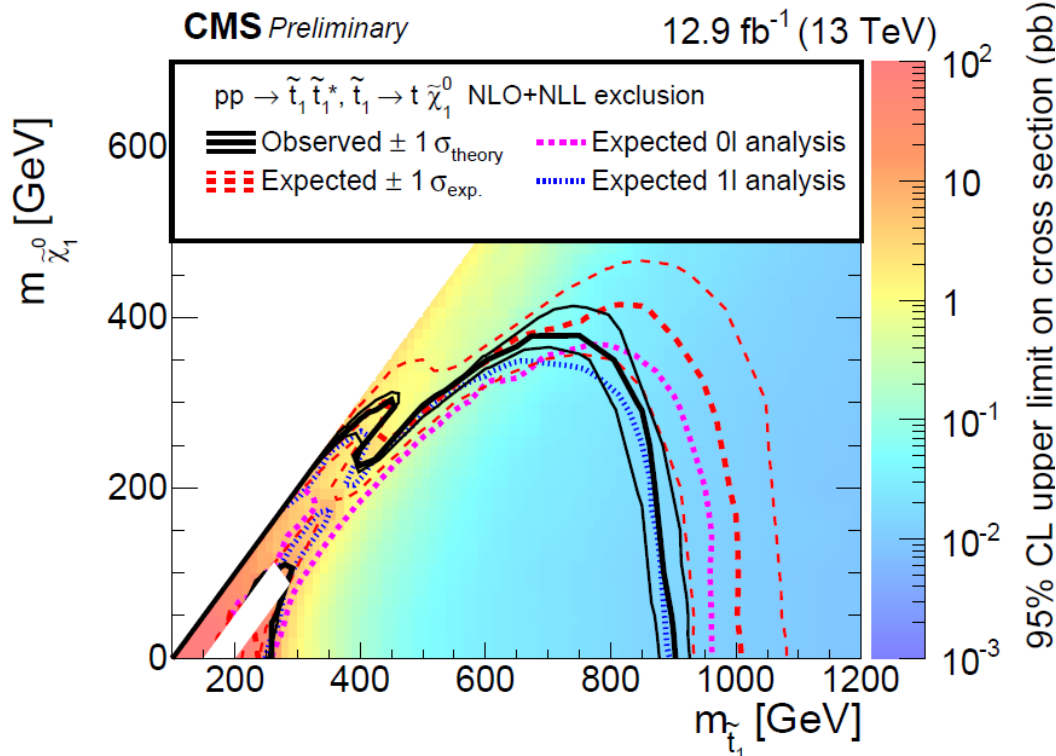




# Complementarity with all-jets analysis

- To believe a new-physics discovery we will also want to see multiple decay modes
  - Higgs boson discovery both in ZZ and  $\gamma\gamma$
- Make sure both analyses are fully disjoint, use same prescriptions for systematic uncertainties and statistical treatment

**Breaking 1 TeV frontier!**



**For intermediate chargino:  
1l always the most sensitive**

- Why would we look for top squarks?
- How would top squarks decay and how can we detect them?
- Looking for top squarks in the single-lepton final state
- **Peek at difficult mass spectra**
- Conclusions and outlook

# Top mass corridor

- Most difficult area to target when

$$\Delta M = M_{\tilde{t}_1} - M_{\chi_1^0} = M_t$$

- In rest frame of top squark

- $E_{\tilde{t}_1} = M_{\tilde{t}_1}$  and  $|\vec{p}_{\tilde{t}_1}| = 0$
- LSP also at rest:  
 $E_{\text{LSP}} = M_{\text{LSP}}$  and  $|\vec{p}_{\text{LSP}}| = 0$

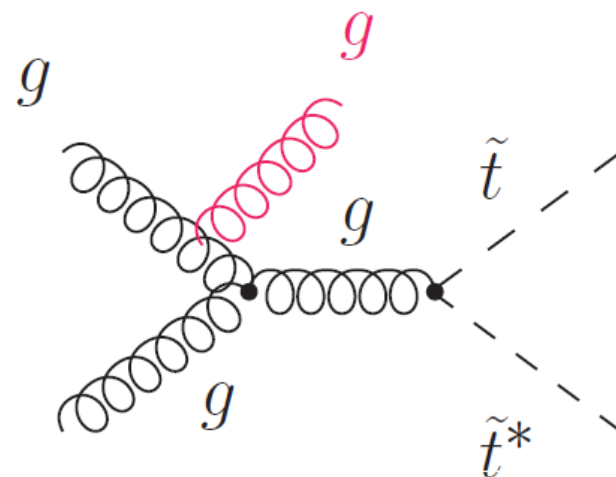
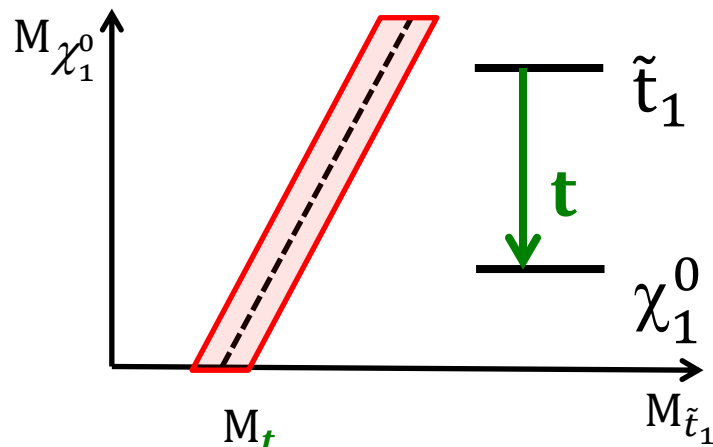
- After Lorentz boost:

- $p_{\text{T,LSP}} = -\gamma v E_{\text{LSP}} = -\gamma v M_{\text{LSP}}$
- $p_{\text{T},\tilde{t}_1} = -\gamma v E_{\tilde{t}_1} = -\gamma v M_{\tilde{t}_1}$

- Combining this information:

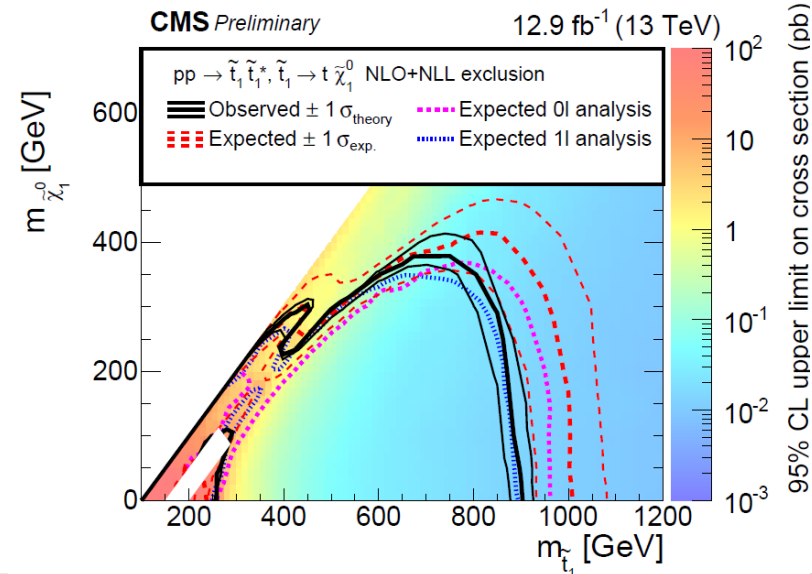
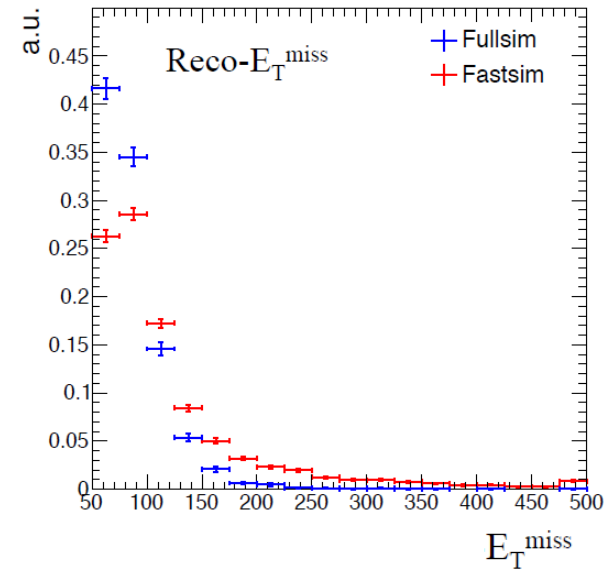
$$p_{\text{T,LSP}} = \frac{M_{\text{LSP}}}{M_{\tilde{t}_1}} p_{\text{T},\tilde{t}_1}$$

1. Boosting the top squarks will lead to large MET
2. For  $M_{\text{LSP}} \ll M_{\tilde{t}_1}$  impossible to get high MET



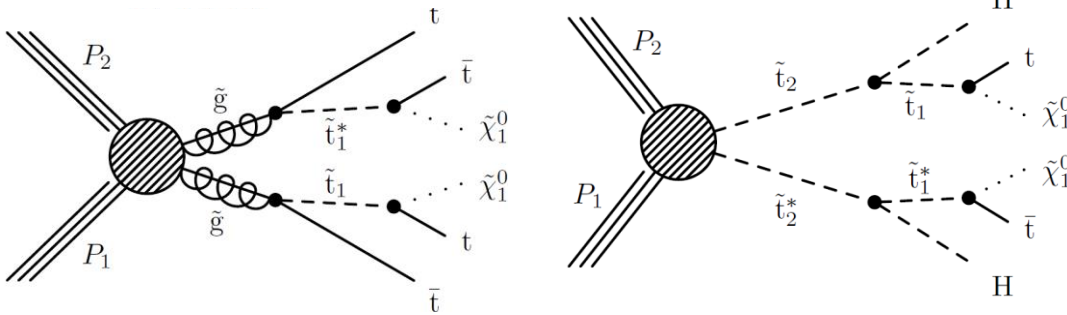
# Top mass corridor

- Dedicated working group to focus on this region:
  - Study the sensitivity in more detail
    - Especially comparing LSP mass of 0 GeV with  $t\bar{t}$  background
  - Signal modeling/simulation
    - Fast simulation had spurious high  $p_T$  jets and much worse MET resolution
- Still a few problems remaining:
  - Different MadGraph settings for background modeling
  - Strong dependence on polarization
- All these effects most sensitive for low LSP masses
- Conclusion:
  - New guidelines for FastSim usage
  - **Light LSPs still needs further study**

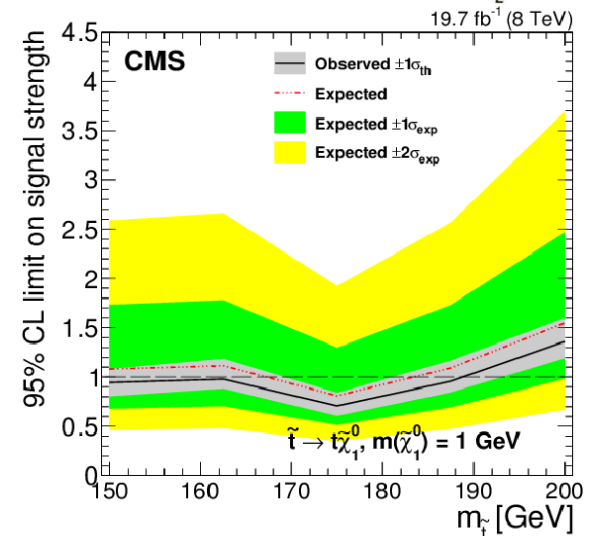
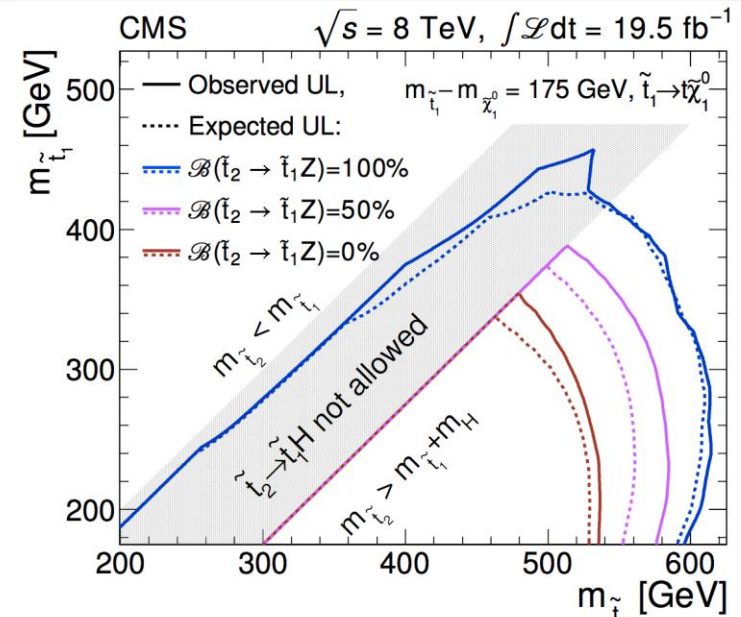


# Other ways to target the top corridor

- Cascade decays
  - Gluinos or heavier stop partners



- Indirect search
  - Excess in  $t\bar{t}$  cross section measurement
  - Measurement of spin correlations between the tops (e.g.  $\Delta\phi(l,l)$ ) since top squarks have spin 0

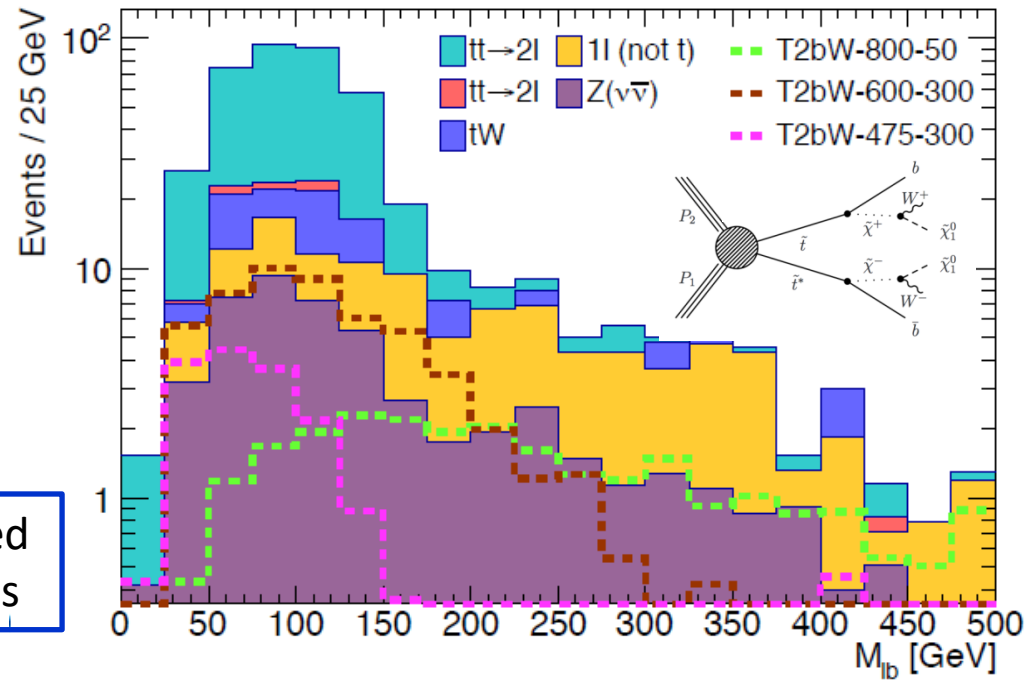


- Why would we look for top squarks?
- How would top squarks decay and how can we detect them?
- Looking for top squarks in the single-lepton final state
- Peek at difficult mass spectra
- **Conclusions and outlook**

# Possible improvements

- Search regions for decays with intermediate chargino
  - More energetic b jets
    - » Use b jet  $p_T$ ,  $M(l,b)$ ,  $N_b=2$  category,...
- Compressed mass splittings:
  - Ask for ISR jet activity
  - Use soft leptons
- Better lost lepton background selection:
  - Start from modified topness
    - » Include detector resolution
    - » Use MET covariance matrix
    - » Softer jets

Dedicated analyses

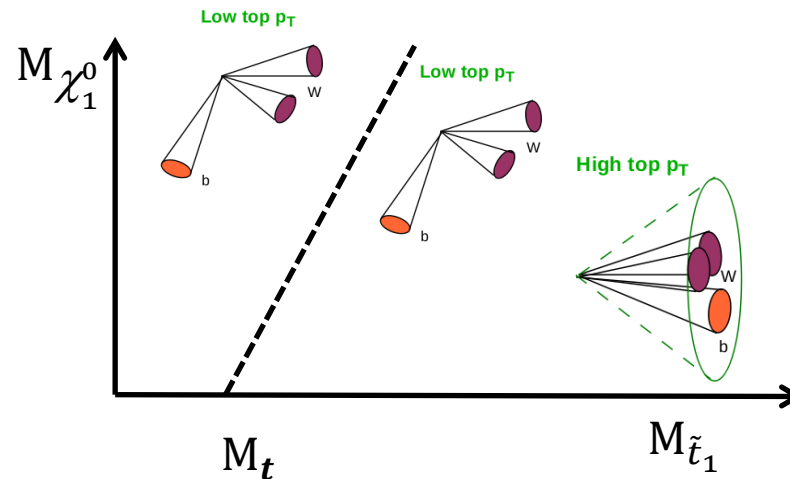


**'theoretical width' parameters**

$$t_{\text{mod}} = \ln(\min S) \quad \text{with} \quad S(\vec{p}_W, p_{\nu,z}) = \underbrace{\frac{(m_W^2 - (p_\nu + p_l)^2)^2}{a_W^4}}_{\text{Constrained to } M_W} + \underbrace{\frac{(m_t^2 - (p_b + p_W)^2)^2}{a_t^4}}_{\text{Constrained to } M_{\text{top}}}$$

# Possible improvements

- Identifying hadronic top decays
  - Both boosted and resolved top quark decays
    - » **Low  $p_T$  top tagging the most important!**
  - Also crucial for all-hadronic top squark search

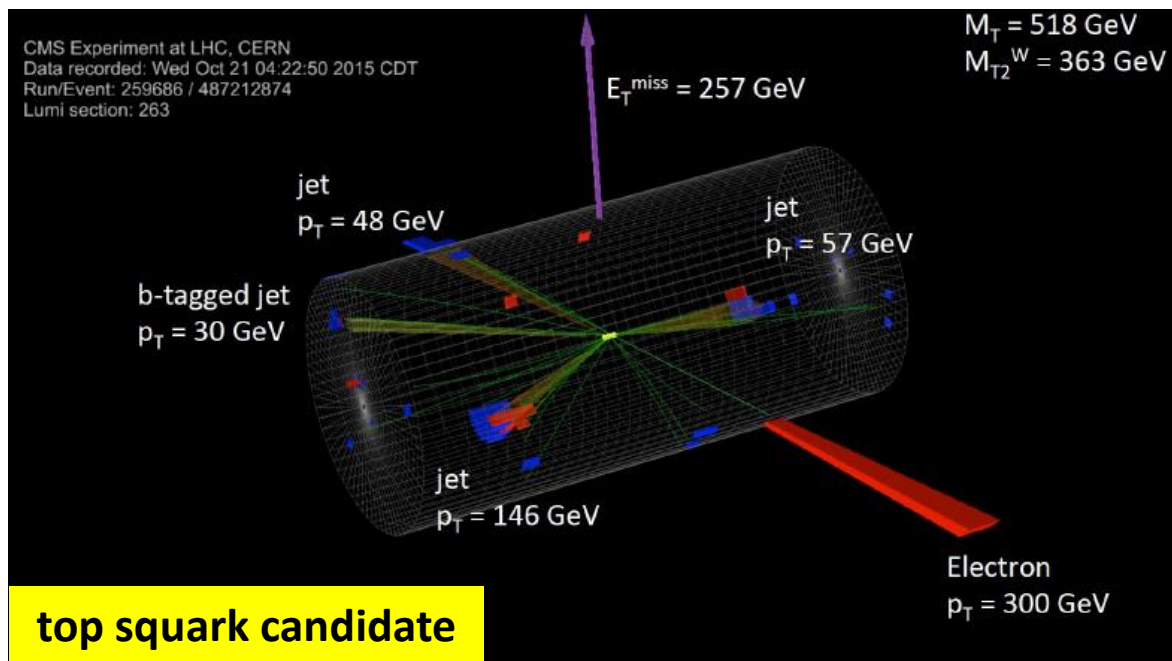


- Improve background prediction
  - Rare backgrounds
  - Current data-driven methods dominated by statistical uncertainties:
    - » Can we replace some by (smaller) systematic uncertainties?



# Conclusions

- Light top squarks provide a clean solution for the hierarchy problem
  - And connects nicely to dark matter searches
- With  $12.9 \text{ fb}^{-1}$  CMS puts stringent limits on top squark masses
  - Start probing 1 TeV mass top squarks
  - Full dataset 3 times larger
- Took a first look at possible improvements



Any questions?



# Back-up

# CMS DETECTOR

Total weight : 14,000 tonnes  
Overall diameter : 15.0 m  
Overall length : 28.7 m  
Magnetic field : 3.8 T

STEEL RETURN YOKE  
12,500 tonnes

SILICON TRACKERS  
Pixel ( $100 \times 150 \mu\text{m}$ )  $\sim 16\text{m}^2 \sim 66\text{M}$  channels  
Microstrips ( $80 \times 180 \mu\text{m}$ )  $\sim 200\text{m}^2 \sim 9.6\text{M}$  channels

SUPERCONDUCTING SOLENOID  
Niobium titanium coil carrying  $\sim 18,000\text{A}$

MUON CHAMBERS  
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers  
Endcaps: 468 Cathode Strip, 432 Resistive Plate Chambers

PRESHOWER  
Silicon strips  $\sim 16\text{m}^2 \sim 137,000$  channels

FORWARD CALORIMETER  
Steel + Quartz fibres  $\sim 2,000$  Channels

CRYSTAL  
ELECTROMAGNETIC  
CALORIMETER (ECAL)  
 $\sim 76,000$  scintillating  $\text{PbWO}_4$  crystals

HADRON CALORIMETER (HCAL)  
Brass + Plastic scintillator  $\sim 7,000$  channels

# Other models

