Imperial College London



Search for HIdden Particles Ulrik Egede, on behalf of SHIP collaboration

Lund, 28 May 2014

A contradiction?

The Standard Model seems perfect and can exist without corrections to the Planck scale

But the same Standard Model can't explain dark matter, neutrino masses and baryogenesis.

I will propose a set of possible ways out of this

28 May 2014 Ulrik Egede 2/51

Outline

The physics landscape

Hidden sector theories

The neutrino minimal Standard Model

Design of a new beam-dump experiment

Sensitivity and future plans

28 May 2014 3/51

Who are we?

www.cern/ch/ship

CERN-SPSC-2013-024 / SPSC-EOI-010 October 8, 2013

Proposal to Search for Heavy Neutral Leptons at the SPS

W. Bonivento^{1,2}, A. Boyarsky³, H. Dijkstra², U. Egede⁴, M. Ferro-Luzzi², B. Goddard², A. Golutvin⁴, D. Gorbunov⁵, R. Jacobsson², J. Panman², M. Patel⁴, O. Ruchayskiy⁶, T. Ruf², N. Serra⁷, M. Shaposhnikov⁶, D. Treille^{2 (‡)}

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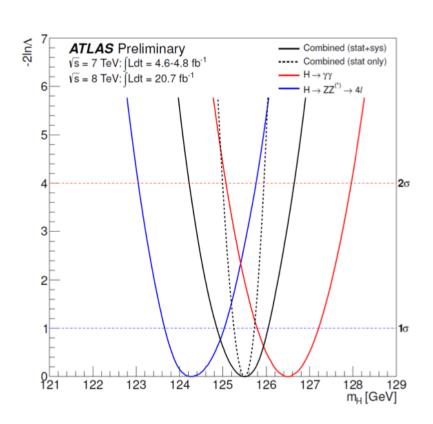
^(‡) retired

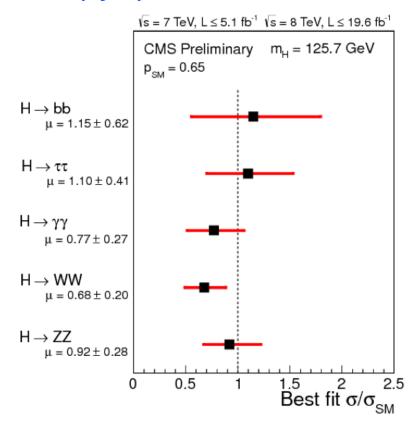
The triumph of the Standard Model

Boson consistent with the SM-Higgs has been found!

ATLAS : $M_{\text{H}} = 125.5 \pm 0.2$ (stat) +0.5-0.6 (syst) GeV/ c^2

CMS : $M_{\perp} = 125.7 \pm 0.3$ (stat) ± 0.3 (syst) GeV/c^2





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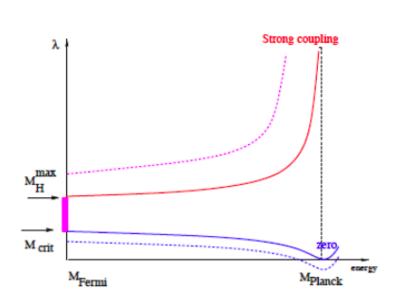
The triumph of the Standard Model

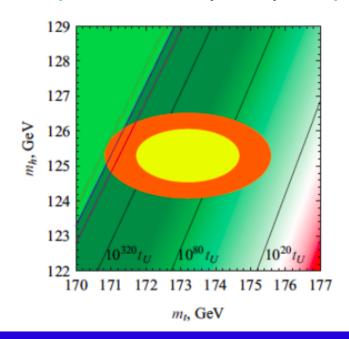
Mass value important for the stability of the vacuum:

$$M_H < 175 \text{ GeV}$$

SM weakly coupled up to the Plank energies!

EW vacuum is stable or metastable with a lifetime greatly exceeding the age of our Universe (JHEP 1208 (2012) 098)





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The limitations of the Standard Model

But we still have a number of significant problems Theory

Radiative corrections to Higgs mass fine-tuning

Experiment

Matter anti-matter asymmetry in the Universe

Neutrino masses and oscillations

Non-baryonic dark matter

Dark Energy

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Direct searches ...

ATLAS SUSY Searches* - 95% CL Lower Limits

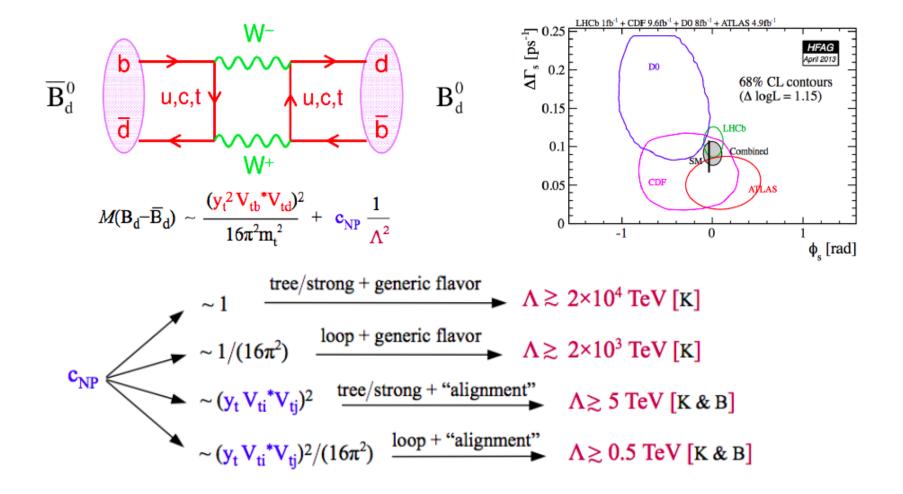
ATLAS Preliminary

5	Status: Moriond 2014						$\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$	\sqrt{s} = 7, 8 TeV
	Model	e, μ, au, γ	Jets	$E_{ m T}^{ m miss}$	$\int \mathcal{L} dt [\mathbf{fl}$	Mass limit	3	Reference
Inclusive Searches	GGM (higgsino-bino NLSP) GGM (higgsino NLSP) Gravitino LSP	$\begin{array}{c} 0 \\ 1 e, \mu \\ 0 \\ 0 \\ 0 \\ 2 e, \mu \\ 2 e, \mu \\ 2 e, \mu \\ 1 \cdot 2 \tau \\ 2 \gamma \\ 1 e, \mu + \gamma \\ \gamma \\ 2 e, \mu (Z) \\ 0 \end{array}$	2-6 jets 3-6 jets 7-10 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 2-4 jets 0-2 jets - 1 b 0-3 jets mono-jet	Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.7 20.7 20.7 20.3 4.8 4.8 5.8	### 1.7 TeV ### 1.2 TeV ### 1.1 TeV ### 740 GeV #### 1.3 TeV #### 1.18 TeV ##### 1.18 TeV ####################################	$ \begin{aligned} & & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & \\ & & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & & & \\ & & & \\ & & & & \\ & &$	ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 1308.1841 ATLAS-CONF-2013-047 ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 ATLAS-CONF-2013-069 1208.4688 ATLAS-CONF-2013-026 ATLAS-CONF-2014-011 ATLAS-CONF-2014-011 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-147
3 rd gen.	$\begin{array}{ccc} \tilde{\mathbf{g}} & \tilde{\mathbf{g}} \rightarrow b \tilde{b} \tilde{\chi}^0_1 \\ \tilde{\mathbf{g}} \rightarrow t \tilde{\lambda}^0_1 \\ \tilde{\mathbf{g}} \rightarrow t \tilde{\lambda}^0_1 \\ \tilde{\mathbf{g}} \rightarrow b \tilde{b} \tilde{\chi}^1_1 \end{array}$	0 0 0-1 <i>e</i> , μ 0-1 <i>e</i> , μ	3 <i>b</i> 7-10 jets 3 <i>b</i> 3 <i>b</i>	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	\$\bar{x}\$ 1.2 TeV \$\bar{x}\$ 1.1 TeV \$\bar{x}\$ 1.34 TeV \$\bar{x}\$ 1.3 TeV	$m(\tilde{k}_1^0)$ <600 GeV $m(\tilde{k}_1^0)$ <350 GeV $m(\tilde{k}_1^0)$ <400 GeV $m(\tilde{k}_1^0)$ <400 GeV	ATLAS-CONF-2013-061 1308.1841 ATLAS-CONF-2013-061 ATLAS-CONF-2013-061
3 rd gen. squarks	$\tilde{t}_1 \tilde{t}_1 \text{ (heavy)}, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$ $\tilde{t}_1 \tilde{t}_1 \text{ (heavy)}, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$	$\begin{array}{c} 0 \\ 2e,\mu (\mathrm{SS}) \\ 1\text{-}2e,\mu \\ 2e,\mu \\ 2e,\mu \\ 0 \\ 1e,\mu \\ 0 \\ 0 \\ \end{array}$	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b nono-jet/c-ta 1 b 1 b	Yes	20.1 20.7 4.7 20.3 20.3 20.1 20.7 20.5 20.3 20.3 20.3	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{split} & m(\vec{k}_1^0) < 90 \text{ GeV} \\ & m(\vec{k}_1^0) = 2 \text{ m}(\vec{k}_1^0) \\ & m(\vec{k}_1^0) = 5 \text{ GeV} \\ & m(\vec{k}_1^0) = m(\vec{k}_1) - m(W) < 50 \text{ GeV}, m(\vec{k}_1) < < m(\vec{k}_1^0) \\ & m(\vec{k}_1^0) = 1 \text{ GeV} \\ & m(\vec{k}_1^0) < 200 \text{ GeV}, m(\vec{k}_1^0) - m(\vec{k}_1^0) = 5 \text{ GeV} \\ & m(\vec{k}_1^0) = 0 \text{ GeV} \\ & m(\vec{k}_1^0) = 0 \text{ GeV} \\ & m(\vec{k}_1^0) - 10 \text{ GeV} \\ & m(\vec{k}_1^0) > 150 \text{ GeV} \\ & m(\vec{k}_1^0) > 150 \text{ GeV} \\ & m(\vec{k}_1^0) < 200 \text{ GeV} \end{split}$	1308.2631 ATLAS-CONF-2013-007 1208.4305, 1209.2102 1403.4853 1403.4853 1403.2631 ATLAS-CONF-2013-037 ATLAS-CONF-2013-034 ATLAS-CONF-2013-068 1403.5222
EW.	$\begin{array}{c} \tilde{\ell}_{L,R}\tilde{\ell}_{L,R},\tilde{\ell} \to \ell \tilde{\chi}^0_1 \\ \frac{1}{\tilde{\chi}^1_1}\tilde{\chi}^1_1,\tilde{\chi}^1_1 \to \tilde{\ell}\nu(\ell\tilde{\nu}) \\ \frac{2}{\tilde{\chi}^1_1}\tilde{\chi}^1_1,\tilde{\chi}^1_1 \to \tilde{\ell}\nu(\ell\tilde{\nu}) \\ \tilde{\chi}^1_1\tilde{\chi}^0_1 \to \tilde{\ell}_L\nu_L^2\ell(\tilde{\nu}\nu),\ell\tilde{\nu}\tilde{\ell}_L\ell(\tilde{\nu}\nu) \\ \tilde{\chi}^1_1\tilde{\chi}^0_2 \to \tilde{\ell}_L\nu_L^2\ell(\tilde{\nu}\nu),\ell\tilde{\nu}\tilde{\ell}_L\ell(\tilde{\nu}\nu) \\ \tilde{\chi}^1_1\tilde{\chi}^0_2 \to W^0_1 X^0_1 \\ \tilde{\chi}^1_1\tilde{\chi}^0_2 \to W^0_1 X^0_1 \end{array}$	2 e, μ 2 e, μ 2 τ 3 e, μ 2-3 e, μ 1 e, μ	0 0 - 0 0 2 <i>b</i>	Yes Yes Yes Yes Yes Yes	20.3 20.3 20.7 20.3 20.3 20.3	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{split} m(\vec{k}_1^0) &= 0 \text{GeV} \\ m(\vec{k}_1^0) &= 0 \text{GeV}, m(\vec{\ell}, \vec{r}) = 0.5 (m(\vec{k}_1^+) + m(\vec{k}_1^0)) \\ m(\vec{k}_1^0) &= 0 \text{GeV}, m(\vec{\ell}, \vec{r}) = 0.5 (m(\vec{k}_1^+) + m(\vec{k}_1^0)) \\ &= m(\vec{k}_2^0), m(\vec{k}_1^0) = 0, m(\vec{\ell}, \vec{r}) = 0.5 (m(\vec{k}_1^+) + m(\vec{k}_1^0)) \\ m(\vec{k}_1^+) &= m(\vec{k}_2^0), m(\vec{k}_1^0) = 0, \text{sleptons decoupled} \\ m(\vec{k}_1^+) &= m(\vec{k}_2^0), m(\vec{k}_1^0) = 0, \text{sleptons decoupled} \end{split}$	1403.5294 1403.5294 ATLAS-CONF-2013-028 1402.7029 1403.5294, 1402.7029 ATLAS-CONF-2013-093
Long-lived	Direct $\tilde{X}_1^{\dagger}\tilde{X}_1^{\dagger}$ prod., long-lived \tilde{X}_1^{\dagger} Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{X}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e$ GMSB, $\tilde{X}_1^1 \rightarrow \gamma \tilde{e}, \tilde{u}) + \tau(e$ GMSB, $\tilde{X}_1^1 \rightarrow \gamma \tilde{e}, \tilde{u})$ GMSB, $\tilde{X}_1^1 \rightarrow qq\mu$ (RPV)	0	1 jet 1-5 jets - - -	Yes Yes - Yes -	20.3 22.9 15.9 4.7 20.3	$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\begin{array}{l} m(\tilde{k}_{1}^{+}) \cdot m(\tilde{k}_{1}^{0}) = 160 \; \text{MeV}, \; \tau(\tilde{k}_{1}^{+}) = 0.2 \; \text{ns} \\ m(\tilde{k}_{1}^{0}) = 100 \; \text{GeV}, \; 10 \; \mu \text{s} < \tau(\tilde{g}) < 1000 \; \text{s} \\ 10 \cdot \tan \beta < 50 \\ 0.4 < \tau(\tilde{k}_{1}^{0}) < 2 \; \text{ns} \\ 1.5 < c\tau < t56 \; \text{mm}, \; BR(\mu) = 1, \; m(\tilde{k}_{1}^{0}) = 108 \; \text{GeV} \end{array}$	ATLAS-CONF-2013-069 ATLAS-CONF-2013-057 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092
RPV	$ \begin{array}{l} LFV \; pp \! \to \! \bar{v}_{\tau} + X, \bar{v}_{\tau} \! \to \! e + \mu \\ LFV \; pp \! \to \! \bar{v}_{\tau} + X, \bar{v}_{\tau} \! \to \! e(\mu) + \tau \\ Billinear \; RPV \; CMSSM \\ \bar{X}_{1}^{\dagger} \bar{X}_{1}^{\dagger}, \bar{X}_{1}^{\dagger} \to \! W \bar{X}_{0}^{\dagger}, \bar{X}_{0}^{\dagger} \! \to \! e \bar{v}_{\mu}, e \mu \bar{v}_{e} \\ \bar{X}_{1}^{\dagger} \bar{X}_{1}^{\dagger}, \bar{X}_{1}^{\dagger} \to W \bar{X}_{0}^{\dagger}, \bar{X}_{1}^{\dagger} \to \tau \tau \bar{v}_{e}, e \tau \bar{v}_{\tau} \\ \bar{X}_{1}^{\dagger} \bar{X}_{1}^{\dagger}, \bar{X}_{1}^{\dagger} \to W \bar{X}_{0}^{\dagger}, \bar{X}_{1}^{\dagger} \to \tau \tau \bar{v}_{e}, e \tau \bar{v}_{\tau} \\ \bar{x}_{2}^{\dagger} \bar{q} q q \\ \bar{g} \to \bar{t}_{1}^{\dagger} t, \bar{t}_{1}^{\dagger} \to b s \end{array} $	$\begin{array}{c} 2e,\mu \\ 1e,\mu+\tau \\ 1e,\mu \\ 4e,\mu \\ 3e,\mu+\tau \\ 0 \\ 2e,\mu(\text{SS}) \end{array}$	- 7 jets - - 6-7 jets 0-3 b	- Yes Yes Yes - Yes	4.6 4.6 4.7 20.7 20.7 20.3 20.7	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	λ'_{31} = 0.10, λ_{132} = 0.05 λ'_{311} = 0.10, λ_{1233} = 0.05 $m(\phi)$ = m(ϕ), c_{135} = 0.1 mm $m(\tilde{k}_1^0)$ > 300 GeV, λ_{121} > 0 $m(\tilde{k}_1^0)$ > 80 GeV, λ_{133} > 0 BR(r) = BR(r) = BR(r) = 0%	1212.1272 1212.1272 ATLAS-CONF-2012-140 ATLAS-CONF-2013-036 ATLAS-CONF-2013-036 ATLAS-CONF-2013-091 ATLAS-CONF-2013-007
Other	_	0 2 e, μ (SS) 0	4 jets 2 b mono-jet		4.6 14.3 10.5	sgluon 100-287 GeV sgluon 350-800 GeV M* scale 704 GeV	incl. limit from 1110.2693 $ m(\chi) < 80 \; {\rm GeV, limit \; of} < 687 {\rm GeV \; for \; D8} $	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
		\sqrt{s} = 8 TeV partial data	$\sqrt{s} = 0$	8 TeV data		10^{-1} 1	Mass scale [TeV]	

^{*}Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 or theoretical signal cross section uncertainty.

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... and indirect searches



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A hidden sector?

Rather than being heavy, could new particles be light but very weakly interacting?

A new, light "hidden sector" of particles

Singlets with respect to gauge group of the SM

How could we have missed it?

Key is that it only interacts with SM particles through some kind of mixing through a "portal" particle

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A hidden sector?

Several possibilities for renormalisable singlet operators

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Vector portal, U(1) B_{\mu\nu}
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Massive vector photon (paraphoton, secluded photon...) mixing with regular photon $\rightarrow \epsilon B_{\mu\nu} F^{\mu\nu}$

Higgs portal

Scalar field χ , $(\mu \chi + \lambda \chi^2)H^{\dagger}H$

Axial portal

Pseudo Nambu-Goldstone bosons

Axion like vector field , (a/F) $G_{\mu\nu}G^{\mu\nu}$, $(\partial_{\mu}a/F)\psi^{\dagger}\gamma_{\mu}\gamma_{5}\psi$

Neutrino portal

Heavy neutral leptons (HNL), YH^TN'L

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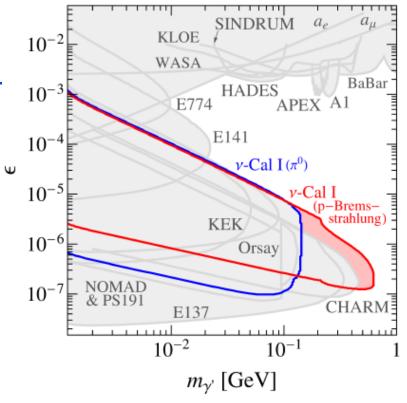
Vector Portal

Produce a virtual photon in fixed-target, mix into a dark photon ($\sim\epsilon$), dark photon, mix back into a SM photon ($\sim\epsilon$), decay into e⁺e⁻, μ ⁺ μ ⁻, π ⁺ π ⁻ etc

No other interactions with SM particles - "light-shining-through-awall" experiments

Constraints from a wide range of experiments – note results from CERN CHARM experiment

[arXiv:1311.3870]



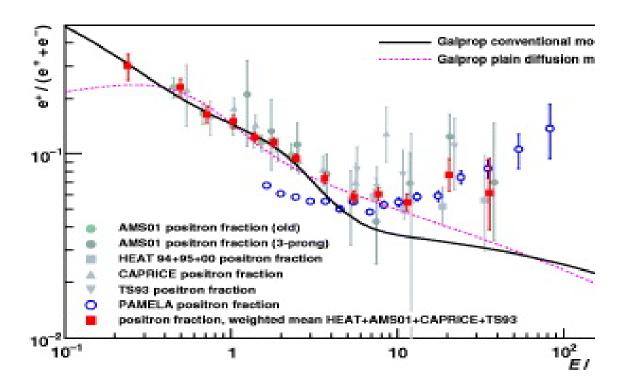
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Vector Portal

PAMELA cosmic ray experiment observes a sharp upturn in the positron fraction

DM annihilation?

Requires large cross section, but not in hadron signature



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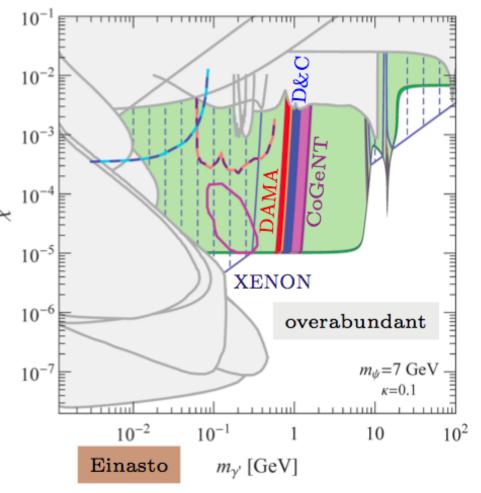
Vector Portal

ATIC, WMAP, EGRET all suggest excess electronic production

INTEGRAL 511 keV suggests far more annihilation in the galactic centre than can be explained by supernovae

Annihilation can be into a new force carrier m~1 GeV

Low mass can make hadronic modes kinematical inaccessible, forcing decay dominantly into leptons



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Higgs Portal

arXiv:1403.4638

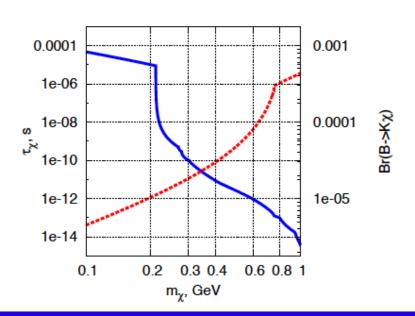
Example of inflaton

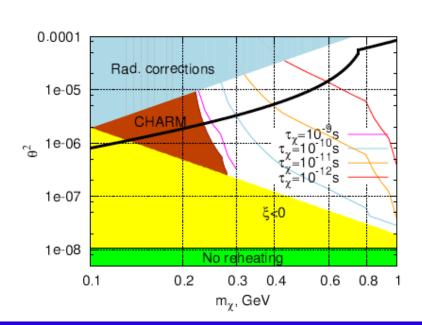
Together with Higgs, generates inflation of the early Universe

Model has a 7 keV (warm) DM candidate and respects constraints from BICEP2 and Planck

Interesting mass region 0.3 GeV $< m_{\chi} < 1$ GeV

Little experimental exploration of interesting region...

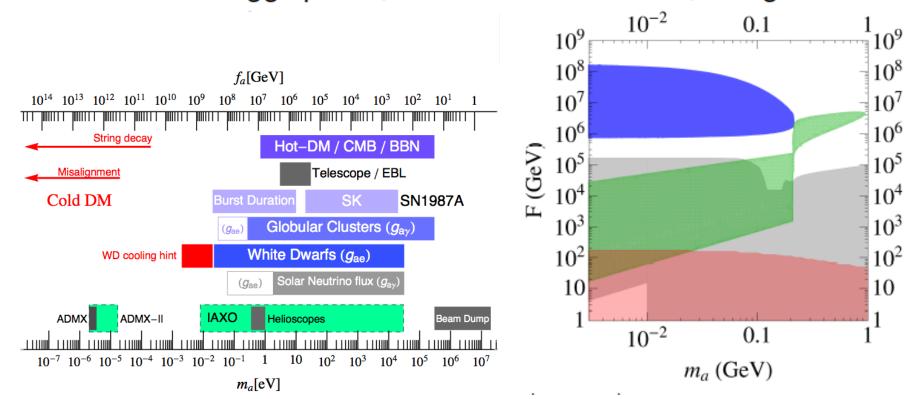




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Axion Portal

Similar to the Higgs portal, from arXiv:1008.0636, Essig et al



Right: Gray: the combined exclusion region from meson decays; green: CHARM; blue: supernova SN 1987a; red: muon anomalous magnetic moment.

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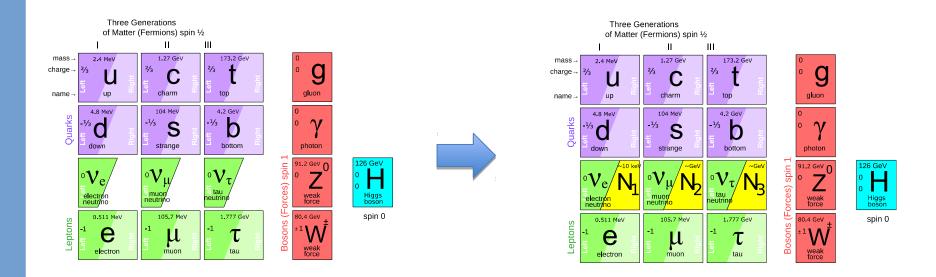
Neutrino Portal

[T.Asaka, M.Shaposhnikov, Phys. Lett B620 (2005) 17

The neutrino Minimal Standard Model (vMSM) aims to explain

Matter anti-matter asymmetry in the Universe, neutrino masses and oscillations, non-baryonic dark matter

Adds three right-handed, Majorana, Heavy Neutral Leptons (HNL), N_1 , N_2 and N_3



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Neutrino Portal

[T.Asaka, M.Shaposhnikov, Phys. Lett B620 (2005) 17

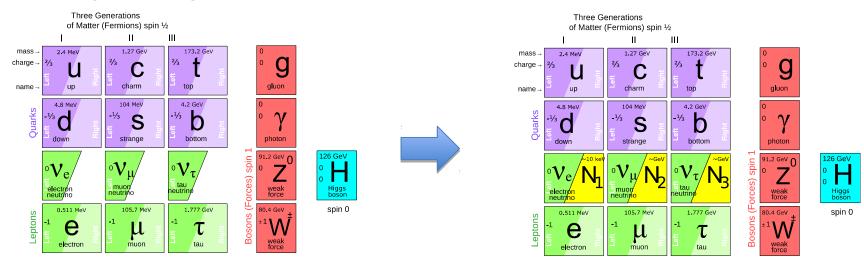
 N_1

Mass in keV region, (warm) dark matter candidate

 $N_{2,3}$

Mass in 100 MeV – GeV region

Generate neutrino masses via see-saw and produce baryon asymmetry of the Universe



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See-saw for v mass

Most general renormalisable Lagrangian of all SM particles (+3 singlets wrt the SM gauge group):

$$L_{\rm singlet} = i \bar{N}_I \partial_\mu \gamma^\mu N_I - Y_{I\alpha} \bar{N}_I^c \tilde{H} L^c_\alpha - M_I \bar{N_I}^c N_I + {\rm h.c.},$$

Yukawa term: mixing of N₁ with active neutrinos to explain oscillations

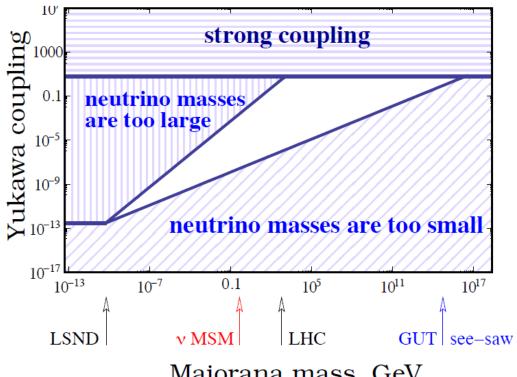
Majorana term which carries no gauge charge

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See-saw for v mass

The scale of the active neutrino mass is given by the seesaw formula, m₀=m₀²/M

Typical value of the Dirac mass term is linked to the Yukawa coupling of the I-th neutrino by m_D~Y_{Ia}v



Majorana mass, GeV

20/51 28 May 2014

Constraints on N₁ as DM

Stability

Must have a lifetime larger than that of the Universe

Production

Created in the early Universe in reactions $I^+I^- \rightarrow \nu N_1$, $qq \rightarrow \nu N_1$ etc.

Need to provide correct DM abundance

Structure formation

Should be heavy enough to not erase non-uniformities at small scales

Decay

Should not produce decays we have already excluded!

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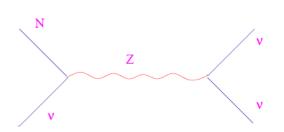
N₁ – dark matter candidate

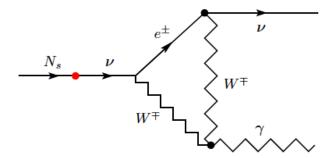
Small Yukawa couplings mean that N₁ can be very stable

$$au_{N_1}=10^{14}\, ext{years}\left(rac{10\, ext{keV}}{M_N}
ight)^5\left(rac{10^{-8}}{ heta_1^2}
ight) \qquad \qquad heta_1=rac{m_D}{M_N}$$

Main decay mode $N \rightarrow 3v$, clearly unobservable

Subdominant radiative decay $N \rightarrow v\gamma$ would give a monoenergetic photon with $E_v = M_N/2$

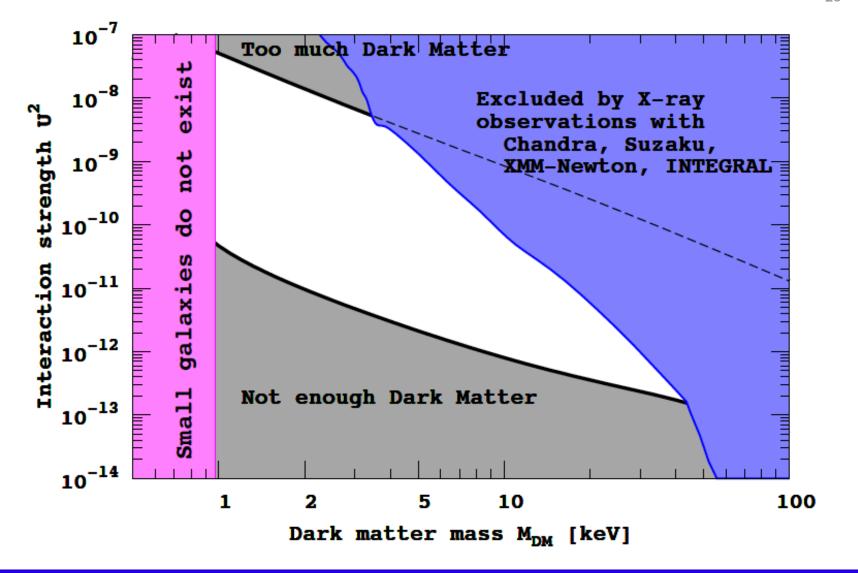




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N₁ allowed parameter space





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Space-based searches for N₁

Existing searches by XMM-Newton, Chandra, Suzaku, INTEGRAL

Relatively poor spectral resolution – need $\Delta E/E \sim 10^{-3}$ to separate from lots of line from iron

Four future satellite missions planned:

Astro-H – approved for 2015 launch but has rather limited field of view

Athena+

LOFT

Origin/Xenia ... timescale ~2020

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New line in galaxy spectrum?

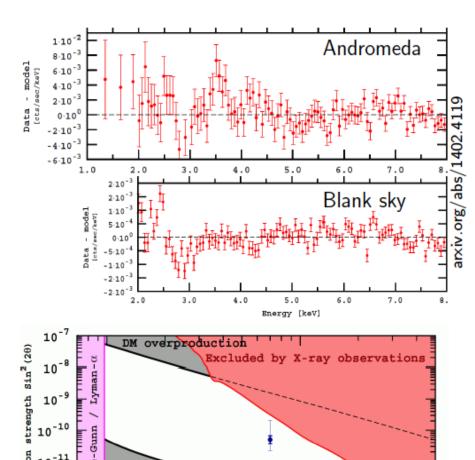
ArXiv:1402.4119

An unidentified line in the xray spectrum of the Andromeda galaxy and Perseus galaxy cluster E_y = ~3.5 keV

ArXiv:1402.2301

Detection of an unidentified emission line in the stacked x-ray spectrum of galaxy clusters E_v =~3.56 keV

Astro-H will be able to check these claims with better energy resolution



Dark matter mass M_{DM} [keV]

Not enough DM

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Generating the baryon asymmetry

CP is not conserved in the vMSM:

6 new CP-violating phases in lepton sector

Process for Baryon asymmetry

HNL are created in the early Universe

CPV in the interference of HNL production and decay

Lepton number asymmetry goes from HNL to active neutrinos

Asymmetry transferred to baryons via "sphaleron processes"

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N_{2,3} production and decay

 $M(N_2) \approx M(N_3) \sim$ a few GeV \rightarrow can dramatically increase amount of CPV to explain Baryon Asymmetry of the Universe (BAU)

Explanation of DM with N₁ reduces number of free parameters, need degeneracy to ensure sufficient CPV

Very weak $N_{2,3}$ to ν mixing ($\sim U^2$) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles

Typical lifetimes >10 ms for $(N_{2,3})$ ~ 1 GeV → decay distance O(km) too large U erases any BAU

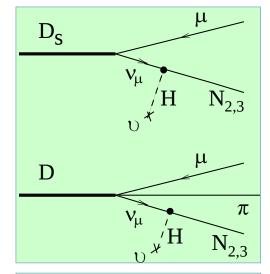
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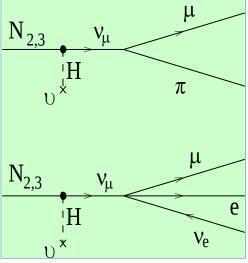
N_{2,3} production and decay

Production in charm decays...

... and decays to lighter SM particles

For what follows will focus on $N \rightarrow \mu - \pi +$ decay, BF~0.1-50% Similar BF for $N \rightarrow \mu^- \rho^+$, BF($N \rightarrow \nu \mu e$) ~ 1-10%

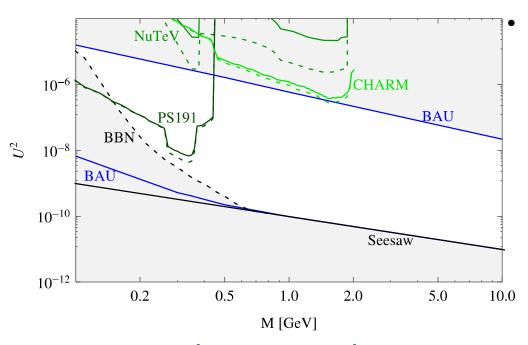




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Experimental and cosmological constraints

BAU, See-saw and Big Bang Nucleosynthesis (BBN) constraints indicate that previous experiments probed the interesting region only below the kaon mass:



Previous searches:

- PS191('88)@PS 19.2 GeV
 1.4×10^{19 pot, 128} m from target
- CHARM('86)@SPS 400 GeV, 2.4×10^{18 pot, 480} m from target
- NuTeV('99)@Fermilab 800 GeV, 2.5×10^{18} pot, 1.4 km from target

BBN, BAU and Seesaw give stronger constraints than experimental searches for $M_N > 400 \text{ MeV}$

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Experimental and cosmological constraints

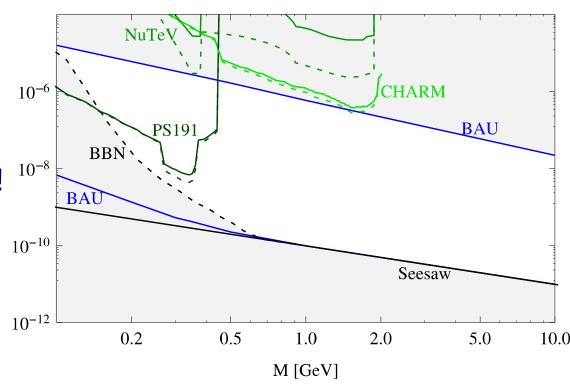
BAU, Seesaw and Big Bang Nucleosynthesis (BBN) constraints indicate that previous experiments probed the interesting region only below the kaon mass:

Mixing at both production and decay

BF(D \rightarrow NX) around $10^{-8} - 10^{-12}$

Lifetime can give further factor 10⁻⁴

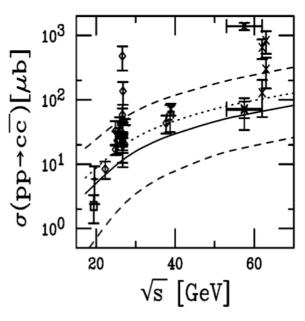
Need >10¹⁶ D meson⁵! ¹⁰⁻⁸

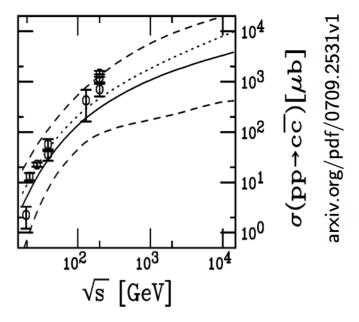


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Where to produce charm?

cc cross-section:





LHC ($\sqrt{s} = 14 \text{ TeV}$)

1 ab⁻¹ (i.e. 3-4 years), $\sim 2 \times 10^{16}$ in 4π

SPS (400 GeV p-on-target (pot) $\sqrt{s} = 27$ GeV)

 2×10^{20} pot (i.e. 3-4 years): ~ 2×10^{17}

Fermilab: 120 GeV, 10× smaller σ_{cc} , 10×pot by 2025 for

Could mass range be extended?

Would neutrinos from B-decays extend the mass range for $N_{2,3}$ upwards?

Produced with factor 20-100 smaller cross-section

Dominant semi-leptonic decay B → Dμν

Similar limit for neutrino mass

Charmless $B \rightarrow \pi \mu \nu$ heavily Cabibbo suppressed

B decays are not at all competitive

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Experimental design

Propose a beam dump experiment at the CERN SPS with a total of $\sim 2 \times 10^{20}$ protons on target

Crucial design parameters:

Minimise residual neutrino and muon fluxes

Can produce K^o that decay in detector and mimic signal events

Short-lived resonances generate 109 muons/spill

Muon shield

Neutrinos from light meson decays

Dense target/hadron absorber

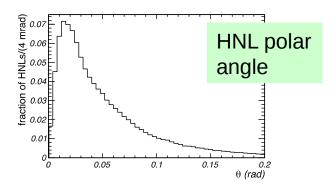
Prevent neutrino interactions from mimicking HNL decay

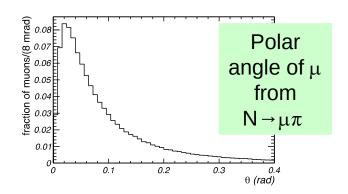
Evacuate decay volume

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Experimental design

HNLs produced in charm decays have significant p_{τ}





Detector must be close to target to maximise geometrical acceptance

Shielding for muons must be as short as possible

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Experiment

Secondary beam line

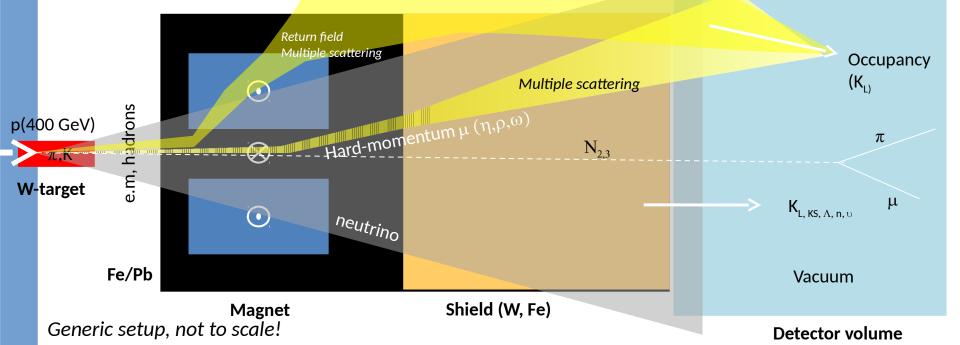
Initial reduction of beam induced backgrounds

- Heavy target (50 cm of W)
- Hadron absorber
- Muon shield: optimization of active and passive shields is underway Low-mid-momentum μ from fast decays of π,Κ

Acceptable occupancy <1% per spill of $5\times10^{13 p.o.t.}$

spill duration 10ms [] < 50×103 muons

spill duration $10\mu s$ II < 500 muons



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Muon Shield

Without μ -filter: 5×10^9 / spill (5×10^{13} pot)

Idea to reduce background from μ -interactions to below ν -background

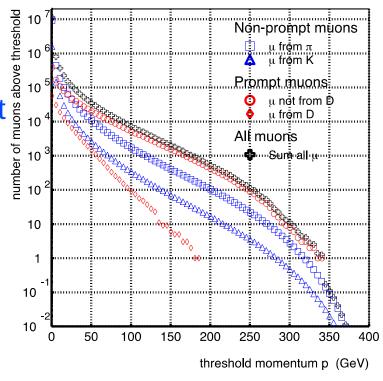
Acceptable rate $\sim 10^5 \,\mu$ / 2×10^{20} pot

Main sources of muons simulated using PYTHIA

Two alternatives for shield:

Passive: i.e. use high Z material: need 54 m of W to stop 400 GeV μ

Active (+passive): need 40 Tm to deflect 400 GeV μ outside acceptance



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Neutrino backgrounds

Neutrino interactions in the decay volume :

After shield expect 2×10⁴ per 2×10²⁰ pot at atmospheric pressure

Negligible at 0.01 mbar

Neutrino interactions in the final part of the muon shield:

Use GEANT and GENIE to simulate the CC and NC neutrino interactions

CC(NC) rate of ~6(2)×10⁵ per interaction length per 2×10²⁰ pot

Use veto-station to suppress short lived

 v_{μ} + p \rightarrow X + K_{L} \rightarrow $\mu \pi \nu$ main background

Requiring μ -id. for one of the two decay products

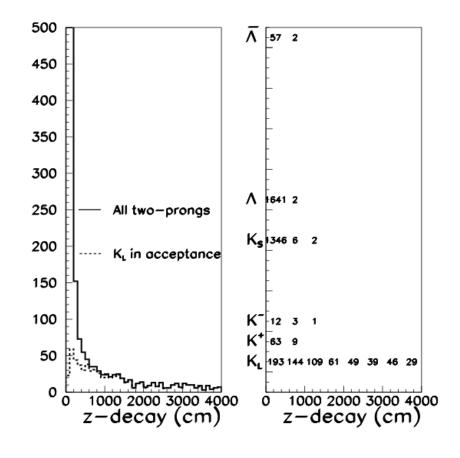
→ 150 two-prong vertices in 2×10²⁰ pot

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Neutrino backgrounds

Neutrino interactions in the decay volume

~10% of neutrino interactions in the muon shield just upstream of the decay volume produce Λ or K^0

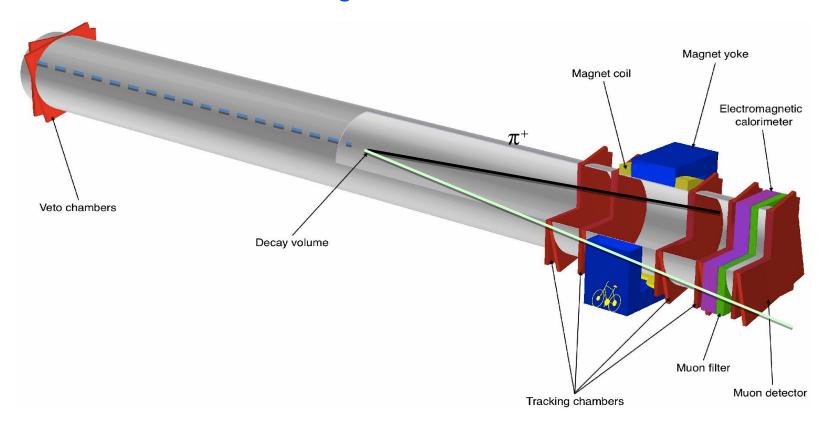


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Detector Concept

Aim to reconstruct HNL decays into the final states: $\mu^-\pi^+$, $\mu^-\rho^+$, $e^-\rho^+$

Require long decay volume, magnetic spectrometer, muon detector and electromagnetic calorimeter



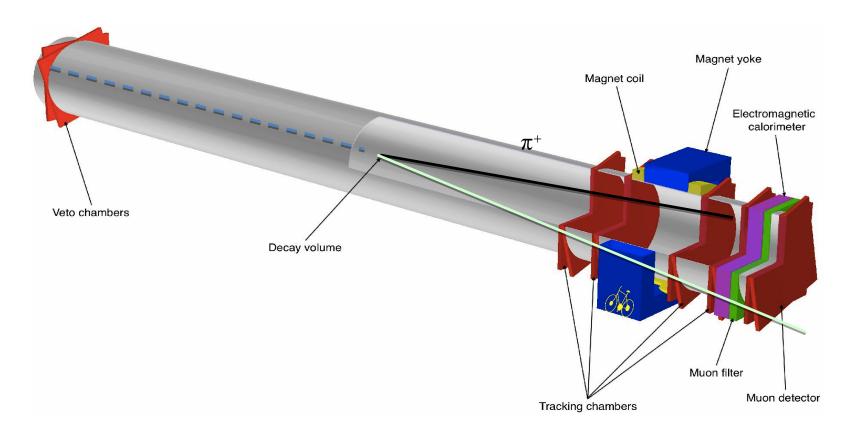
μ

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Detector Concept

5 m diameter, 50 m length vacuum vessel

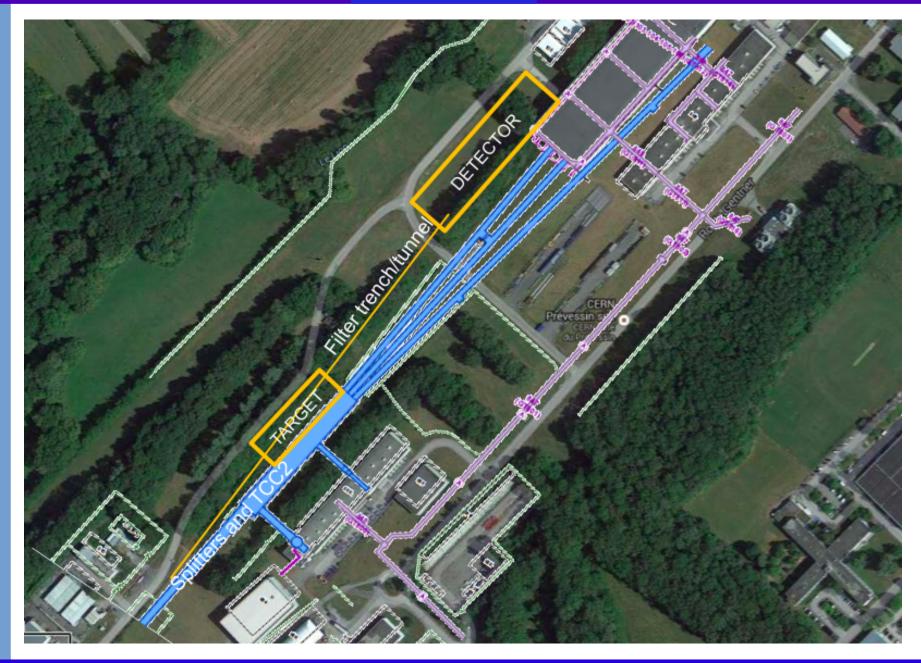
10 m long magnetic spectrometer with 0.5 Tm dipole magnet and four tracking chambers



μ

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Experiment



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Detector Technologies

Dipole magnet

Magnet similar to LHCb design required, but with ~40% less iron and 3× less power

Free aperture of ~16 m²

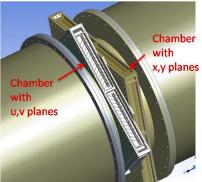
Field integral ~ 0.5 Tm over 5 m length

Vacuum tank and straw tracker

NA62 has 10⁻⁵ mbar pressure, only 10⁻² mbar here

Have demonstrated gas tightness of straw tubes with 120 μm spatial resolution and 0.5% X^0 material budget in long term tests







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Detector Technologies

Electromagnetic calorimeter

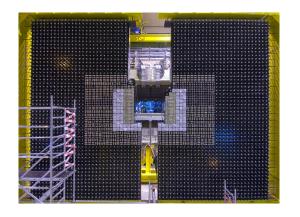
Shashlik technology used in LHCb would provide economical solution with good energy and time resolution

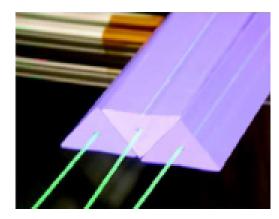
Muon detector

Scintillator strips with WLS fibres and Silicon Photomultiplier (SiPM) an attractive option

Trigger and DAQ

Requirements on both are very modest due to low data rate



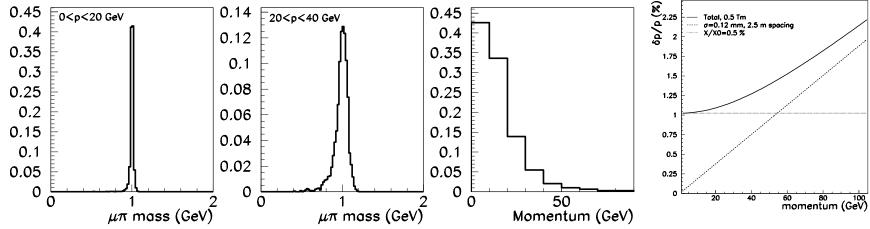


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Spectrometer resolution

Arrange spectrometer such that multiple scattering and spatial resolution of straw tubes give similar contribution to the overall $\Delta p/p$

For m(N_{2,3}) = 1 GeV, 75% of $\mu^-\pi^+$ decay products have p < 20 GeV



For 0.5 Tm field integral $\sigma(mass) \sim 40$ MeV for p < 20 GeV

Good discrimination between high mass tail from small number of residual $K_{_{\! I}} \to \mu^{_{\! T}} \pi^{_{\! T}} \upsilon$ and 1 GeV HNL

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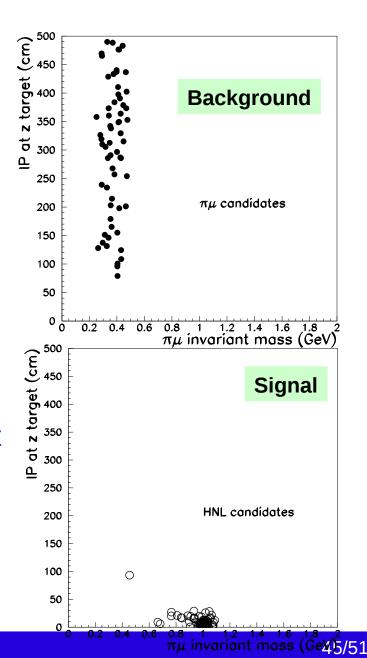
Residual Backgrounds

K_L produced in the final part of the muon shield have very different pointing to the target compared to signal events

Use Impact Parameter (IP) to further suppress K₁ background

IP < 1 m is 100% efficient for signal and leaves only a handful of background events

The IP cut will also be used to reject backgrounds induced in neutrino interactions in the material surrounding the detector



Expected event yield

Integral mixing angle U^2 is given by $U^2 = U_e^2 + U_\mu^2 + U_\tau^2$

Make a conservative estimate of the sensitivity by only considering the decay $N_{2,3} \rightarrow \mu^- \pi^+ - \text{probes } U_\mu^{-2}$

Expected number of signal events then,

$$N_{\text{signal}} = n^{\text{pot}} \times 2\chi_{\text{cc}} \times BR(U_{\mu}^{2}) \times \epsilon_{\text{det}}(U_{\mu}^{2})$$

Strongest experimental limit for $M_N \sim 1$ GeV at $U_{\mu}^{2} = 10^{-7}$

Would then expect $\tau_N = 1.8 \times 10^{-5} \, \text{s}$ and ~12k fully reconstructed N $\rightarrow \mu^- \pi^+$

For cosmologically favoured region $U_{\mu}^{2} = 10^{-8} (\tau_{N} = 1.8 \times 10^{-4} \text{ s})$

Would expect 120 fully reconstructed events

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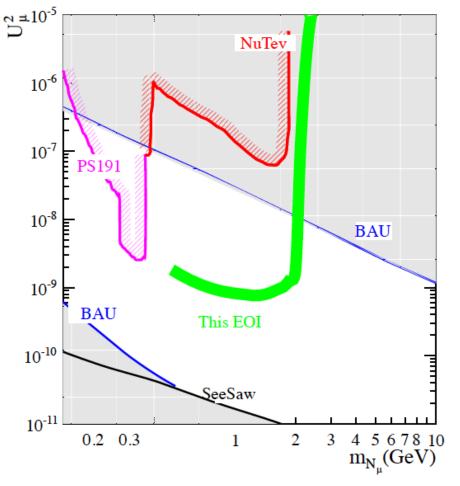
Expected sensitivity

For M_N < 2 GeV the proposed experiment has discovery potential for the cosmologically favoured region with

 $10^{-7} < U_{\mu}^{2} < a \text{ few} \times 10^{-9}$

Limit from decay channels with electromagnetic not studied yet.

Will extend search on U^2 and limit on U_a^2



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Status of the SPSC review

Submitted our EOI in Oct 2013 [CERN-SPSC-2013-024 / SPSC-EOI-010 / arXiv:1310.1762]

SPSC assigned four referees – provided answers to their questions [http://ship.web.cern.ch/ship/EOI/SPSC-EOI-010_ResponseToReferees.pdf]

SPSC discussed our proposal Jan 2015, official feedback

"The Committee received with interest the response of the proponents to the questions raised in its review of EOI010.

The SPSC recognises the interesting physics potential of searching for heavy neutral leptons and investigating the properties of neutrinos.

Considering the large cost and complexity of the required beam infrastructure as well as the significant associated beam intensity, such a project should be designed as a general purpose beam dump facility with the broadest possible physics programme, including maximum reach in the investigation of the hidden sector.

To further review the project the Committee would need an extended proposal with further developed physics goals, a more detailed technical design and a stronger collaboration."

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A strengthening collaboration

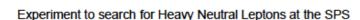


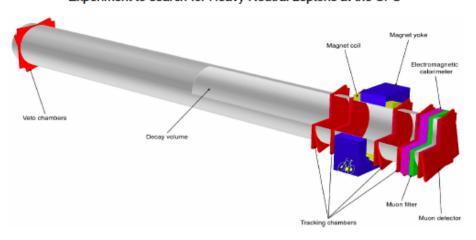






CERN, Universität Zürich, EPFL Lausanne, INFN Cagliari, Università Federico II and INFN Napoli, Imperial College London





We propose a new fixed-target experiment at the CERN SPS accelerator to search for hidden particles. In particular, to search for Heavy Neutral Leptons (HNLs) produced in charm decays. HNLs are right-handed partners of the Standard Model neutrinos. The existence of such particles is strongly motivated by theory, as they can simultaneously explain the baryon asymmetry of the Universe, account for the pattern of neutrino masses and oscillations, and provide a Dark Matter candidate.

SHIP is a collaboration of six institutes: CERN, Universität Zürich, École Polytechnique Fédérale de Lausanne, INFN Sezione di Cagliari, Università Federico II and INFN Napoli, Imperial College London. Groups interested in joining should contact Andrey Golutvin and Jaap Panman. The extension of the collaboration will be discussed at the First SHIP Workshop that will be take place in Zürich the 10-12 June 2014.

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First SHiP Workshop, 10-12th June 2014

Meeting at University of Zurich to discuss:

Physics reach of SHiP detector

Detector requirements and technologies

Please get in touch if you are interested ...

Theoretical Overview (10th June)

Review of heavy neutral leptons, with discussions about leptogenesis and cosmological constraints

Theory review (11th June Morning)

Discussion of theoretical status and present experimental constraints

Facility and Experiment (11th June Afternoon)

Discussion on the primary beam line, target and detector design for the SHIP experiment

Tau neutrinos and SHIP detector (12th June Morning)

Discussion on the electronics and DAQ system for the SHIP experiment and on the detector for tau neutrinos

Summary and discussion (12th June Afternoon)

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Conclusions

The proposed experiment will search for NP in the largely unexplored domain of new, very weakly interacting particles with masses below the Fermi scale

Detector is based on existing technologies

Ongoing discussion of the beam line with CERN experts

The discovery of a HNL would have enormous impact – could solve several of the significant problems of the SM

The origin of the baryon asymmetry of the Universe

The origin of neutrino mass

The results of this experiment, together with cosmological and astrophysical data, could be crucial to determine the nature of Dark Matter

Wide range of other hidden sector physics under investigation

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